Dissertation zur Erlangung des Doktorgrades der Fakultät für Chemie und Pharmazie der Ludwig-Maximilians-Universität München

# NH<sub>4</sub>F-Assisted High-Pressure Synthesis of Novel Nitrides and Atomic-Resolution Chemical Mapping by STEM

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#### Erklärung

Diese Dissertation wurde im Sinne von § 7 der Promotionsordnung vom 28. November 2011 von Herrn Prof. Dr. Wolfgang Schnick betreut.

#### **Eidesstattliche Versicherung**

Diese Dissertation wurde eigenständig und ohne unerlaubte Hilfe erarbeitet.

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Für alle, die mich zu dem gemacht haben, der ich heute bin.

"Es wird ja fleißig gearbeitet und viel mikroskopiert, aber es müßte mal wieder einer einen gescheiten Gedanken haben." Rudolf Ludwig Carl Virchow (\*1821–†1902), Physician

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# 1 Introduction

### 1.1 Nitride Materials

The ongoing increase in population demands new technologies to bring into line the finiteness of resources and the claim to, at least, maintain overall wealth. The role of fundamental science, in this case, is the discovery of new materials suited to future technologies. The compound class of nitrides attracted interest as their versatility spans many fields of applications. Silicon nitride  $Si_3N_4$  is used in the valves of high-performance combustion engines, in ball bearings, as an abrasive material, and as a passivating layer in integrated circuits.<sup>[1-4]</sup> This range of applications is enabled by the covalent character and therefore strong Si–N bond with a molar bond energy of 335 kJ·mol<sup>-1[5]</sup> as well as a high degree of condensation ( $\kappa^*$ ) of the underlying tetrahedral network which combines to a material of high hardness with thermal as well as chemical stability. If we increase the complexity by introducing at least a third element we obtain a compound class called nitridosilicates. Noteworthy representatives of nitridosilicates are the compounds  $AE_2Si_5N_8$  (AE = Ca, Sr, Ba) which, upon doping with Eu<sup>2+</sup> and irradiation with blue light, emit orange light.<sup>[6]</sup> Based on the phosphor-converted InGaN LED this led to the mass production of the first warm white pc-LED (phosphor-converted light-emitting diode) marketed to a wide consumer base. The aforementioned properties like chemical and thermal stability are similarly beneficial even for these more elaborate applications.

Another binary main group nitride,  $P_3N_5$ , although having a higher degree of condensation lacks thermal stability.  $P_3N_5$  readily decomposes at 850 °C<sup>[7]</sup> under ambient pressure which can be attributed to the weaker P–N bond with a molar bond energy of 290 kJ·mol<sup>-1[5]</sup> if compared to the previously mentioned Si–N bond. Despite the differences, there are likewise multinary compounds called nitridophosphates.

The last binary group of nitrides relevant for this thesis are the interstitial transition metal nitrides, for example, TiN and ScN. These compounds show remarkable thermal stability and in the case of TiN also mechanical stability hence its use as a coating material for drill heads and also bicycle chains.<sup>[8]</sup> For a comprehensive overview of some applications of nitride-based materials see Figure 1.1.

<sup>&</sup>lt;sup>\*</sup>  $\kappa$  is defined as the ratio of the number of tetrahedra centers and the number of coordinating atoms. For example, Sr<sub>2</sub>Si<sub>5</sub>N<sub>8</sub> has a  $\kappa$  of 5/8.

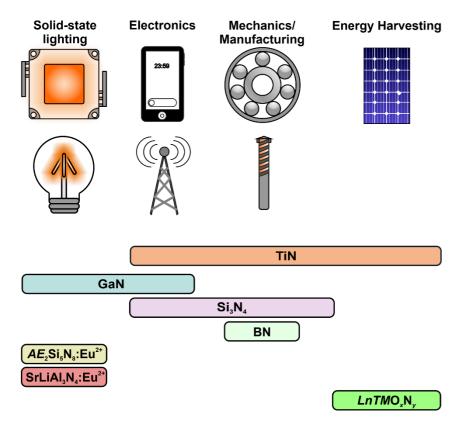


Figure 1.1. Applications of some nitridic materials ranging from solid-state lighting to manufacturing.

### **1.2** Syntheses of Nitridosilicates and -phosphates

The synthesis of Si<sub>3</sub>N<sub>4</sub> is either performed from the elements or by the imide-decomposition process by the pyrolysis of Si(NH)<sub>2</sub>.<sup>[9,10]</sup> The synthesis of nitridosilicates is likewise readily performed by conventional solid-state methods such as the heating of the starting materials in a tube- or either radio-frequency furnaces delivering the activation energy needed to surpass the stability of Si<sub>3</sub>N<sub>4</sub>.<sup>[6,11-14]</sup> The metal-bearing starting materials, according to literature, range from nitrides, hydrides, amides to even the elemental form.<sup>[15]</sup> The limits of this intuitive approach however are reached when transition metals shall be incorporated. As mentioned in literature the transition metals are prone to reduction while the nitride-ion is in turn oxidized.<sup>[16,36]</sup> A problem expected to become more severe as the oxidation state of the metal is increased. Bielec et. al. synthesized the compound Fe<sub>2</sub>Si<sub>5</sub>N<sub>8</sub> from Ca<sub>2</sub>Si<sub>5</sub>N<sub>8</sub> and FeCl<sub>2</sub> in an ion-exchange type of reaction.<sup>[16]</sup> It is suggested that Ca<sup>2+</sup> migrates through channels in the nitridic network, upon heating, to form CaCl<sub>2</sub> while cation sites are refilled by Fe<sup>2+</sup>. Bielec et al. not only managed aliovalent substitution, in the case of Sc<sub>0.2</sub>Ca<sub>1.7</sub>Si<sub>5</sub>N<sub>8</sub> but also avoided the use of refractory ScN by their ion-exchange approach.<sup>[17]</sup> However, a drawback of this approach is its dependence on the preformed network if only topotactical reactions are

considered as "ion-exchange" reactions. Wendl et al. demonstrated post-synthetic modifications, via metathesis, altering the underlying nitridic network during the reaction.<sup>[18]</sup> The synthesis of  $P_3N_5$  is performed at 850 °C from  $P_4S_{10}$  in flowing ammonia, which yields an amorphous bulk sample as described by Stock and Grüneberg as early as 1907.<sup>[19]</sup> Suitable samples for powder X-ray diffraction structure elucidation from synchrotron radiation were achieved by Horstmann et al. only after the pyrolysis of [P[NH<sub>2</sub>]<sub>4</sub>]I in 1997.<sup>[20]</sup> Further research yielded compounds like the nitridic sodalites  $Zn_8[P_{12}N_{24}]X_2$  (X = 0, S, Se, Te), and quasi-binary HPN<sub>2</sub> from syntheses in tube furnaces but the success story of nitridophosphates was yet to come.<sup>[21-25]</sup> With the advent of multianvil Walker-type large volume presses the thermal decomposition of P<sub>3</sub>N<sub>5</sub> could be prevented according to Le Chateliers' principle supported by the use of group 1 or 2 azides which upon heating decompose under the evolution of  $N_{2}$ .<sup>[26]</sup> This enabled systematic screening for nitridophosphates but the obstacle of missing single-crystals for X-ray diffraction remained. After Marchuk et al. employed NH<sub>4</sub>Cl as a mineralizer, the high-pressure high-temperature (HP/HT) approach yielded the desired quality and size of single crystals, reliably.<sup>[27]</sup> The ability to obtain accurate and precise structural data of nitridophosphates also coincided with a beginning interest in potential luminescent materials among nitridophosphates.<sup>[28-31]</sup> Despite tremendous efforts, expanding the structural and compositional variety of nitridophosphates beyond the group 1 and 2 elements to transition metal (TM) or rare-earth metals (RE) was not feasible. In analogy to nitridosilicate syntheses, the *TM*- or *RE* nitrides were believed to be too unreactive and the reduction of the metal or P and the oxidation of the nitride ion seemed to impede the discovery of new compounds. In addition, the investigated TMs tend to form stable phosphides like FeP. Kloß et al. developed the HP/HT solid-state metathesis that employs a metal halide and LiPN<sub>2</sub> resulting in a nitridophosphate and the lithium halide which is the driving force of the reaction. Following this approach, the compound class of RE nitridophosphates was made accessible and a few representatives of *TM* nitridophosphates were synthesized.<sup>[32-38]</sup> Surprisingly, the syntheses of compounds of group 4 elements resulted exclusively in oxonitridophosphates. Most of the presented syntheses aimed at negatively charged networks that needed metal cations for balancing charges. Vogel et al. delivered access to neutral main group nitrides BP<sub>3</sub>N<sub>6</sub> and SiP<sub>2</sub>N<sub>4</sub>NH. This was achieved by employing the binary nitrides BN or  $Si_3N_4$  respectively, (PNCl<sub>2</sub>)<sub>3</sub> and NH<sub>4</sub>N<sub>3</sub> which led to the development of quantitative amounts of HCl in contrast to Marchuk's mineralizer

approach.<sup>[39,40]</sup> SiP<sub>2</sub>N<sub>4</sub>NH constitutes the second nitride of Si and P after Baldus et al. discovered SiPN<sub>3</sub> both having syntheses from molecular precursors in common.<sup>[41]</sup> A unified synthetic approach that allows for the reaction of P<sub>3</sub>N<sub>5</sub> with refractory nitrides like Si<sub>3</sub>N<sub>4</sub> and *TM*N, the first objective of this thesis, was still missing. We employed NH<sub>4</sub>F as a mineralizer either exploiting the fluorophilicity of Si or the industrially known processes of nitride etching with HF. The syntheses were performed at 8 GPa of pressure and 1400–1700 °C hindering the decomposition of P<sub>3</sub>N<sub>5</sub> while delivering thermal energy for the reaction of either Si<sub>3</sub>N<sub>4</sub> or the *TM*N.

#### **1.3 Modern Methods of Structure Elucidation**

Explaining the properties of a new material is linked to the knowledge of its crystal structure. Challenges arise if routine structure elucidation by SCXRD (single crystal X-ray diffraction) fails due to microcrystallinity and/or real structure phenomena. Structure elucidation from PXRD (powder X-ray diffraction) has been the most common method in the case of microcrystalline samples as shown for  $\alpha$ -P<sub>3</sub>N<sub>5</sub>. However, the loss of information from SCXRD to PXRD leads to a loss in the precision of atomic site parameters. In addition, the presence of side phases, especially unknown ones, can severely complicate structure elucidation. Fahrnbauer et al. reported the collection of single-crystal X-ray diffraction data sets from microcrystallites, which have been precharacterized using a transmission electron microscope (TEM) and been subjected to subsequent diffraction experiments carried out employing microfocused synchrotron radiation.<sup>[40]</sup> This approach enabled the structure elucidation of e. g. Pb<sub>2</sub>Si<sub>5</sub>N<sub>8</sub>. The claim of an unusually short Pb-Pb distance of 3.190(1) Å, compared to Pb metal with 3.49 Å, was supported by accurate diffraction data.<sup>[41]</sup> Other examples of this approach include the crystal structures of U<sub>0.5x</sub>Sr<sub>2-0.75x</sub>Si<sub>5</sub>N<sub>8</sub> with  $x\approx1.05$ , Lu<sub>4</sub>Ba<sub>2</sub>[Si<sub>12</sub>O<sub>2</sub>N<sub>16</sub>C<sub>3</sub>]:Eu<sup>2+</sup>, BaP<sub>8</sub>N<sub>14</sub>, or La<sub>21</sub>P<sub>40</sub>O<sub>46</sub>N<sub>57</sub>.<sup>[17,31,38,44]</sup>

The imaging mode of high-resolution TEM aided in clarifying ambiguous results from SCXRD as shown for  $CaMg_2P_6O_3N_{10}$ .<sup>[29]</sup> SCXRD analysis led to a Ca position with a rather oblate displacement ellipsoid coordinated octahedrally by six N atoms. The interpretation of a randomly displaced Ca position in the octahedron was supported by HRTEM images compared to simulations based on SCXRD data. Aside from nitrides, this approach was also used for the chalcogenide Pb<sub>5</sub>Sb<sub>4</sub>S<sub>11</sub> to distinguish a disordered and an ordered polymorph

based on simulations.<sup>[45]</sup> This already shows the drawback of HRTEM as images are not readily interpretable and contrast maxima do not necessarily coincide with atomic positions. The advent of aberration-corrected scanning TEM (STEM) with a high-angle annular darkfield (HAADF) detector enabled an imaging mode with a direct correlation of contrast maxima and atomic positions with sub-Angstrom resolution. The dependence of image contrast from  $\sim Z^2$  enabled the distinction between ordered and disordered structure models for Ca<sub>1-x</sub>Li<sub>x</sub>Al<sub>1-</sub>  $_x$ Ge<sub>1+x</sub>N<sub>3</sub> ( $x \approx 0.2$ ) and in the case of Ce<sub>4</sub>Li<sub>3</sub>P<sub>18</sub>N<sub>35</sub>, the slight displacement of Ce atoms could be identified as the cause for superstructure reflections.<sup>[46,47]</sup> Considering that contrast scaling depends on  $\sim Z^2$  the limitations of this method become apparent if the atomic numbers of two elements become increasingly similar or are far apart. The interpretation of images can be complicated by insufficient dynamic range of the detector in both cases. A combination of atomic resolution STEM-HAADF and energy dispersive X-ray spectroscopy (EDX) yields a 3D dataset with image contrast in two dimensions and pixel-wise EDX spectra as the third. As shown for SrTiO<sub>3</sub> and predominantly perovskite-related compounds this method finally delivers images that are comparable to those from structure visualization software.<sup>[48,49]</sup> The emphasis on oxides in the literature can be most likely attributed to the potential beam damage caused by this method, especially as X-ray collection is less efficient for elements with Z < 30 thus either demanding longer acquisition times or increased beam currents and therefore aggravating the issue of beam damage even more.<sup>[50]</sup>

#### **1.4** Scope of this Thesis

Despite the tremendous efforts in synthesizing novel nitrides that have been achieved by high-pressure high-temperature chemistry, the reaction between multiple refractory nitrides had not been successful. In the two cases where double nitrides of Si and P were formed, SiPN<sub>3</sub> and SiP<sub>2</sub>N<sub>4</sub>NH, there was a need for molecular precursors, SiPN(NH)(NH<sub>2</sub>)<sub>4</sub> or (PNCl<sub>2</sub>)<sub>3</sub> respectively.<sup>[41,39]</sup> To overcome these limitations and employ the nitrides P<sub>3</sub>N<sub>5</sub> and Si<sub>3</sub>N<sub>4</sub>, the previously reported mineralizer-assisted approaches were developed further by employing NH<sub>4</sub>F as mineralizing agent.<sup>[27-29]</sup> This led to the compound classes of alkaline earth nitridosilicatephosphates and silicon imidonitridophosphates. The elemental combination Si and P is troublesome concerning X-ray diffraction as the atomic scattering factors of Si and P are similar. STEM-EDX with atomic resolution enables reliable statements on the distribution

of elements like Si and P in nitridic networks. The refractory interstitial transition metal nitrides were used as starting materials to evaluate the limits of the NH<sub>4</sub>F-assisted approach. This yielded (oxo)nitridophosphates with band gaps in the visible part of the electromagnetic spectrum.

In summary, this thesis addresses two needs of materials sciences. The first is developing a systematic access to novel compounds and materials: The NH<sub>4</sub>F-assisted HP/HT syntheses for mixed nitridic nitridosilicatephosphates and silicon imidonitridophosphates. Both classes of compounds may include promising phosphors for solid-state-lighting and are discussed in Chapters 3 and 4. Following this approach, the transition metal (oxo)nitridophosphates Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub>, Ti<sup>III</sup><sub>4</sub>Ti<sup>IV</sup>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub>, and Ti<sup>IV</sup>P<sub>4</sub>N<sub>8</sub> were synthesized from the highly stable binary nitrides ScN and TiN as presented in Chapters 5 and 6. Accordingly, refractory nitrides have finally been included in the staple of starting materials and semiconductor-like band gaps in (oxo)nitridophosphates have been obtained. The second demand addressed is the accurate structural investigation of occupational and positional disorder. Aberration-corrected STEM was performed on all compounds including Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> to elucidate real structure effects and order/disorder effects which will be discussed beginning from chapter 2.

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# 2 A Layered Tin Bismuth Selenide with Three Different Building Blocks that Account for an Extremely Large Lattice Parameter of 283 Å

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**Abstract:** The layered compound  $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27}$  exhibits an extraordinarily long-periodic 150R stacking sequence. The crystal structure contains three different building blocks, which form upon the addition of Sn to a Bi-rich bismuth selenide. Sn-doped Bi<sub>2</sub> double ("2") layers similar to those in elemental bismuth,  $Sn_{0.3}Bi_{1.7}Se_3$  quintuple ("5") layers and  $Sn_{0.4}Bi_{2.6}Se_4$ septuple ("7") layers are arranged in a 7525757525/75255/7525/75255/7525 sequence, which corresponds to a structure with a = 4.1819(4) and c = 282.64(6) Å in space group  $R\bar{3}m$ . The structure of a microcrystal was determined using microfocused synchrotron radiation and refined as a formally commensurately modulated structure in (3+1)D superspace (superspace group  $R\bar{3}m(00\gamma)00$ ), with a trivial basic structure that contains just one atom. The stacking sequence as well as the cation distribution are confirmed by aberration-corrected scanning transmission electron microscopy (STEM) in combination with chemical mapping by X-ray spectroscopy with atomic resolution. Stacking faults are not typical but have been observed occasionally.

## 2.1 Introduction with Results and Discussion

Owing to their pronounced compositional and structural diversity, chalcogenides represent one of the most intriguing classes of compounds in solid-state chemistry and physics.<sup>[1]</sup> Besides Bi<sub>2</sub>Se<sub>3</sub>, which is one of the most well-known topological insulators,<sup>[2]</sup> recent studies of for example, Bi<sub>2</sub>Te<sub>3</sub>,<sup>[3]</sup> MnBi<sub>2</sub>Te<sub>4</sub>,<sup>[4]</sup> and SnBi<sub>2</sub>Te<sub>4</sub>,<sup>[5]</sup> revealed their potential as topologically

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non-trivial materials. In addition, layered chalcogenides like BiSe,<sup>[6]</sup> SnSe,<sup>[7]</sup> and Bi-doped SnSe,<sup>[8]</sup> have attracted much attention in the field of thermoelectrics. Such chalcogenides also constitute a fascinating field of fundamental research. For example, single-crystal data of Ge<sub>4</sub>Se<sub>3</sub>Te revealed unexpected Ge–Ge bonds, which have been confirmed by STEM (scanning transmission electron microscopy) and theoretical calculations.<sup>[9]</sup> Modular stacking can lead to new inorganic-organic hybrid materials, for example, by inserting polymer chains in tin selenides.<sup>[10]</sup>

Layered chalcogenides often form homologous series of compounds, where variations of the slab thicknesses between van der Waals gaps characterize a plethora of new phases.<sup>[11]</sup> This can lead to compounds with large lattice parameters, which can be expected to be beneficial for thermoelectric properties as phonon scattering on the nanometer scale decreases thermal conductivity.<sup>[12]</sup> Compounds in the binary M/X systems with M =Sb, Bi and X = Se, Te form trigonal layered structures built up from  $M_2$  double layers (symbolized by "2")—comparable to those in the element structures of Sb and Bi—and tetradymite-like  $M_2X_3$  quintuple blocks symbolized by "5". Complex stacking results in large lattice parameters, for example, *c* = 103 Å for Bi<sub>8</sub>Te<sub>9</sub>, which features a 2555/2555/2555 sequence.<sup>[13]</sup> So far, no single-crystal X-ray data have been reported for such layered chalcogenides with long-periodic stacking sequences that exhibit lattice parameters larger than these 103 Å. Electron diffraction data and high-resolution transmission electron microscopy (HRTEM) of quenched samples, however, showed the presence of a rhombohedral bismuth telluride  $Bi_{2+\delta}Te_3$  with a 138*R*-Å) stacking sequence (δ 0.11 - 0.75, С ≈ 275 of type  $\approx$ 5552552552[5552552552]5552552552[5552552552.<sup>[14]</sup> In ternary systems such as *Tt*Te(Bi<sub>2</sub>Te<sub>3</sub>)<sub>n</sub> (*Tt* = Ge, Sn, Pb), the structural chemistry is extended by building blocks with septuple ("7") *Tt*Bi<sub>2</sub>Te<sub>4</sub> slabs in addition to quintuple ("5") Bi<sub>2</sub>Te<sub>3</sub> slabs.<sup>[15]</sup> The structure with the longest periodicity has been found in a 159R-type germanium bismuth telluride that exhibits a 557575757[557575757]557575757 stacking sequence with  $c \approx 318$  Å as deduced from electron diffraction patterns.<sup>[16]</sup> However, crystal structure refinements have not been reported for such extreme cases.

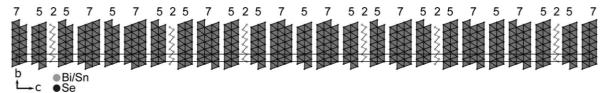
With respect to inorganic compounds, unit-cell dimensions of more than 100 Å are very unusual in general. They have been reported for a few binary compounds like the well-known polytypes of SiC,<sup>[17]</sup> ZnS,<sup>[18]</sup> CdI<sub>2</sub>,<sup>[19]</sup> and PbI<sub>2</sub>.<sup>[20]</sup> Varying sequences and orientations of the same structural entities result in polytypes with huge lattice parameters of up to 990 Å in

SiC.<sup>[21]</sup> Such crystallographic phenomena have usually been observed only in small domains by electron microscopy; and the structures were assigned by plausibility or trial and error methods in case diffraction data were available. Only very few actual structure refinements based on single-crystal data of structure models with very large lattice parameters do exist. In this respect, even lattice parameters of around 57 Å as recently observed for cesium rareearth silicates Cs<sub>3</sub>*RE*Si<sub>6</sub>O<sub>15</sub> (*RE* = Dy-Lu, Y, In) have been reported as being unusually large,<sup>[22]</sup> although for example, for the mineral turtmannite with *c* = 204 Å, a full structure refinement has been carried out.<sup>[23]</sup> Hexaferrites are also prone to form anisotropic structures with large translation periods along the stacking direction of Ba-rich and Fe-rich slabs, which can be arranged with varying sequences that exhibit translation periods of up to 1577 Å for Ba<sub>70</sub>Zn<sub>66</sub>Fe<sub>444</sub>O<sub>802</sub>.<sup>[24]</sup> While this extraordinary layer stacking has only been identified by electron microscopy, Rosseinsky et al. fully refined the crystal structure of Ba/Fe/Zn oxide hexaferrite polytypes, for example,  $Ba_{10}Fe_{72}Zn_8O_{126}$  with c = 488 Å, from single-crystal synchrotron data.<sup>[25]</sup> Hexaferrites can be identified by the 00*l* reflections in electron diffraction patterns and described by a unified (3+1)D superspace model.<sup>[26]</sup> Similar descriptions have been reported for perovskite-like compounds.<sup>[27]</sup> Lidin et al. developed a similar superspace formalism for the system Bi/Se described above.<sup>[28]</sup> Such superspace descriptions are well known from modulated structures, but the (3+1)D description of longperiodic layered structures does not mean that a somehow simpler structure is modified by wave-like displacements or occupation modulations. in fact, the hypothetical basic structure would consist of just one or a few atoms. Therefore, the superspace formalism focuses, on the one hand, on a unified description of a series of compounds. On the other hand, it gives access to elegant structure refinements in cases where the reflections of extremely large unit cells are too closely spaced to be integrated from area-detector data based on a 3D periodic indexing.

In the pseudobinary system  $(SnSe)_xBi_2Se_3$ , compounds with very diverse structures are formed.<sup>[29]</sup> Besides cubic  $Sn_4Bi_2Se_7$  (x = 4) with defect NaCl-type structure, and layered  $SnBi_4Se_7$  (x = 0.5) with defect  $GeSb_2Te_4$ -type structure, at least four compounds ( $0.8 \le x \le 3$ ) with structures derived from lillianite ( $Pb_3Bi_2S_6$ )<sup>[30]</sup> have been discovered. Here we report the single-crystal structure of a new compound in the system Sn/Bi/Se, which features an extraordinarily long lattice parameter of c = 282.6 Å in a 150R stacking sequence. A complex layered phase with the chemical composition  $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27}$  was formed during the decomposition of samples of  $SnBi_4Se_7$  with defect  $GeSb_2Te_4$ -type structure.<sup>[29]</sup> The fact that this decomposition occurred during repeated slow heating indicates that the new phase is most likely thermodynamically stable at temperatures below  $\approx 500$  °C. The heterogeneous product was examined by scanning electron microscopy (SEM) and energy-dispersive X-ray spectroscopy (EDX), which revealed regions of different contrast and composition (Figure A1 and Table A1, "A" denotes Figures and Tables in the Supporting Information). In addition to a minority phase, which according to EDX corresponds to  $Sn_{11.49}Bi_{12.39}Se_{30}$ ,<sup>[29]</sup> and a main phase of  $Sn_{0.85}Bi_{2.15}Se_4$ ,<sup>[29]</sup> the title compound has been identified as a main component that is slightly Sn-depleted compared to the nominal composition  $SnBi_4Se_7$ .

A single crystal (Figure A2) of this unknown phase has been characterized by transmission electron microscopy (TEM) and EDX (Table A2). Using a fluorescence detector, this crystallite has been centered in a microfocused synchrotron beam at the European Synchrotron Radiation Facility (ESRF) and single-crystal diffraction data were recorded.<sup>[31]</sup> The crystallite exhibits a complex diffraction pattern. Initial attempts to determine the crystal structure from the single-crystal data by "routine data processing" resulted in a structure close to that of BiSe.<sup>[32]</sup> However, the fact that 60 % of all reflections remained unindexed indicated that this was a wrong model. Taking into account the precise positions of the reflections leads to a unit cell with c = 282.6(6) Å; however, reasonable data cannot be extracted as many very weak reflections lie very close to strong ones that impede intensity integration. This problem can be overcome by formally describing the diffraction pattern based on the (3+1) dimensional superspace group  $R^{3}m(00\gamma)00$ . This involves a formal basic cell of a = 4.1819(4) and c =5.6528(5) Å, which contains just one atom, with a commensurate modulation vector q =87/50 = 1.74. This way, all reflections in the reciprocal lattice can be indexed (Figure A3–A5). Following the concept outlined by Lidin et al.,<sup>[28]</sup> the structure solution and refinement were based on modulation parameters of the single atom site at the origin of the small unit cell of the basic structure (for details on the refinement see the Supporting Information). Initially assuming one Bi and one Se atom on this site, both fully occupied, the additional dimension in (3+1)D superspace can be used to modulate the occupancy using discontinuous functions. Crenel functions were used to describe the occupancy modulation by defining intervals in the additional dimension x4 where the corresponding atom exists. The refinement of the crenel function width  $\Delta$  defines the overall chemical composition.

The refined value of  $\Delta$ [Se1] = 1- $\Delta$ [Bi1] = 0.5398(2) was constrained to 0.54 in order to provide an integer number of atoms in the unit cell. The resulting sum formula Bi<sub>0.46</sub>Se<sub>0.54</sub> and a *q*-vector of q = 1.74 = 87/50 correspond to a commensurate structure with a stacking sequence of 7525757525|7525757525|7525757525 (Figure 2.1). A refinement of positional modulation waves for Se and Bi shows their displacements along [001] (×3 direction), which became obvious from *F*<sub>obs</sub> Fourier maps (Figure A6). Mixed occupancy of Sn and Bi was taken into account by a further modulation of site occupancy factors, again constrained to full occupancy. The refinement reveals a varying Sn fraction on the cation sites. The occupancy of the atoms as a function of the t variable in superspace (Figure A7) shows where the atoms are present in the ×4 dimension. The Sn fraction amounts to 4% in the double layers, whose structure corresponds to that of elemental Bi, up to 16% at the central position in septuple layer (Figure A8). This results in a sum formula of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub>, well in line with the EDX analyses (Table A1 and A2). The range of interatomic distances is comparable to those found in other bismuth selenides and tin bismuth selenides (Table A3, visualized as a function of tin Figure A9). Crystallographic data<sup>[33]</sup> are shown in Table A4, structure parameters including modulation parameters are given in Tables A5 and A6 (note that R values for some "satellite orders" are affected by adjacent very strong reflections of other orders and that the intensity does not decrease with increasing "satellite order" as there is no wave-like modulation).

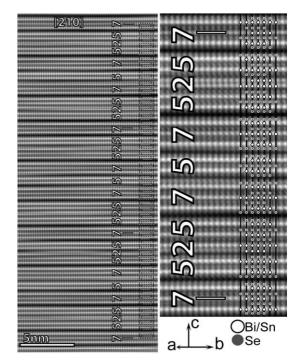


**Figure 2.1.** Crystal structure of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> viewed along [100]: 150*R* stacking sequence of 7525757525 [7525757525]7525757525 (cf. text); Se atoms are displayed in black, Bi/Sn atoms in gray; (Bi,Sn)Se<sub>6</sub> octahedra are highlighted in order illustrate the different thicknesses of the rock-salt-like slabs; the zigzaglines correspond to (Bi,Sn)<sub>2</sub> layers; the unit cell is indicated.

The long-periodic trigonal structure of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> with a 150*R* stacking sequence is built up from three different kinds of slabs: double and quintuple layers known from the system Bi/Te, an additional septuple layers obtained by the introduction of Sn. The 7525757525|7525757525|7525757525 sequence corresponds to an extremely large lattice parameter of the 3D-periodic structure, which amounts to 282.6 Å. This is the first layered chalcogenide with three different building blocks, and it features by far the longest translation

period found in the system Sn/Bi/Se. Taking into account the slabs, the formula can be written as  $(Sn_{0.1}Bi_{1.9})_2(Sn_{0.3}Bi_{1.7}Se_3)_5(Sn_{0.4}Bi_{2.6}Se_4)_3$ . Assuming Bi<sup>0</sup> and Sn<sup>0</sup> in the double layers and Sn<sup>2+</sup>, Bi<sup>3+</sup> and Se<sup>2-</sup> in the quintuple and septuple slabs, the overall charge of the compound is neutral within standard deviations. Sn/Bi cation disorder in septuple slabs has also nicely been confirmed by X-ray diffraction and Mössbauer spectroscopy of SnBi<sub>4</sub>Se<sub>7</sub>.<sup>[34]</sup>

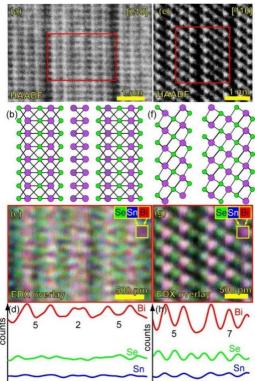
A structure model in 3D derived from the description in superspace (Figure 2.1) can be only be refined tentatively (Table A7) using a 3D dataset. Since such a 3D dataset corresponds to a (3+1)D one with satellite reflections up to the 30<sup>th</sup> order, the model has too many parameters without a reasonable amount of observed data. Thus, the commensurately modulated description is inevitable although there is no wave-like modulation as known from typical modulated structures. However, the refined atom positions of such a 3D model do not deviate much from those directly derived from the (3+1)D superspace model (Table A8). Referring to the superspace description of bismuth selenides by Lidin et al.,<sup>[28]</sup> the *q*-vector of 87/50 = 1.74 can be correlated to the stacking sequence and the resulting number of atom layers in the unit cell (Table A9). STEM-HAADF images (Figure 2.2) confirm the long-periodic stacking sequence of  $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27}$  as derived from the single crystal data. The presence of Sn on all metal atom sites as indicated by the X-ray data corresponds well with aberrationcorrected STEM-HAADF and chemical mapping by atomic-resolution EDX.



**Figure 2.2.** Fourier-filtered STEM-HAADF image (300 kV) along [210] with an image of the crystal structure of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> in the corresponding projection as an overlay; the asymmetric unit (7525757525 sequence, cf. text) is shown on the right in an enlarged image. In the HAADF images, brighter contrast

corresponds to Bi/Sn atom columns and darker contrast to Se atomcolumns. Se atoms are displayed as gray circles and Bi/Sn as white circles.

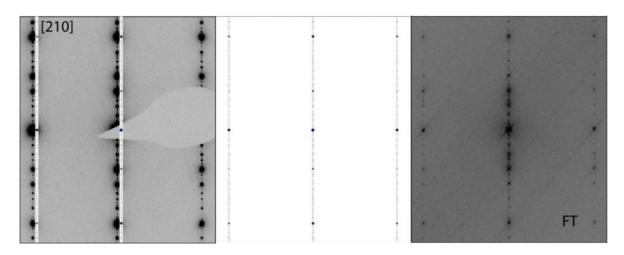
Figure 2.3 shows a 525 stacking sequence with a double layer in the center (a–d) and a 57 stacking sequence with a van der Waals gap between the two blocks (e–h). Intensities in STEM images and their projection perpendicular [001] further confirm the atom distribution along the sequence (Figure A10). Selected-area electron diffraction (SAED) patterns and Fourier transforms of STEM images along [210] match the diffraction pattern calculated from the 3D model derived from the (3+1)D structure refinement (Figure 2.4) and thus additionally confirm the lattice parameter of 286.6(6) Å. A further SAED pattern along [100] with the corresponding simulation is shown in Figure A11.



**Figure 2.3.** STEM-HAADF images along zone axes [210] (a) and [110] (e), brighter contrast corresponds to Bi/Sn atom columns and darker contrast toSe atom columns; (c, g) EDX maps of the indicated (red box) section of the STEM-HAADF images with (b, f) corresponding structure projections and (d, h) integrated EDX signal projected perpendicular to the layers (projection along [210] or [110], that is, parallel to the layers, Se signals and atoms are highlighted in green and Bi/Sn signals and atoms in red and blue), showing mixed occupancy on all cation positions as indicated by the violet color. (b), (c) and (d) show a 525 stacking sequence with a double layer in the center; (f), (g) and (h) depict a 57 stacking sequence with a van der Waals gap between the two blocks. (a, c) were recorded using 200 kV, (e, g) using 300 kV accelerating voltage.

Occasionally, stacking faults and different stacking sequences have locally been observed during extensive electron microscopy studies (Figure A12). Yet, the presented structure is by far the predominant stacking sequence.

In conclusion, repeated annealing afforded Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub>, a new compound in the system Sn/Bi/Se with an extraordinarily long-periodic 150*R* stacking sequence 7525757525|7525757525|7525757525, resulting in a lattice parameter of c = 282.6(6) Å. The structure is built up from three different types of slabs, which significantly expands the plethora of known structure models for layered chalcogenides. In Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub>, the structural features of binary pnictogen-rich chalcogenides, that is, double and quintuple layers and those of ternary tetradymite-like chalcogenides, that is, quintuple and septuple layers, are combined and lead to the unusual long-range ordering. It remains, however, an open and very intriguing question, why such structures form and why they are seemingly thermodynamically stable. The large unit cell and the mixed occupancies impede theoretical calculations as alternative models with different arrangements of the same building blocks will show only tiny differences in the strength of van der Waals interactions between the slabs. For example, in the related compound Ge<sub>2</sub>Sb<sub>2</sub>Te<sub>5</sub>, which features much smaller translation periods, activation energies for transitions between ordered and disordered models, which differ by ca. 1 eV in energy, have been calculated to be as low as 0.005 eV.<sup>[35]</sup> The structure model was obtained by means of X-ray diffraction using a microfocused synchrotron beam and confirmed by STEM-HAADF measurements with atomic resolution and EDX mappings. The synergism of these methods enables unprecedented accuracy of structure determinations of microcrystalline compounds and promotes the discovery of numerous new related layered compounds with mixed site occupancies. These may exhibit intriguing physical properties such as thermoelectricity or non-trivial topological behavior.



**Figure 2.4.** Comparison of an experimental SAED pattern (top, zone axis [210]) (left), a Fourier transform of the STEM-HAADF image displayed in Figure 2.2 (right) with a simulated SAED pattern (middle, based on the structure model from single crystal data); for better comparability, cutouts of the simulation are repeated as insets in the experimental pattern

#### **Experimental Section**

 $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27}$  was formed during the decomposition of a quenched sample of  $SnBi_4Te_7^{[29]}$ upon repeated heating. For TEM measurements, the powdered sample was drop-cast on a copper grid coated with a holey carbon film. This grid was fixed on a glass capillary for single crystal data collection at the European Synchrotron Radiation Facility (ESRF). For STEM measurements, polycrystalline pieces were embedded in epoxy resin and mechanically cut and thinned to a thickness of  $\approx 20 \ \mu m$  in the center using a dimple grinder. Electron transparency was achieved by polishing a hole in the center of the disc by Ar-ion milling (Figure A13). Further details are given in the Supporting Information.

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Information. Deposition number 1949450 contains the supplementary crystallographic data for this paper. These data are provided free of charge by the joint Cambridge Crystallographic Data Centre and Fachinformationszentrum Karlsruhe Access Structures service.

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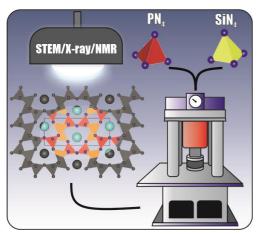
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# High-Pressure High-Temperature Synthesis of Mixed Nitridosilicatephosphates and Luminescence of *AE*SiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> (*AE* = Sr, Ba)

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**Abstract:** Tetrahedra-based nitrides with network structures have emerged as versatile materials with a broad spectrum of properties and applications. Both nitridosilicates and nitridophosphates are well-known examples of such nitrides that upon doping with Eu<sup>2+</sup> exhibit intriguing luminescence properties, which makes them attractive for applications. Nitridosilicates and nitridophosphates show manifold



structural variability; however, no mixed nitridosilicatephosphates except SiPN<sub>3</sub> and SiP<sub>2</sub>N<sub>4</sub>NH have been described so far. The compounds  $AESiP_3N_7$  (AE = Sr, Ba) were synthesized by a highpressure high-temperature approach using the multianvil technique (8 GPa, 1400—1700 °C) starting from the respective alkaline earth azides and the binary nitrides P<sub>3</sub>N<sub>5</sub> and Si<sub>3</sub>N<sub>4</sub>. The latter were activated by NH<sub>4</sub>F, probably acting as a mineralizing agent. SrSiP<sub>3</sub>N<sub>7</sub> and BaSiP<sub>3</sub>N<sub>7</sub> were obtained as single crystals. They crystallized in the barylite-10 (M = Sr) and barylite-20 structure types (M = Ba), respectively, with P and Si being occupationally disordered. Cation disorder was further supported by solid-state NMR spectroscopy and energy-dispersive X-ray spectroscopy (EDX) mapping of BaSiP<sub>3</sub>N<sub>7</sub> with atomic resolution. Upon doping with Eu<sup>2+</sup>, both compounds showed blue emission under UV excitation.

## 3.1 Introduction with Results and Discussion

Emerging environmental consciousness has pushed the development of solid-state lighting solutions forward. The invention of efficient InGaN-based blue LEDs (light-emitting diodes) enabled the development of pc-LEDs (phosphor-converted) with remarkable properties in terms of color temperature, color rendition, and efficacy. Significant improvements in the aforementioned properties were possible due to nitride compounds such as  $M_2Si_5N_8$ :Eu<sup>2+</sup> (M = Sr, Ba),  $MSi_2O_2N_2$ :Eu<sup>2+</sup> (M = Ca, Sr, Ba),  $SrLiAl_3N_4$ :Eu<sup>2+</sup>, and  $MAlSiN_3$ :Eu<sup>2+</sup>.<sup>[1-4]</sup>

Materials properties concerning solid-state lighting can be tuned by dopant concentration to a limited extent, affecting Stokes shifts in emission spectra, or by a variation of the size of coordination polyhedra by substitution such as introducing Sr on Ba sites. Completely shifted emission properties, however, can only be achieved by a fundamental alteration of the host lattice.<sup>[5, 6]</sup>

The main goal of this work was to expand the compositional and structural diversity of tetrahedra-based luminescent materials. Thus, discovery of the title compounds  $SrSiP_3N_7$  and  $BaSiP_3N_7$  opens up the novel compound class of mixed nitridosilicatephosphates, which can now be further explored as innovative host lattices. While nitridosilicates have been investigated thoroughly and nitridophosphates show similar promising structures and properties, only two compounds that contain both  $SiN_x$  (x = 4, 6) and  $PN_4$  units have been reported so far, that is,  $SiPN_3$  and  $SiP_2N_4NH.^{[7-9]}$  The crystal structure of  $SiPN_3$  corresponds to a defect wurtzite-type arrangement with mixed occupation of Si and P at the tetrahedral sites. The crystal structure of  $SiP_2N_4NH$  is related to sillimanite-type  $Al_2SiO_5$ . It is built up from edge-sharing  $SiN_6$  octahedra interconnected by all-side vertex-sharing  $PN_4$  tetrahedra. A possible explanation for the challenges concerning syntheses that are involved in the preparation of mixed nitridosilicatephosphates could be the chemical inertness of  $Si_3N_4$ , while  $P_3N_5$  already decomposes at temperatures above 850 °C if no external pressure is applied. As shown in previous publications, according to Le Chatelier's principle, the decomposition of  $P_3N_5$  under the formation of  $N_2$  is suppressed by external pressure.<sup>[10]</sup>

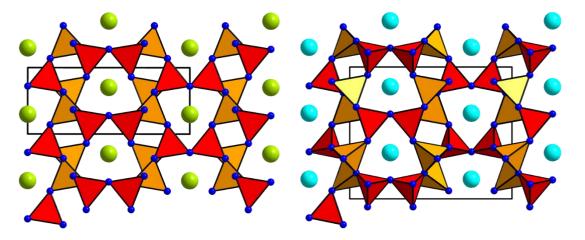
NH<sub>4</sub>Cl has been successfully employed as a mineralizer facilitating crystal growth of nitridophosphates. HCl formed in situ most likely leads to reversible P–N bond formation and cleavage. <sup>[11, 12]</sup> After nitridosilicatephosphates proved to be not accessible with the help of NH<sub>4</sub>Cl, changing the mineralizing agent to NH<sub>4</sub>F afforded the title compounds *AE*SiP<sub>3</sub>N<sub>7</sub> (*AE* = Sr, Ba). This may be explained by the fact that HF cannot only reversibly cleave and form P–N bonds, but also Si–N bonds. The surface of Si<sub>3</sub>N<sub>4</sub> features SiNH<sub>2</sub> groups that can be attacked by F-in a nucleophilic substitution.<sup>[13]</sup>

The nitridosilicatephosphates  $AESiP_3N_7$  (AE = Sr, Ba) were synthesized by high-pressure hightemperature (HP/HT) reactions at 8 GPa and 1400 °C (Ba) and 1700 °C (Sr), respectively, using a modified Walker-type multianvil apparatus.<sup>[14]</sup> The synthesis of SrSiP<sub>3</sub>N<sub>7</sub> at temperatures below 1700 °C resulted in samples with significant amounts of unknown side phases. Reactions followed the so-called azide route using P<sub>3</sub>N<sub>5</sub>, Si<sub>3</sub>N<sub>4</sub>, and the respective metal azide as starting materials with additional NH<sub>4</sub>F ( $\approx$ 5 wt %) as a mineralizing agent [Eq. (3.1)]. To investigate luminescence properties, samples with the addition of approximately 1 mol % of EuF<sub>3</sub> (concerning *AE*<sup>2+</sup>) to the starting mixture were prepared.

$$3 AE(N_3)_2 + 3 P_3N_5 + Si_3N_4 \rightarrow 3 AESiP_3N_7 + 8 N_2$$
(3.1)

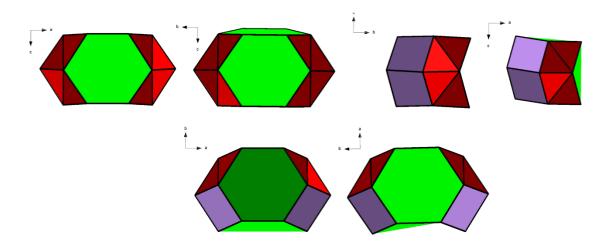
The title compounds were obtained as colorless powders (Eu<sup>2+</sup>-doped samples of SrSiP<sub>3</sub>N<sub>7</sub> show a yellow tint) and showed no sensitivity to air or moisture. More detailed information on the HP/HT synthesis is given in the Supporting Information.

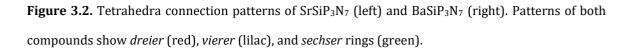
The crystal structures were elucidated by single-crystal X-ray diffraction (SCXRD) using direct methods and least-squares refinement.  $SrSiP_3N_7$  crystallizes in space group  $Pmn2_1$  (no. 31) with Z = 2. BaSiP<sub>3</sub>N<sub>7</sub> crystallizes in space group Pnma (no. 62) with Z = 4; details are given in Tables 1 and B2–7. In addition, Rietveld refinements indicate the presence of BaSiP<sub>3</sub>N<sub>7</sub> crystallizing in space group  $Pmn2_1$  (no. 31) with Z = 2 as a side phase. Both compounds are isotypic to the two polymorphs of barylite BaBe<sub>2</sub>Si<sub>2</sub>O<sub>7</sub>. SrSiP<sub>3</sub>N<sub>7</sub> corresponds to the barylite-1*O* polymorph, whereas BaSiP<sub>3</sub>N<sub>7</sub> features the structure of barylite-2*O*. The structures of barylite-1*O* and barylite-2*O* represent the maximum degree of order (MDO) polytypes of their polytype family. Both structures consist of a network of all-vertex-sharing PN<sub>4</sub> and (Si<sub>0.5</sub>P<sub>0.5</sub>)N<sub>4</sub> tetrahedra and elongated square pyramid (J<sub>8</sub>)  $AEN_9$  (AE = Sr, Ba) polyhedra (Figure B1).<sup>[15]</sup>



**Figure 3.1.** Crystal structures of SrSiP<sub>3</sub>N<sub>7</sub> (left) and BaSiP<sub>3</sub>N<sub>7</sub> (right) both along [001]. For SrSiP<sub>3</sub>N<sub>7</sub> PN<sub>4</sub> tetrahedra (red) and (Si<sub>0.5</sub>P<sub>0.5</sub>)N<sub>4</sub> (orange) all vertices point in the same direction (behind the plane of projection). For BaSiP<sub>3</sub>N<sub>7</sub> the orientation of tetrahedra vertices alternate.

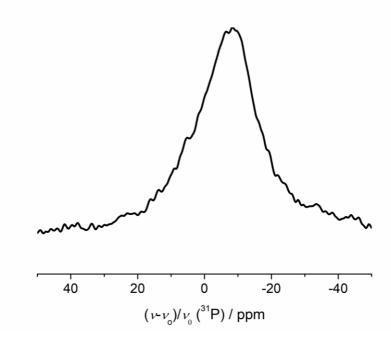
The main difference between the two polymorphs concerns the arrangement of tetrahedra. While in  $SrSiP_3N_7$  all tetrahedra vertices point in the same direction, those in  $BaSiP_3N_7$ alternate, which results in a doubled unit cell with 2b (SrSiP<sub>3</sub>N<sub>7</sub>) = a (BaSiP<sub>3</sub>N<sub>7</sub>) (Figure 3.1). Although tetrahedra orientation differs in both compounds, the environment of AE atoms is strikingly similar. The tetrahedra connection patterns show that both compounds consist of dreier, vierer, and sechser rings that, apart from slight distortions, are arranged and distributed in the same manner (Figure 3.2), leading to the same topology point symbol {3<sup>2</sup>.4<sup>3</sup>.5.6<sup>4</sup>}{3<sup>4</sup>.4<sup>5</sup>.5<sup>4</sup>.6<sup>2</sup>}.<sup>[16]</sup> Both compounds exhibit one tetrahedrally coordinated site shared by Si and P, while the other site is solely occupied by P. CSD 2050660 and 2050661 contain supplementary crystallographic data for this paper. These data can be obtained free of charge from FIZ Karlsruhe via www.ccdc.cam.ac.uk/structures. In the case of SrSiP<sub>3</sub>N<sub>7</sub>, potential ordering of Si and P was considered by symmetry reduction and refinement of the structure against SCXRD data in the subgroups P2<sub>1</sub>, Pn, and Pm of space group Pmn2<sub>1</sub>. However, no indications of complete ordering were found. In all structure models, the volumes of the resulting four symmetrically independent tetrahedra were compared as (P,Si)-N bond lengths were not sufficiently meaningful for discrimination.<sup>[17]</sup> This investigation led to two different kinds of tetrahedra.





Structure models in subgroups *P*2<sub>1</sub>, *Pn*, and *Pm* reveal two tetrahedra exhibiting a volume of 2.21–2.24 Å<sup>3</sup> that coincides with the volume of PN<sub>4</sub> tetrahedra from known nitridophosphates in literature and corresponds to one site in the final structure model in *Pmn*2<sub>1</sub>. The other two

tetrahedra have a volume of 2.37–2.40 Å<sup>3</sup>, which lies between the volumes of PN<sub>4</sub> (2.13–2.28 Å<sup>3</sup>) and SiN<sub>4</sub> tetrahedra (2.52–2.76 Å<sup>3</sup>) (Tables B10 and 11).<sup>[8, 9, 26–28, 18–25]</sup> BVS (bond valence sum) calculations performed on all structure models revealed two tetrahedral sites fully occupied by P and two tetrahedral sites occupied by Si and P in a 1:1 ratio (Tables B13–15).<sup>[29]</sup> In the case of BaSiP<sub>3</sub>N<sub>7</sub>, the ordering of tetrahedra was considered by symmetry reduction and refinement of the structure against SCXRD data in subgroups *Pna2*<sub>1</sub>, *P2*<sub>1</sub>*ma*, and *P2*<sub>1</sub>/*c* of space group *Pnma*. Only subgroups retaining the extinction condition of the *a* glide planes present in *Pnma* were taken into account because electron diffraction parallel to [001] showed no violation of the extinction conditions. Additional electron diffraction patterns parallel to [100] showed no violation of the extinction conditions of the *n* glide, too, further supporting the structure model in space group *Pnma* (a comparison of experimental diffraction patterns with simulated ones based on the structure model in space group *Pnma* is given in Figure B5).



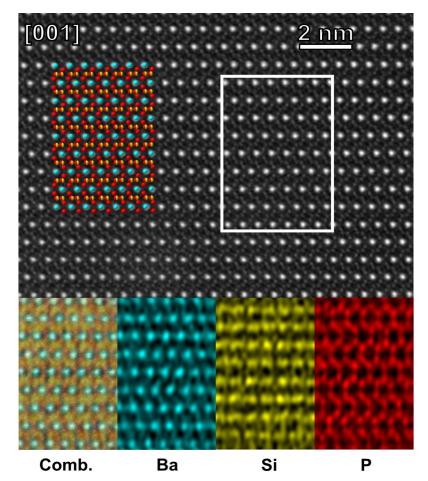
**Figure 3.3.** <sup>31</sup>P NMR spectrum showing one broad signal for BaSiP<sub>3</sub>N<sub>7</sub> instead of two signals for the different crystallographic sites most likely due to the occupational disorder

formula	SrSiP <sub>3</sub> N <sub>7</sub>	BaSiP <sub>3</sub> N <sub>7</sub>
molar mass /g·mol⁻¹	306.69	356.41
crystal system	orthorhombic	orthorhombic
space group	<i>Pmn</i> 2 <sub>1</sub> (no. 31)	<i>Pnma</i> (no. 62)
lattice parameters / Å	a = 11.979(2)	a = 9.9048(3)
	b = 4.9040(10)	b = 12.1858(3)
	c = 4.6870(9)	c = 4.73580(10)
cell volume / ų	275.34(10)	571.60(3)
formula units/ unit cell	2	4
density / g·cm <sup>3</sup>	3.699	4.142
μ / mm <sup>-1</sup>	10.807	7.927
temperature / K	296(2)	297(2)
absorption correction	semiempirical	
radiation	Mo- $K_{\alpha}$ ( $\lambda = 0.71073$ Å)	
F(000)	292	656
heta range / °	$3.4 \le \theta \le 44.09$	3.34≤ θ≤38.44
total no. of reflections	9890	10523
Independent reflections [I≥2σ(I) / all]	1253/ 1379	959/ 1128
$R_{\sigma}, R_{ m int}$	0.0393, 0.0952	0.0234, 0.0481
refined parameters	60	58
Goodness of fit	1.100	1.044
<i>R</i> -values [I≥2σ(I)]	$R_1 = 0.0293$	$R_1 = 0.0261$
	$wR_2 = 0.0731$	$wR_2 = 0.0576$
<i>R</i> -values (all data)	$R_1 = 0.0343$	$R_1 = 0.0340$
	$wR_2 = 0.0750$	$wR_2 = 0.0599$
$\Delta  ho_{ m max}$ , $\Delta  ho_{ m min}$ / e·Å <sup>-3</sup>	2.18, -1.49	0.81, -1.53

**Table 3.1:** Crystallographic data of the single-crystal structure refinements of AESiP<sub>3</sub>N<sub>7</sub> (AE = Sr, Ba).Standard deviations are given in parentheses.

The comparison of resulting tetrahedral volumes showed the same features as for SrSiP<sub>3</sub>N<sub>7</sub> (Table B12). BVS calculations performed for the different structure models again suggested two sites completely occupied by P and two sites shared by Si and P (Tables B16–18). The simple approach of comparing tetrahedra volume as a tool for assigning Si and P, which lack scattering contrast, to the respective sites was indeed confirmed by scanning transmission electron microscope energy-dispersive X-ray spectroscopy (STEM-EDX) mapping with atomic resolution

for BaSiP<sub>3</sub>N<sub>7</sub>. These data support the model in space group *Pnma*, showing two sites with mixed Si/P occupation (Figure 3.4, enlarged version see Figure B6). This result is also corroborated by <sup>31</sup>P solid-state magic angle spinning (MAS)-NMR spectra, which show a broad signal [full width at half-maximum (fwhm) = 19.7 ppm] that is consistent with a disordered model (Figure 3.3, Figure B4).

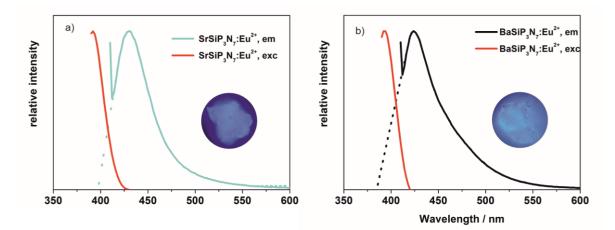


**Figure 3.4.** Atomic resolution STEM EDX of BaSiP<sub>3</sub>N<sub>7</sub> along [001]. STEM HAADF image (top) with structure overlay (Ba cyan, P red, Si yellow). The inset shows the corresponding area for EDX maps (bottom) with a combined color map, Ba map (cyan), Si map (yellow), and P map (red).

Line broadening in the NMR spectrum is probably due to disorder in the second coordination sphere of P atoms. In contrast, ordered nitrides like  $BP_3N_6$  or  $Li_{12}P_3N_9$  show very narrow signals in their <sup>31</sup>P NMR spectra.<sup>[8, 30]</sup>

In both structures, the connectivity of the tetrahedra via their vertices can explain the presence of different tetrahedral volumes. The smaller tetrahedra have three vertices occupied by twofold bridging nitrogen atoms  $N^{[2]}$  and one vertex occupied by a threefold bridging nitrogen atom  $N^{[3]}$ . The larger tetrahedra, in contrast, feature two vertices occupied

by N<sup>[2]</sup> and two vertices occupied by N<sup>[3]</sup>. Chemical analysis by EDX agrees with the sum formulas. Due to ambiguous O contents (as indicated by EDX measurements) either surface hydrolysis or slight compositional variations cannot completely be ruled out so that a phase width according to  $AESi_{1+x}P_{3-x}N_{7-x}O_x$  (AE = Sr, Ba) with x < 1 could also be considered (Table B8) even though some analyses show no O.



**Figure 3.5.** Emission spectra of a) SrSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> (blue) and b) BaSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> (black), measured data in solid lines and extrapolation in dotted lines, respective excitation spectra (red) (insets: micrographs of luminescent particles).

Upon doping with Eu<sup>2+</sup>, both compounds emit blue light under UV excitation. Luminescence spectra show emission maxima of  $\lambda_{max} = 430$  nm for SrSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> and  $\lambda_{max} = 424$  nm for BaSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> upon excitation with  $\lambda_{exc} = 400$  nm. The emission curves were extrapolated to give an estimate of the fwhm, which amount to 45 nm (2404 cm<sup>-1</sup>) for SrSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> and 53 nm (2731 cm<sup>-1</sup>) for BaSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> (Figure 3.5). The corresponding Stokes shifts are 38 nm (2254 cm<sup>-1</sup>) for SrSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> and 32 nm (1925 cm<sup>-1</sup>) for BaSiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup>. The presence of a single narrow emission band for both phosphors can be explained by the emission properties of Eu<sup>2+</sup>and the presence of a single crystallographic site for the alkaline earth ions suitable for doping with Eu<sup>2+</sup> with *AE*–N distances ranging from 2.696(3)–3.270(3) Å (SrSiP<sub>3</sub>N<sub>7</sub>). The similarity of emission properties in terms of fwhm are most likely to be explained by the P–Si<sub>0.5</sub>P<sub>0.5</sub> "cages" around the *AE* position, which are very similar. Thus, only minute deviations are caused by different *AE* cation sizes even though both structures differ with respect to their space groups, unit cell volumes, and tetrahedra orientations.

### Conclusions

High-pressure high-temperature synthesis with the addition of NH<sub>4</sub>F is a suitable approach to the synthesis of mixed nitridosilicatephosphates. The compounds *AE*SiP<sub>3</sub>N<sub>7</sub> (*AE* = Sr, Ba) adopt the structure types of the two polymorphs of the mineral barylite. This structure type has not been observed for nitride compounds so far. Although silicon and phosphorus exhibit little contrast in X-ray diffraction, the comparison of polyhedra volumes led to structure models with an occupationally disordered site that also persists if potential ordering is considered by symmetry reduction. The disordered model for BaSiP<sub>3</sub>N<sub>7</sub> is further supported by solid-state NMR spectroscopy. Scanning transmission electron microscopy energydispersive X-ray spectroscopy (STEM-EDX) mapping with atomic resolution enables to directly observe said disorder, which is additionally in accordance with systematic absences observed in electron diffraction patterns. Therefore, nitridosilicatephosphates have the potential to significantly diversify the structural chemistry of nitrides. Their suitability as host lattices for rare-earth activator ions seems especially intriguing considering the emission properties of other compounds with multiple tetrahedra centers like CaAlSiN<sub>3</sub>:Eu<sup>2+</sup>, Sr[Li<sub>2</sub>Al<sub>2</sub>O<sub>2</sub>N<sub>2</sub>]:Eu<sup>2+</sup>, or Sr[LiAl<sub>3</sub>N<sub>4</sub>]:Eu<sup>2+</sup>,<sup>[2, 31, 32]</sup>

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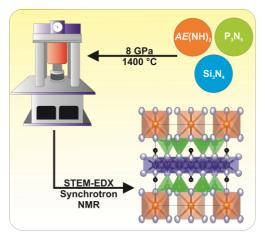
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4

# Nitridic Analogs of Micas *AE*Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (*AE* = Mg, Mg<sub>0.94</sub>Ca<sub>0.06</sub>, Ca, Sr)

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**Abstract:** We present the first nitridic analogs of micas, namely  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg,  $Mg_{0.94}Ca_{0.06}$ , Ca, Sr), which were synthesized under high-pressure high-temperature conditions at 1400 °C and 8 GPa from the refractory nitrides  $P_3N_5$ and  $Si_3N_4$ , the respective alkaline earth amides, implementing  $NH_4F$  as a mineralizer. The crystal structure was elucidated by single-crystal

diffraction with microfocused synchrotron radiation, energy-dispersive X-ray spectroscopic (EDX) mapping with atomic resolution, powder X-ray diffraction, and solid-state NMR. The structures consist of typical tetrahedra–octahedra–tetrahedra (T-O-T) layers with P occupying T and Si occupying O layers, realizing the rare motif of sixfold coordinated silicon atoms in nitrides. The presence of H, as an imide group forming the SiN<sub>4</sub>(NH)<sub>2</sub> octahedra, is confirmed by SCXRD, MAS-NMR , and IR spectroscopy. Eu<sup>2+</sup>-doped samples show tunable narrow-band emission from deep blue to cyan (451–492 nm).

## 4.1 Introduction with Results and Discussion

Silicates offer a broad range of structural diversity ranging from discrete SiO<sub>4</sub> tetrahedra to ribbons, sheets and frameworks. These structural motifs are unmatched in their diversity as compared to other tetrahedral anions like phosphates, sulfates or vanadates. Such structural diversity accompanied by chemical stability and their mechanical properties allows silicates to be employed in several applications, e.g. dielectrics, construction materials, fire retardants,

and as host compounds for activator ions in luminescent materials.<sup>[1–3]</sup> The plethora of structural motifs is enabled by the feasibility of SiO<sub>4</sub> tetrahedra to condense in a manifold of patterns. Whereas oxygen atoms rarely interconnect more than two tetrahedra, in silicate-related structures, nitrogen atoms may even bridge up to four tetrahedra, which for example form star-shaped units and even edge-sharing tetrahedra have been observed. This has led to a multitude of compounds with very diverse structures. The manifold of structures has also been observed when tetrahedra centers were exchanged for P in so-called nitridophosphates. Many of these compounds exhibit promising properties for application in phosphor-converted LEDs.<sup>[4, 5]</sup> Structural diversification can be furthermore enhanced by mixed networks in terms of oxo- and nitridosilicates alike as displayed by alumosilicates and their related nitridoalumosilicates.

Syntheses of nitridosilicates and -phosphates are typically performed at high temperatures (>1000 °C) and, especially for nitriodophosphates often under high pressures (>4 GPa).<sup>[6]</sup> With the recent discovery of SiP<sub>2</sub>N<sub>4</sub>NH, a significant step was taken to accommodate the highpressure motif of SiN<sub>6</sub> octahedra.<sup>[7-9]</sup> Sixfold coordinated Si atoms are also very rare in oxidic compounds where examples include rutile-type stishovite, a high-pressure polymorph of SiO<sub>2</sub>, or K<sub>2</sub>Si[Si<sub>3</sub>O<sub>9</sub>].<sup>[10, 11]</sup> Despite the aforementioned structural diversity, the observation of mineral-analogous nitridosilicates and -phosphates is uncommon if the charge of the counterion is to be preserved. This is simply explained by the high anionic charge of the nitride networks. The incorporation of Si and P allows mitigating the high anionic charge thus enabling the syntheses of mineral analogous compounds like  $AESiP_3N_7$  (AE = Sr, Ba), which crystallize isotypic to the mineral barylite (BaBe<sub>2</sub>Si<sub>2</sub>O<sub>7</sub>).<sup>[12]</sup> Synthetic challenges targeting nitridosilicate phosphates arise from the decomposition of P<sub>3</sub>N<sub>5</sub> at temperatures above 850 °C and the relative chemical inertness of Si<sub>3</sub>N<sub>4</sub>. This problem has been overcome by high partial pressures of HCl (SiP<sub>2</sub>N<sub>4</sub>NH) or employing small amounts of NH<sub>4</sub>F ( $AESiP_3N_7$  (AE = Sr, Ba)) as mineralizing agents and applying high external pressures of 8 GPa at 1100–1700 °C realized by a multianvil press.<sup>[7, 12]</sup> As shown previously, NH<sub>4</sub>F seems to be able to reversibly cleave the bonds in refractory nitrides allowing straightforward synthesis of nitridic compounds.<sup>[12,</sup> <sup>13]</sup> NH<sub>4</sub>F however cannot be found in the reaction products as side reactions with the BN crucible material seem possible.

Following the NH<sub>4</sub>F mineralizer-assisted approach we have now found a simple way to access mica-like layered silicon imidonitridophosphates<sup>†</sup>  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg,  $Mg_{0.94}Ca_{0.06}$ , Ca, Sr) at high-pressure/high-temperature conditions (details are described in Ref. 14) from the respective AE-amide,  $P_3N_5$ , and  $Si_3N_4$  with NH<sub>4</sub>F as a mineralizer according to Equation (4.1). To further investigate luminescence properties,  $Eu^{2+}$ -doped samples have been synthesized by addition of  $\approx 1$  mole% of EuF<sub>3</sub> (with respect to AE) to the starting mixtures.

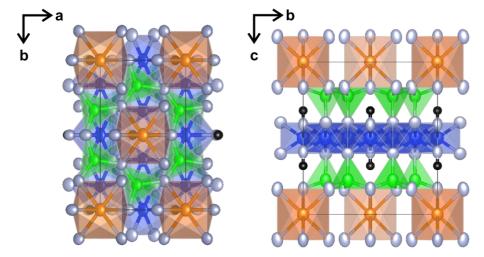
$$3 AE(NH_2)_2 + 4 P_3N_5 + 3 Si_3N_4 \rightarrow 3 AESi_3P_4N_{10}(NH)_2 + 2 NH_3$$
(4.1)

The reactions yielded the title compounds as colorless, microcrystalline powders, which are stable towards air and moisture. Samples of  $AESi_3P_4N_{10}(NH)_2$  ( $AE = Mg_{0.94}Ca_{0.06}$ , Ca, Sr) doped with  $Eu^{2+}$  are yellow. More details on the synthesis are given in the Supporting Information. The crystal structure of  $CaSi_3P_4N_{10}(NH)_2$ , was elucidated by single-crystal diffraction (SCXRD) with microfocused synchrotron radiation (C2/m (no. 12), a = 4.91790(10), b = 8.5159(2), c =9.8267(2) Å,  $\beta = 99.476(3)^{\circ}$ , Z = 2,  $R_1 = 0.0805$ ). For this purpose, pre-characterized crystallites on TEM-grids (Figure C1) were used at beamline ID11 of the ESRF (Grenoble, France).<sup>[15-17]</sup> Data from two twinned crystallites were merged to increase completeness. Structure elucidation from single-crystal X-ray diffraction (SCXRD) enabled Rietveld refinements, also for the compounds with AE = Mg,  $Mg_{0.94}Ca_{0.06}$ , and Sr. The silicon imidonitridophosphates *AE*Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> are isotypic to e.g. clintonite, a brittle mica with composition  $Ca(Mg,Al)_3(Al_3Si)O_{10}(OH)_2.$ <sup>[18]</sup> Single-crystal diffraction patterns of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> show diffuse streaks and signs of twinning, which are typical for mica-type materials.<sup>[19, 20]</sup> Twinning by rotation of 120° around [310] for both crystals was taken into account and the position of H was determined from difference Fourier maps. The N–H bond length was restrained at 0.89 Å.<sup>[21]</sup>

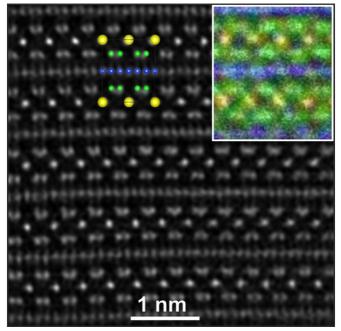
Elemental compositions of the title compounds were confirmed by X-ray spectroscopy (EDX) (Table C6) and phase compositions of respective samples were analyzed by Rietveld refinements (Figures C2–S5, Tables C7–S11). The structures of  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg,  $Mg_{0.9}Ca_{0.1}$ , Ca, Sr) consist of layers of  $AEN_6$  octahedra, PN<sub>4</sub> tetrahedra and SiN<sub>4</sub>(NH)<sub>2</sub> octahedra following the general scheme of tetrahedra–octahedra–tetrahedra (*T-O-T*) arrangement for mica-like structures (Figure 4.1).<sup>[22, 23]</sup>

<sup>&</sup>lt;sup>†</sup> According to IUPAC nomenclature the compounds are called silicon imidonitridophosphate in this reprint.

Interatomic distances P–N range from 1.614(6) to 1.702(5) Å with the latter corresponding to a surprisingly long P–N bond that is comparable to those in compounds like  $Sr_3P_3N_7$  (1.683(11) Å), Mg<sub>2</sub>PN<sub>3</sub> (1.693(5) Å) and  $\beta$ -HP<sub>4</sub>N<sub>7</sub> (1.697(2) Å).<sup>[21, 24, 25]</sup> Bond lengths Si–N range between 1.837(8) and 1.923(7) Å; similar to those reported for the high-pressure compounds  $\gamma$ -Si<sub>3</sub>N<sub>4</sub> (1.8626(1) Å) or SiP<sub>2</sub>N<sub>4</sub>NH (1.8031(9)–2.0146(10) Å).<sup>[7–9]</sup>



**Figure4.1:** Structure of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> with coordination polyhedra of Ca displayed in orange, P green, Si blue, N gray and H black. Displacement ellipsoids are displayed with 99% probability (except for H).

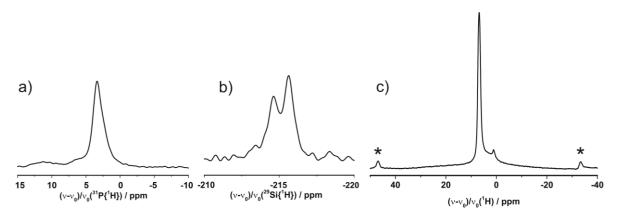


**Figure 4.2:** STEM HAADF image of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> along [100] with structure projection (top middle) and EDX map top right. Ca yellow, P green and Si blue. H and N were omitted for clarity. Further experimental details are given in the Supporting Information.

The title compounds incorporate the alkaline earth metals Mg, Ca and Sr, which results in the lattice parameter *c* varying by  $\approx 0.8$  Å. Synthesis aiming at BaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>, however, yielded BaSiP<sub>3</sub>N<sub>7</sub>. Possibly, Ba cannot be accommodated, which can be explained by the limited space along [100] and [010] compared with micas with high Ba-content like kinoshitalite (BaMg<sub>3</sub>[Al<sub>2</sub>Si<sub>2</sub>O<sub>10</sub>](OH)<sub>2</sub>).<sup>[26]</sup>

As the correct assignment of atom types to crystallographic positions is impeded by the similar X-ray scattering form factors of Si and P, STEM-EDX mappings with atomic resolution were performed (Figure 4.2). The overlay of STEM-EDX maps with an HAADF image shows the ordering of Si and P and led to the conclusion that the title compounds consist of PN<sub>4</sub> tetrahedra and SiN<sub>4</sub>(NH)<sub>2</sub> octahedra.

IR spectra (Figure C17) show absorption bands for each of the title compounds at 3313–3334 cm<sup>-1</sup>, indicating the presence of N–H stretching vibrations.<sup>[27]</sup> The positions of the corresponding H atom were localized from single-crystal diffraction data and confirmed by solid-state NMR as imide groups adjacent to Si, forming SiN<sub>4</sub>(NH)<sub>2</sub> octahedra.



**Figure 4.3:** Solid-state NMR spectra of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> at 20 kHz MAS speed. One signal in the <sup>31</sup>P{1H} (a) and two signals in the <sup>29</sup>Si{1H} (b) spectra agree with the structure model. <sup>1</sup>H NMR reveals one intense signal of the imide group of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> while the weaker one belongs to MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>. Rotation sidebands are marked with asterisks. Full spectra are provided in the Supporting Information.

All signals in <sup>1</sup>H, <sup>31</sup>P, cross-polarized (CP) <sup>1</sup>H $\rightarrow$ <sup>31</sup>P and <sup>1</sup>H $\rightarrow$ <sup>29</sup>Si (indicated by the notation {1H}) <sup>31</sup>P{1H} and <sup>29</sup>Si{1H} MAS NMR spectra of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> shown in Figure 4.3 are consistent with the structure model. Additional weak peaks are attributed to the side phase MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>, which forms by reaction with MgO spacer disks used in the multianvil assembly. The sharp <sup>31</sup>P signal at 3.5 ppm corresponds to the single Wyckoff site and agrees

with NMR data for P in SiP<sub>2</sub>N<sub>4</sub>NH. The persistence of this signal in the <sup>31</sup>P{1H} measurements indicate the vicinity of H to the P site. <sup>29</sup>Si{1H} spectra show two signals with an estimated integral ratio of 1 : 2 centered at 214.6 and 215.6 ppm, respectively. Although the use of integrated intensities is problematic for CP spectra, it is warranted here in good approximation as the average distances of the <sup>29</sup>Si to the four neighboring protons in the structure are similar, see Table C13. The two <sup>29</sup>Si resonances can accordingly be attributed to the Wyckoff sites 2*d* and 4*h* occupied by sixfold coordinated Si, comparable to  $\gamma$ -Si<sub>3</sub>N<sub>4</sub> and SiP<sub>2</sub>N<sub>4</sub>NH with resonances at –225 and –205 ppm, respectively.<sup>[7, 28]</sup> Again, the presence of both signals in cross-polarization experiments indicates the vicinity of H to both Si sites. <sup>1</sup>H NMR shows a strong signal at 6.8 ppm, consistent with H localized above the SiN<sub>4</sub>(NH)<sub>2</sub> layers and centered in the void formed by the PN<sub>4</sub> *sechser* rings. BVS calculations (Table C14) are in agreement with the structure model.<sup>[29]</sup>

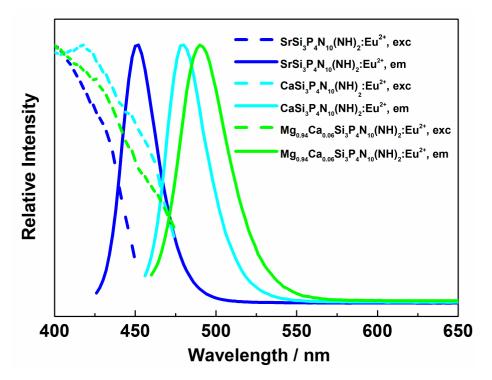
The thermal behavior of  $CaSi_3P_4N_{10}(NH)_2$  was analyzed by temperature-dependent powder X-ray diffraction (PXRD), revealing thermal stability up to 900 °C with exceptionally low thermal expansion of the unit cell volume (Figure C10–S12).

The direct optical band gaps of the undoped title compounds derived from Tauc plots amount to 4.6, 4.2 and 3.9 eV for AE = Mg, Ca, Sr, respectively, with decreasing band gap towards heavier homologs (Figure C18).<sup>[30, 31]</sup> The large band gaps are beneficial concerning luminescence of Eu<sup>2+</sup>-doped samples. The only reported luminescent imidonitride so far is BaP<sub>6</sub>N<sub>10</sub>NH:Eu<sup>2+</sup> with  $\lambda_{max} = 451$  nm and a FWHM of 52 nm (2423 cm<sup>-1</sup>).<sup>[32]</sup> SrSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>:Eu<sup>2+</sup> and CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>:Eu<sup>2+</sup> show narrow emission bands upon excitation with UV light at  $\lambda_{max} = 451$  nm and 478 nm with FWHMs of 26 nm (1300 cm<sup>-1</sup>) and 30 nm (1298 cm<sup>-1</sup>), respectively (Figure 4.4).

MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> showed no luminescence since the smaller size of the coordination polyhedron impedes Eu<sup>2+</sup> incorporation. However, intrigued by the natural solid solution series of micas like the phlogopite–aspidolite series (K(Mg)<sub>3</sub>AlSi<sub>3</sub>O<sub>10</sub>(F,OH)<sub>2</sub>-NaMg<sub>3</sub>AlSi<sub>3</sub>O<sub>10</sub>(OH)<sub>2</sub>),<sup>[33]</sup> we have synthesized the compound Mg<sub>1-x</sub>Ca<sub>x</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>:Eu<sup>2+</sup> (*x*≈0.06). Since the Ca content in this compound is below 0.5 at%, the Ca content was estimated by extrapolation of unit cell volumes between the end members MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> and CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>. This compound showed the most red-shifted, narrow emission of the series at  $\lambda_{max}$  = 492 nm with an FWHM of 35 nm (1444 cm<sup>-1</sup>).

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Low-temperature emission spectra were recorded at 6 K (Figure C19), revealing the zerophonon-line and giving insights into the vibrational modes of the layered crystal structure with an estimated phonon frequency of ca. 430 cm<sup>-1</sup>. The rather high phonon-frequency may explain the strong thermal quenching (Figure C20).



**Figure 4.4:** Emission spectra of  $AESi_3P_4N_{10}(NH)_2$  ( $AE = Mg_{0.94}Ca_{0.06}$ , Ca, Sr) in solid lines. Emission maxima and fwhms are:  $Mg_{0.94}Ca_{0.06}$  492 nm, 35 nm (1444 cm<sup>-1</sup>), Ca 478 nm, 30 nm (1298 cm<sup>-1</sup>) and Sr 451 nm, 26 nm (1300 cm<sup>-1</sup>). Corresponding excitation spectra in dashed lines.

Summarizing, based on the approach of employing NH<sub>4</sub>F as a mineralizing agent, we were able to synthesize the first nitridic analogous mica through HP/HT syntheses. Structure determination was performed by a combination of diffraction of microfocused synchrotron radiation on twinned crystallites, STEM-EDX and solid-state NMR. Eu<sup>2+</sup>-doped samples showed narrow band emission from blue (451 nm) to cyan (492 nm). These findings represent the possibility of mimicking one of the most abundant and important aluminum silicates offering new scope for structural diversity and materials properties of nitrides. We expect that nitridic micas can act as model compounds to investigate the influence of aliovalent substitution of the cations and the influence of mixed anionic frameworks on physical properties such as luminescence and dielectric constants, e.g., by exchange of the imide group against OH groups or fluoride as this compositional range is already observed in natural micas.

### Acknowledgements

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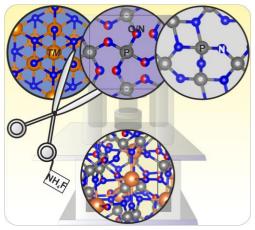
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## High-Pressure Synthesis of Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub> and Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub> by Activation of Binary Nitrides ScN and TiN with NH<sub>4</sub>F

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**Abstract:** Multinary transition metal nitrides and oxonitrides are a versatile and intriguing class of compounds. However, they have been investigated far less than pure oxides. The compounds  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$  have now been synthesized from the binary nitrides ScN and TiN, respectively, by following a high-pressure high-temperature approach at 8 GPa and 1400 °C. NH<sub>4</sub>F acts as a mineralizing agent that



supports product formation and crystallization. The starting materials ScN and TiN are seemingly an uncommon choice because of their chemical inertness but, nevertheless, react under these conditions.  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$  crystallize isotypically with  $Ti_5B_{12}O_{26}$ , consisting of solely vertex-sharing  $P(O/N)_4$  tetrahedra forming two independent interpenetrating diamond-like nets that host  $TM(O/N)_6$  (TM = Sc, Ti) octahedra.  $Ti_5P_{12}N_{24}O_2$  is a mixed-valence compound and shows ordering of  $Ti^{3+}$  and  $Ti^{4+}$  ions.

## 5.1 Introduction with Results and Discussion

A plethora of nitridophosphates was discovered in the last decades. However, only few examples containing transition metals have been described so far which is rather intriguing bearing the ongoing research on transition metal phosphates and their applications in mind.<sup>[1-5]</sup> Oxonitridophosphates like *M*<sup>1</sup>*M*<sup>11</sup>P<sub>3</sub>O<sub>9</sub>N (*M*<sup>1</sup> = Na, K, *M*<sup>11</sup> = Al, Ga, Cr, Fe, Mn) were the first compounds paving the way for transition metal (oxo)nitridophosphates.<sup>[6,7]</sup> By using high-pressure conditions accessible by Walker-type multianvil presses, compounds like MnP<sub>2</sub>N<sub>4</sub>, CdP<sub>2</sub>N<sub>4</sub> and Zn<sub>2</sub>PN<sub>3</sub> were synthesized. The metal-containing starting materials were either the *TM* azides (Cd), metal powders (Cd, Mn), or binary nitride (Zn).<sup>[8-11]</sup> Despite these achievements, no systematic access to this class of compounds was reported as illustrated by

the different starting materials. Common chemicals such as azides and nitrides are either too sensitive and detonate on impact or are too inert, respectively. The first systematic approach to a variety of *TM* compounds is represented by metathesis reactions under high-pressure high-temperature conditions as shown for  $MP_8N_{14}$  (M = Fe, Co, Ni),  $M_{1-x}PO_{3+4x}N_{1-4x}$  ( $x\approx 0.05$ ),  $M_{0.75}PO_4$  (M = Zr, Hf) and Hf<sub>9-x</sub>P<sub>24</sub>N<sub>52-4x</sub>O<sub>4x</sub> ( $x\approx 1.84$ ).<sup>[12-15]</sup>

High-pressure metathesis with the formation of lithium halides from LiPN<sub>2</sub> and *TM* halides, however, features inherent drawbacks. Like many HP/HT syntheses, this approach typically yields microcrystalline powders which impede precise structure elucidation by conventional single-crystal X-ray diffraction. The addition of Li can also pose a problem if the latter is incorporated unintentionally, especially in presence of heavy metal atoms. This combination is rather problematic in terms of structure determination by X-ray methods as it is hard to distinguish between heavy metal vacancies or occupational disorder of the heavy metal and Li.

To facilitate crystal growth of nitridophosphates in high-pressure high-temperature reactions, the addition of small amounts of NH<sub>4</sub>Cl to the mixture of starting materials as mineralizer proved to be successful. NH<sub>4</sub>Cl most likely aids in reaction and crystal growth by reversible P–N bond cleavage and formation.<sup>[16]</sup> As shown for group 1, 2 and 13 elements, this kind of activation seems sufficient to provide access to a number of different structures and compounds.<sup>[17-19]</sup> As previously shown for  $AESiP_3N_7$  (AE = Sr, Ba), the capability of NH<sub>4</sub>Cl is limited as some starting materials such as  $Si_3N_4$  are less reactive and the decomposition, as well as activation temperatures of the starting materials, may differ significantly; for example, thermal decomposition of P<sub>3</sub>N<sub>5</sub> may compete with the activation of Si<sub>3</sub>N<sub>4</sub>.<sup>[20]</sup> The question arose if previous findings on the stability of transition metal nitrides remain valid if the mineralizing agent is changed. The starting materials ScN and TiN were chosen because of their refractory character and notorious inertness. For example, TiN is used as a wear- and heat-resistant coating for drill heads.<sup>[21]</sup> Etching of TiN films is heavily investigated in semiconductor fabrication. Solutions for etching consisting of a source of hydrogen fluoride and an oxidizing agent have proven to be successful in reducing TiN film thickness. The oxidizing agent in these cases is needed to oxidize Ti<sup>3+</sup> to Ti<sup>4+</sup> which forms volatile TiF<sub>4</sub>, in contrast to non-volatile TiF<sub>3</sub>.<sup>[22]</sup> As the volatility of the fluoride is not an issue in HP/HT syntheses, and we assumed a similar behavior for ScN, we explored the capability of NH<sub>4</sub>F as

a safe and convenient HF source in reactions yielding the transition-metal oxonitridophosphates  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$ .

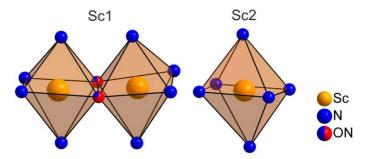
Both compounds were synthesized at 1400 °C and 8 GPa by using a modified Walker-type multianvil apparatus.<sup>[23]</sup> Reactions followed the so-called nitride-route using  $P_3N_5$  and *TM*N (*TM* = Sc, Ti) with additional NH<sub>4</sub>F as a mineralizing agent [Eqs. (5.1) and (5.2)]. PON or TiO<sub>2</sub> were added as oxygen sources. The synthesis of Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub>, however, was performed with only 0.8 equivalents of TiO<sub>2</sub> as stoichiometric amounts lead to microcrystalline samples and split reflections in the PXRD possibly due to excess oxygen from surface hydrolysis of the BN crucible.

$$5 \text{ ScN} + 3 P_3 N_5 + 3 \text{ PON} \rightarrow \text{Sc}_5 P_{12} N_{23} O_3$$
(5.1)

$$4 \operatorname{TiN} + 4 \operatorname{P_3N_5} + \operatorname{TiO_2} \to \operatorname{Ti_5P_{12}N_{24}O_2}$$
(5.2)

Both reactions yielded moisture- and air-resistant crystalline powders with gray and black color for  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$ , respectively. More detailed information on the synthesis is given in the Supporting Information.

Structure elucidation was performed by single-crystal X-ray diffraction (SCXRD). Deposition Numbers 2084626 (for  $Sc_5P_{12}N_{23}O_3$ ) and 2084627 (for  $Ti_5P_{12}N_{24}O_2$ ) contain the supplementary crystallographic data for this paper. These data are provided free of charge by the joint Cambridge Crystallographic Data Centre and Fachinformationszentrum Karlsruhe. Both compounds are isotypic and crystallize in space group  $I4_1/acd$  (no. 142) with Z = 8 (Table 5.1). Additional Rietveld refinements indicate that the title compounds are the main constituents of the obtained samples (Figures D2 and D3, Table D9). Both compounds are isotypic to  $Ti_5B_{12}O_{26}$ .<sup>[24]</sup> The transition metals are coordinated octahedrally by O/N. The  $TM1(O/N)_6$  octahedra share one common edge and form pairs in contrast to the  $TM2(O/N)_6$  octahedra, which are not condensed (Figure 5.1).

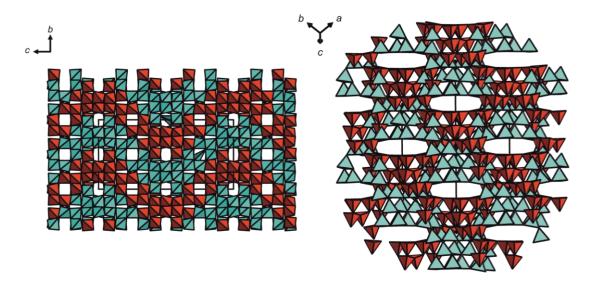


**Figure 5.1.** Edge-sharing ScO<sub>0.75</sub>N<sub>5.25</sub> and isolated ScN<sub>6</sub> octahedra.

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formula	Sc <sub>5</sub> P <sub>12</sub> N <sub>23</sub> O <sub>3</sub>	Ti <sub>5</sub> P <sub>12</sub> N <sub>24</sub> O <sub>2</sub>
molar mass /g⋅mol <sup>-1</sup>	966.67	979.38
crystal system	tetragonal	
space group	<i>I</i> 4 <sub>1</sub> / <i>acd</i> (no. 142)	
lattice parameters / Å	a = 12.3598(2)	a = 12.1214(2)
	c = 24.0151(4)	c = 23.8458(5)
cell volume / ų	3668.66(13)	3503.62(14)
formula units/ unit cell	8	
density / g·cm <sup>3</sup>	3.500	3.713
μ / mm <sup>-1</sup>	2.893	3.369
temperature / K	296(2)	298(2)
absorption correction	semiempirical	
radiation	Mo- $K_{\alpha}$ ( $\lambda = 0.71073$ Å)	
F(000)	3760	3792
hetarange / °	3.3 - 36.3	2.9 - 33.1
total no. of reflections	66300	52537
Indep. reflections [I≥2σ(I) / all]	1971 / 2226	1567 / 1676
$R_{\sigma}, R_{ m int}$	0.0158. 0.0616	0.0108, 0.0353
refined parameters	100	99
Goodness of fit	1.187	1.432
<i>R</i> -values [I $\geq 2\sigma(I)$ ]	$R_1 = 0.0210$	$R_1 = 0.0222$
	$wR_2 = 0.0474$	$wR_2 = 0.0680$
R-values (all data)	$R_1 = 0.0264$	$R_1 = 0.0240$
	$wR_2 = 0.0489$	$wR_2 = 0.0690$
$\Delta ho_{ m max}$ , $\Delta ho_{ m min}$ / e·Å <sup>-3</sup>	0.552 / -0.715	0.921 / -1.636

**Table 5.1.** Crystallographic data for the single-crystal structure refinements of  $Sc_5P_{12}N_{23}O_3$  and $Ti_5P_{12}N_{24}O_2$  Standard deviations are given in parentheses.



**Figure 5.2** Structure projections viewed from [100] (left) and [111] (right). The different diamond-like nets of tetrahedra are colored in red or blue, respectively. *TM* atoms are omitted for clarity.

The  $P(O/N)_4$  network consists of edge-sharing tetrahedra forming truncated, hollow supertetrahedra which in turn form two separate interpenetrating networks (Figure 5.2). As differentiation between O and N is not unambiguous from X-ray data, and ordering of O and N is not possible in terms of charge neutrality, mixed occupied sites were assigned by considering a combination of bond lengths, bond valence sums and charge neutrality. Following this combination, the position 05/N5 and N2 in  $Sc_5P_{12}N_{23}O_3$  both showed potential mixed occupation according to BVS but position 05/N5 was considered more likely because of the short P-O/N bond lengths of 1.5722(9) Å and N2 forming unusually large Sc-N distances explained in the following. In  $Ti_5P_{12}N_{24}O_2$  assigning mixed occupied positions led to positions 02/N2, 05/N5 and 06/N6 fully consistent with BVS calculations as only these positions showed deviations from the ideal value of 3 for N. Distances *TM*–O/N are between 2.1421(9)-2.4424(9) Å for Sc-O/N and 2.0435(11)-2.2911(11) Å for Ti-O/N (Tables D6 and D7). Both compounds exhibit rather large interatomic distances concerning the *TM*2 position. In the equatorial plane of the  $TM2(O/N)_6$  octahedron, the distances Sc2-N2 (2.4424(9) Å) deviate quite significantly from those in ScN (2.25 Å).<sup>[25]</sup> The Ti2 position exhibits distances  $Ti-(O_{0.15}N_{0.85})$  of 2.2911(11) Å which are also larger than expected with respect to TiN or rutile-type TiO<sub>2</sub>, where they are 2.13 and 1.81–2.06 Å, respectively.<sup>[26,27]</sup> The long interatomic distances result in a "rattling" effect which is reflected by low BVS values for these positions

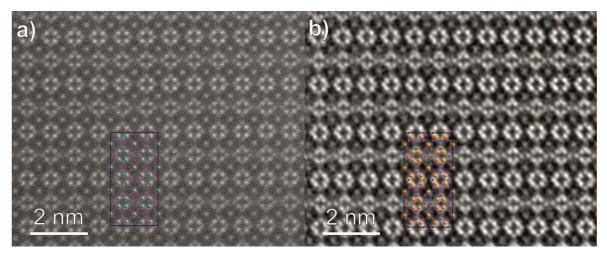
and their rather large and prolate displacements ellipsoids (Figure D1). The BVS values are 2.42 for Sc2 and 2.64 for Ti2 instead of the expected values of 3 for Sc and 4 for Ti2.

The chemical compositions of  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$  are supported by energydispersive X-ray spectroscopy (Table D8) and FTIR spectra indicate the absence of N–H groups (Figure D4). Scanning transmission electron microscope high-angle annular dark-field (STEM-HAADF) images correspond well to structure projections (Figure 5.3).

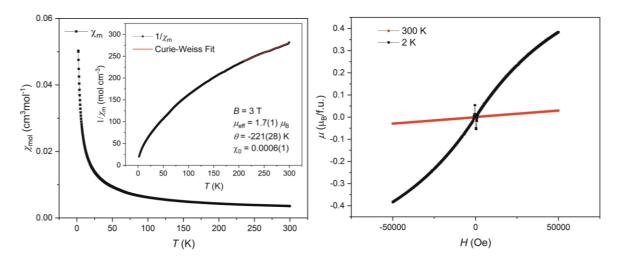
In order to confirm the presence of both Ti<sup>3+</sup> and Ti<sup>4+</sup> in Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub>, UV-Vis absorption spectra and magnetic measurements were performed. UV-Vis spectra show a broad absorption band centered at 420 nm, which corresponds well with the presence of Ti<sup>3+</sup> (Figure D6).<sup>[28]</sup> Figure 5.4 shows the magnetic susceptibility with a linear paramagnetic behavior for temperatures >200 K. The fit indicates an effective paramagnetic moment of  $\mu_{eff}$ = 1.7(1)  $\mu_B$  for each Ti cation on Wyckoff position 32 *g* assuming Ti<sup>4+</sup> on Wyckoff position 8*b*. This is in agreement with the theoretical spin-only value for Ti<sup>3+</sup> with  $\mu_{eff}$  = 1.73  $\mu_{B}$ .<sup>[29]</sup> Below 200 K, the curve deviates from Curie–Weiss behavior with a continuously increasing magnetic moment. Magnetization isotherms at 300 K with purely paramagnetic behavior and with a small saturation effect at 2 K towards higher fields are shown in the bottom part of Figure 5.4.

In analogy to  $Ti_5B_{12}O_{26}$ , we could assign crystallographic sites for  $Ti^{3+}$  and  $Ti^{4+}$  according to the presence of both valence states.<sup>[24]</sup> Although the similarity of the ordering of  $Ti^{3+}$  and  $Ti^{4+}$  is quite remarkable given the fact that the composition of the tetrahedral network is completely altered.

The optical bandgap of  $Ti_5P_{12}N_{24}O_2$  was approximated by converting reflectance spectra to the Kubelka-Munk function and calculating a Tauc plot under the assumption of a direct bandgap. Linear regression between 2.1 and 3.1 eV was used to determine the inflection point yielding a bandgap of about 1.6 eV.<sup>[30,31]</sup> In the same way, the optical bandgap of  $Sc_5P_{12}N_{23}O_3$ was estimated by linear regression between 4.0 and 4.5 eV and determined as a direct bandgap of 3.8 eV (Figure D7).



**Figure 5.3.** STEM-HAADF images of  $Ti_5P_{12}N_{24}O_2$  (a) and  $Sc_5P_{12}N_{23}O_3$  (b) with structure projections as overlays (Ti blue, Sc orange, P violet, O/N omitted for clarity



**Figure 5.4.** Magnetic susceptibility of  $Ti_5P_{12}N_{24}O_2$  and inverse magnetic susceptibility (inset) with an extended Curie-Weiss fit (red) (left) and Magnetic isotherm (right) of  $Ti_5P_{12}N_{24}O_2$  at 300 (red) and 2 K (black) (bottom). At low field strengths, it reveals a small superconducting impurity of most likely TiN ( $T_c = 5.6$  K).<sup>[5]</sup>

#### Conclusion

The oxonitridophosphates  $Ti_5P_{12}N_{24}O_2$  and  $Sc_5P_{12}N_{23}O_3$  have been synthesized with NH<sub>4</sub>F as a mineralizing agent by high-pressure high-temperature synthesis. The use of NH<sub>4</sub>F as a mineralizing agent seems promising to access a variety of yet unknown compounds. The mechanistic reasons for the apparent superiority of NH<sub>4</sub>F in contrast to NH<sub>4</sub>Cl, still need to be investigated and will certainly be interesting. The presence of  $Ti^{3+}$  and  $Ti^{4+}$  in  $Ti_5P_{12}N_{24}O_2$  was

derived from crystal-chemical considerations and confirmed by magnetic measurements as well as UV-Vis spectroscopy.

The easy accessibility of *TM* oxonitridophosphates from binary nitrides constitutes a significant progress in the exploration of multinary nitride network structures. This, in turn, is now offering a new scope of properties for nitridophophates like magnetic properties for spintronic applications as diluted magnetic semiconductors, for example, by substoichiometric doping with *TM*.<sup>[33]</sup> The incorporation of *TM* in nitridophophates also allows the metal cations to be mixed-valent, as shown for Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub>, which could also be formulated as Ti<sup>III</sup><sub>4</sub>Ti<sup>IV</sup>P<sub>5</sub>N<sub>24</sub>O<sub>2</sub> for clarification. This leads to significantly smaller bandgaps compared to group 1 or group 2 nitridophosphates, which, so far, result in ultrawide-bandgap semiconductors. Research on Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub> regarding solar power harvesting and photocatalysis seems intriguing judging by its bandgap of 1.6 eV and the fact that TiO<sub>2</sub> and TiN are both intensively investigated in those fields.<sup>[34-37]</sup>

The next obstacle to overcome are syntheses of oxygen-free *TM* nitrides, by employing NH<sub>4</sub>F and the binary *TM* nitrides. Thus, delivering straightforward access to a group of multinary nitrides which are, despite their projected potential applications, scarcely investigated.

#### Acknowledgements

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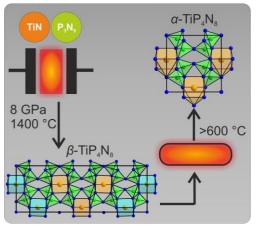
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# 6 Discovery of Two Polymorphs of TiP<sub>4</sub>N<sub>8</sub> Synthesized from Binary Nitrides

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**Abstract:**  $TiP_4N_8$  was obtained from the binary nitrides TiN and  $P_3N_5$  upon addition of  $NH_4F$  as a mineralizer at 8 GPa and 1400 °C. An intricate interplay of disorder and polymorphism was elucidated by in situ temperature-dependent singlecrystal X-ray diffraction, STEM-HAADF, and the investigation of annealed samples. This revealed two polymorphs, that consist of dense networks of PN<sub>4</sub>

tetrahedra (degree of condensation  $\kappa = 0.5$ ) and either augmented triangular TiN<sub>7</sub> prisms or triangular TiN<sub>6</sub> prisms for  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>, respectively. The structures of TiP<sub>4</sub>N<sub>8</sub> exhibit bodycentered tetragonal (bct) framework topology. DFT calculations confirm the measured bandgaps of  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (1.6–1.8 eV) and predict the thermochemistry of the polymorphs in agreement with the experiments.

### 6.1 Introduction with Results and Discussion

A variety of group 1 and 2 nitridophosphates has been characterized over the last decades with silicate-related structural motifs like discrete tetrahedra, layers, or frameworks.<sup>[1]</sup> An explanation for this abundance of group 1 and 2 nitridophosphates is, that P<sub>3</sub>N<sub>5</sub>, at ambient pressure, readily decomposes above 850 °C under the evolution of N<sub>2</sub>.<sup>[2]</sup> To suppress this degradation, high-pressure high-temperature (HP/HT) synthesis emerged as a viable pathway, following Le Chatelier's principle. The use of stable azides of alkali and alkaline earth elements proved to be crucial for syntheses. The amount of  $N_2$  from the decomposition of these azides further prevents the decomposition of  $P_3N_5$ .

The incorporation of transition metals (TM) in nitridic tetrahedral framework structures constitutes a less explored field of research compared to compounds containing group 1 and 2 elements. *TM* bearing compounds are especially interesting given the opportunity to find stable nitrides with suitable band gaps for semiconductor applications like photovoltaics or photocatalysis.<sup>[3–5]</sup> Although many ternary nitrides with outstanding properties have been predicted by theorists, the synthetic limits still restrict experimental confirmation of these claims.<sup>[6,7]</sup> Only on a few occasions, the azide-route could be transferred to the synthesis of transition metal nitridophosphates as the respective *TM* azides are either nonexistent or pose serious safety issues due to being explosive.<sup>[8]</sup> The nitride route, employing the transition metal nitride, was dismissed on many occasions as numerous TM nitrides such as ScN, TiN and VN are refractory and unreactive. The requirement to prevent the oxidation of the nitride ion and reduction of P or the transition metal constitutes a problem that becomes more severe when the oxidation state of the transition metal increases. Two versatile approaches to gain access to ternary transition metal nitrides have recently been described <sup>[1,9,10]</sup> and expanded the compositional range of nitridosilicates and -phosphates beyond the well-investigated group 1 and 2 compounds. The cation-exchange approach relies on preformed nitridic networks. For instance, Ca<sub>2</sub>Si<sub>5</sub>N<sub>8</sub> reacts with a FeCl<sub>2</sub> melt to yield Fe<sub>2</sub>Si<sub>5</sub>N<sub>8</sub> and CaCl<sub>2</sub>.<sup>[11]</sup> Solidstate metathesis of nitridophosphates employs a metal halide and LiPN<sub>2</sub> with the formation of lithium halide as a driving force.<sup>[12]</sup> Both pathways circumvent the formation of stable transition metal phosphides. However, the cation-exchange approach is often limited to the respective nitridic network used in this top-down approach. Solid-state metathesis can be hindered by substantial amounts of Li present. The incorporation of Li beside a heavy transition metal can pose serious problems in terms of structure elucidation.

In our previous studies on the NH<sub>4</sub>F mediated HP/HT synthesis, NH<sub>4</sub>F has shown its ability to overcome the differing reactivities of Si<sub>3</sub>N<sub>4</sub> and P<sub>3</sub>N<sub>5</sub> resulting in mixed nitridic networks like the recently described nitridic barylite- ( $AESiP_3N_7 AE = Sr$ , Ba) as well as mica-type compounds ( $AESi_3P_4N_{10}(NH)_2 AE = Mg$ ,  $Mg_{0.94}Ca_{0.06}$ , Ca, Sr).<sup>[9,13]</sup> In a similar fashion we employed P<sub>3</sub>N<sub>5</sub> and TiN as starting materials although TiN is used as a heat and chemically resistant coating material<sup>[14]</sup> and despite the tendency of group 4 cations to necessitate the

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incorporation of oxygen in oxonitridophosphates.<sup>[14,15]</sup> One may also question if N<sub>2</sub> is a sufficiently strong oxidizing agent at conditions achievable with a multianvil press. Diamond-anvil cell (DAC) experiments resulting in the Ti<sup>IV</sup>-compounds Ti<sub>3</sub>N<sub>4</sub> and TiN<sub>2</sub> from TiN and N<sub>2</sub> had been conducted at 73–75 GPa and 2400 K.<sup>[16,17]</sup> The oxidizing behavior of nitrogen at elevated pressure could originate from a weakening of the triple bond in N<sub>2</sub> as signaled by the existence of polymeric N exhibiting the crystal structure of black phosphorus or non-molecular N polymorphs.<sup>[18-20]</sup>

Two polymorphs of the first  $Ti^{IV}$  nitridophosphate were discovered by HP/HT synthesis employing  $P_3N_5$  and TiN as starting materials with the addition of  $NH_4F$  as a mineralizing agent and  $NH_4N_3$  as nitrogen source (eq. 6.1) at 8 GPa and 1400 °C and subsequent annealing at 700°C in fused silica ampoules.

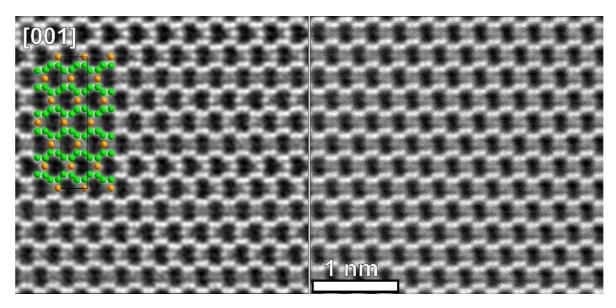
$$24 \operatorname{TiN} + 32 \operatorname{P_3N_5} + 3 \operatorname{NH_4N_3} \rightarrow 24 \operatorname{TiP_4N_8} + 4 \operatorname{NH_3}$$
(6.1)

The HP/HT conditions were achieved by employing a modified Walker-type multianvil press.<sup>[21-23]</sup> The underlying reason for the success of NH<sub>4</sub>F seems to be the reversible elementnitrogen bond cleavage for refractory nitrides that may result in unstable molecular intermediates. In the case of TiN, HF is industrially exploited to convert TiN to the respective Ti fluorides in etching processes.<sup>[24]</sup> There is no NH<sub>4</sub>F with the reaction products, as side reactions with BN, the crucible material, probably occur. HP/HT synthesis yielded  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (*Pmn*2<sub>1</sub>, *a* = 22.9196(5), *b* = 4.58800(10), *c* = 8.0970(2) Å, *Z* = 6, *R*<sub>1</sub> = 0.0221) as amber crystals (Table 6.1, E2, E6–8).<sup>[25]</sup> Ti atoms are coordinated sixfold in TiN<sub>6</sub> triangular prisms with bond lengths *d*<sub>Ti-N</sub> ranging from 2.086(3)–2.241(2) Å. The network of PN<sub>4</sub> tetrahedra features all-side vertex-sharing tetrahedra with bond lengths *d*<sub>P-N</sub> 1.5998(15)–1.6608(16) Å. The topology of the tetrahedral network can be described as *bct* topology with point symbol 4.6<sup>5</sup> as determined by TOPOS, exhibiting *vierer*, *sechser*, and *achter* rings, according to the nomenclature introduced by Liebau, and a degree of condensation  $\kappa = 0.5$  (Figure E1-E2).<sup>[26–28]</sup>

Formula	$\alpha$ -TiP <sub>4</sub> N <sub>8</sub>	$\beta$ -TiP <sub>4</sub> N <sub>8</sub>
molar mass (g·mol <sup>-1</sup> )	283.86	
crystal system	orthorhombic	
space group	<i>Pmn</i> 2 <sub>1</sub> (no. 31)	
lattice parameters (Å)	a = 7.6065(2)	a = 22.9196(5)
	b = 4.63320(10)	b = 4.58800(10)
	c = 7.8601(3)	c = 8.0970(2)
cell volume (ų)	277.009(14)	851.44(3)
formula units/ unit cell	2	6
density (g⋅cm <sup>-3</sup> )	3.403	3.322

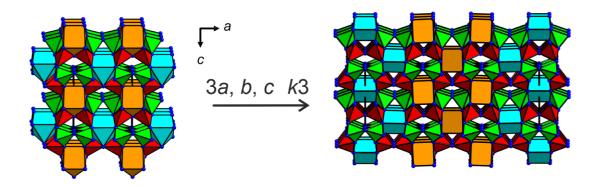
Table 6.1: Selected crystallographic information for  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> from SCXRD data.

Structure elucidation of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> by SCXRD led to occupational disorder on Ti sites, an issue that was encountered on datasets of several crystals. STEM-HAADF investigations showed different cation-site occupations for different crystallites ranging from nearly ordered to severely disordered (Figure 6.1, Table E2, E9–11). *In situ* HTSCXRD investigations showed that Ti disorder was eliminated by heating to 600 °C, which became apparent after accounting for limited short-range order by using separate scale factors for reflections with *h* = 3*n* and *h*  $\neq$  3*n* (Table E2, E12–14).

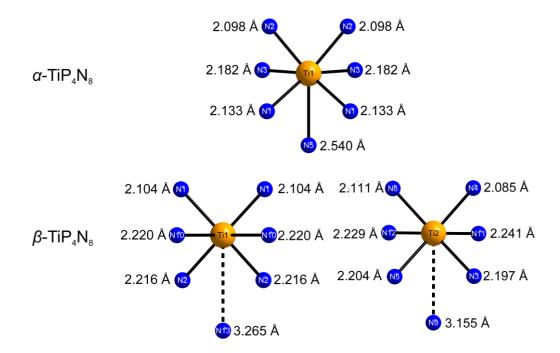


**Figure 6.1:** STEM HAADF images of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> before heat treatment along [001]. Almost ordered crystallite on the left and severely disordered on the right. Structure projection of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> with Ti orange, P green, N omitted for clarity.

Cation ordering was also shown *ex situ* with samples annealed in silica ampoules at 600 °C. Bond lengths  $d_{\text{Ti-N}}$  and  $d_{\text{P-N}}$  are 2.097(4)–2.268(3) Å and 1.605(3)–1.645(3) Å, respectively, at600 °C and change, as well as bonding angles, only slightly compared to ambient temperatures indicating the rigidity of the network. However, products annealed at 700 °C showed a drastic decrease in crystallite sizes and a change of color (Figure E3).



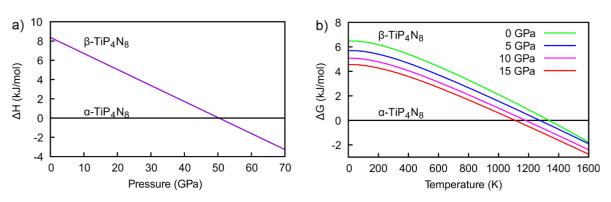
**Figure 6.2:** Structure projections of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (left) and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (right) along ~[010]. PN<sub>4</sub> tetrahedra in green and red and TiN<sub>x</sub> polyhedra in orange and blue, respectively. Different coloring of polyhedra shows different orientations.



**Figure 6.3:** TiN<sub>7</sub> polyhedron in  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (top) and both TiN<sub>6</sub> triangular prisms in  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (bottom): Ti-N bond lengths given next to the corresponding N atom. Dashed Ti-N bonds in  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> correspond to the capping atom in  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>.

Ruby-red crystals of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (*Pmn*2<sub>1</sub>, *a* = 7.6065(2), *b* = 4.63320(10), *c* = 7.8601(3) Å, *Z* = 2,  $R_1 = 0.0256$ ) were isolated from a sample that was annealed at 700 °C. The structure elucidated by SCXRD (Table E2, E3-5)<sup>[25]</sup> is related to that of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> by a cell transformation of 3*a*, *b*, *c* (Table 6.1, Figure 6.2 and E4). This also became evident from lattice parameters and reciprocal lattice sections. The structure of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> strongly resembles that of pseudoorthorhombic svyatoslavite (CaAl<sub>2</sub>Si<sub>2</sub>O<sub>8</sub>, space group P2<sub>1</sub>).<sup>[29,30]</sup> As compared to  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>,  $\alpha$ - $TiP_4N_8$  features no superstructure reflections and a smaller unit cell volume. The orientation of the TiN<sub>x</sub> polyhedra is different (Figure 6.2). Ti is coordinated sevenfold to form  $TiN_7$ augmented triangular prisms enabled by a deformation of the *achter* rings. This leads to a shorter Ti–N distance of 2.540(4) Å (compared to 3.264(4) Å in  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>). The bond lengths in  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> are  $d_{\text{Ti-N}}$  2.098(2)–2.540(4) Å and  $d_{\text{P-N}}$  1.5758(15)–1.644(2) Å (Figure 3). All observed interatomic distances  $d_{P-N}$  agree well with those in comparable compounds.<sup>[31,32]</sup> To the best of our knowledge, there is no SCXRD data of a purely nitridic compound featuring  $Ti^{IV}$  on a cation site in literature to compare the bonding distances  $d_{Ti-N}$  of  $TiP_4N_8$  to. Compositional analyses of bulk samples, performed by Rietveld refinements, show TiP<sub>4</sub>N<sub>8</sub> as the main constituent of all investigated samples. Minor impurities result from black P, resulting from decomposed  $P_3N_5$ , and *h*-BN, which is used as the crucible material (Figure E5– 6, Table E16). Energy-dispersive X-ray spectroscopy confirms the composition of Ti:P:N of 1:4:8 supporting the oxidation state of +IV for Ti (Table E15). IR spectroscopy showed no N-H valence modes (Figure E10). BVS calculations of the polymorphs show almost expected values for the bond-valence sums of the individual atom sites for  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>. For  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>, the lower coordination numbers lead to smaller values for Ti as well as the affected N sites in contrast to the P sites which again show almost regular values (Table E17-18). HTPXRD shows thermal stability up to at least 950 °C in Ar atmosphere (Figure E7). HTPXRD with two subsequent heating cycles showed no structural change once  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is converted to  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (Figure E8)

DFT calculations of the electronic properties identify both  $TiP_4N_8$  polymorphs as semiconductors with indirect bandgaps of 1.35–1.96 eV ( $\alpha$ ) and 1.21–1.73 eV ( $\beta$ ), depending on the exchange-correlation functional used. Figure E12 shows exemplary band structure and



density of state plots. Experimental values from reflectance spectra (Figure E11)<sup>[33]</sup> lead to indirect bandgaps of 1.8 and 1.6 eV, respectively, in agreement with the calculated values.

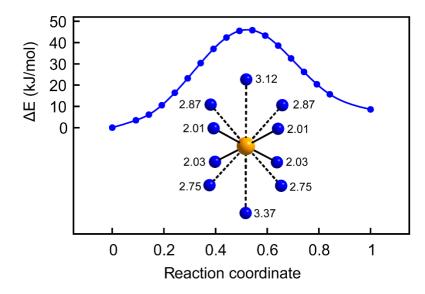
**Figure 6.4:** (a) Enthalpy difference  $\Delta H$  between  $\beta$ - and  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> with increasing pressure. (b) Temperature dependence of the Gibbs free energy difference  $\Delta G$  at different pressures.

Given the synthesis of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> under HP/HT conditions and subsequent conversion to  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> by annealing, it seems obvious that  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is a metastable high-pressure modification. Calculations of the pressure-dependent enthalpy-difference  $\Delta H$  indicate that  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is only about 8 kJ/mol less stable than  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>. However, the pressure dependence of  $\Delta H$  is very small and the  $\beta$ -modification is predicted to form only at pressures around 50 GPa (Figure 6.4a), much higher than the synthesis pressure of 8 GPa. To resolve this discrepancy, we have calculated the pressure and temperature dependence of the free energy G(T,p) using the quasi-elastic approximation (QHA) method.<sup>[34,35]</sup> For that purpose, the phonon contributions to the free energy  $F_{ph}$  at different volumes were calculated, and combined with the electronic energy U(V) to obtain the Gibbs free energy using the expression<sup>[36]</sup>:

$$G(T,p) = \frac{\min}{V} \left[ U(V) + F_{ph}(T,V) + pV \right]$$
<sup>(2)</sup>

Figure 6.4b shows the temperature dependence of the Gibbs free energy at different pressures, which reveals that  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is stable above about 1300 K even at zero pressure and that the transition shifts to lower temperatures with increasing pressure. This matches the experimental results much better and predicts that  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is a quenched high-temperature rather than a high-pressure phase. Finally, we calculated the energy barrier between the polymorphs using the climbing nudged elastic band (Cl-NEB) method.<sup>[37]</sup>

Figure 6.5 shows the NEB energy profile and a cutout of the structure at the transition state. Titanium has four short (2.01–2.03 Å) and four long (2.75–2.87 Å) Ti–N contacts. The position is almost symmetric between the two distant nitrogen atoms (3.12–3.37 Å), one of which becomes the capping atom of the TiN<sub>7</sub> polyhedron in  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (Figure 3).



**Figure 6.5:** Energy barrier of the  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> to  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> transition and the coordination of titanium at the transition state with  $d_{Ti-N}$  (in Å) adjacent to N atoms.

Summing up, TiP<sub>4</sub>N<sub>8</sub> was synthesized by HT/HP synthesis with NH<sub>4</sub>F as a mineralizing agent. A combination of STEM-HAADF, *in situ* HTSCXRD investigations and annealing experiments revealed two polymorphs of TiP<sub>4</sub>N<sub>8</sub>. The occupational disorder of Ti almost vanished at 600 °C resulting in  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>. Further heating leads to a transformation to  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>. DFT calculations including temperature-dependent phonon contributions indicate that  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is a metastable high-temperature polymorph. The presented synthesis of the first Ti<sup>IV</sup>-nitridophosphate from binary nitrides marks a starting point to straightforward syntheses of a multitude of ternary transition metal nitrides. Given the band gaps of  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> both compounds absorb light in the visible part of the spectrum. Band gaps are comparable to those of CdTe, 1.513 eV, and GaAs, 1.423 eV.<sup>[38,39]</sup> It is most likely that the NH<sub>4</sub>F-assisted approach can be adapted for other *TM* nitrides as well, diversifying the compositional range and variability of crystal structures in nitrides, significantly. The present results indicate that the oxidative potential of N in the pressure and temperature range of a multianvil press is sufficient to prevent the formation of *TM* phosphides from reactions of the binary nitrides.

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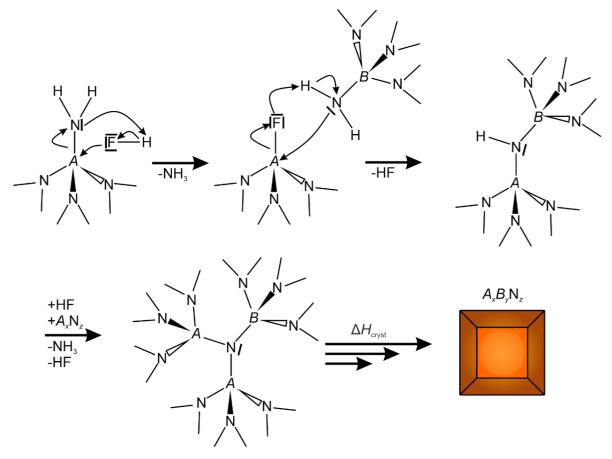
## 7 Discussion and Outlook

### 7.1 Novel Nitrides

A multitude of synthetic pathways has been developed before this thesis to tackle synthesisrelated obstacles in the high-pressure synthesis of nitridophosphates. The advent of the azide route marked systematic access to group 1 and 2 nitridophosphates. The use of NH<sub>4</sub>Cl as a mineralizer reliably yielded single crystals suitable for in-house diffractometers.<sup>[1-3]</sup> The highpressure solid-state metathesis where *RE* or *TM* halides and LiPN<sub>2</sub> form the respective *RE*- or *TM*-nitridophosphate and the Li halide was the first systematic access to these nitridophosphates.<sup>[4-8]</sup> Although this method proved successful, quantitative amounts of Li are present in the starting materials, which may lead to undesired incorporation of Li. Mixed nitridic networks have become accessible through the acid-assisted approach by molecular (PNCl<sub>2</sub>)<sub>3</sub>.<sup>[9,10]</sup> This method, however, involves the release of large amounts of HCl, which enables reactions with the BN crucible and did yet not deliver non-neutral network compounds.

In this context, NH<sub>4</sub>F mediated HP/HT synthesis can be viewed as the synergism of all the aforementioned. Similar to NH<sub>4</sub>Cl, NH<sub>4</sub>F yielded single crystals suitable for laboratory diffractometers as shown by *AE*SiP<sub>3</sub>N<sub>7</sub> (*AE* = Sr, Ba), Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub>, Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub>, and  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>.

The syntheses of nitrdosilicatephosphates  $AESiP_3N_7$  (AE = Sr, Ba) and silicon imidonitridophosphates  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg,  $Mg_{0.94}Ca_{0.06}$ , Ca, Sr) with NH<sub>4</sub>F as a mineralizer yielded anionic mixed nitridic networks. In contrast to the acid-assisted route, no excessive amounts of HF are needed and inertness of the BN crucible seems to be maintained. A possible reaction pathway is given in Figure 7.1. The surface amide groups of a hypothetical  $A_xN_z$  are attacked by HF from NH<sub>4</sub>F decomposition resulting in NH<sub>3</sub> as a leaving group upon the formation of the A-F bond. In the next step, the amide group of a hypothetical  $B_yN_z$  attacks at the A atom, and HF acts as a leaving group thus NH<sub>4</sub>F is regained. After multiple steps, the crystallization enthalpy could act as a driving force of the overall formation of the nitride



 $A_x B_y N_z$ . This is in accordance with investigations of the etching processes of Si<sub>3</sub>N<sub>4</sub> in aqueous HF solutions.<sup>[11]</sup>

**Figure 7.1.** Possible reaction pathway of hypothetical nitrides  $A_xN_z$  and  $B_yN_z$  to  $A_xB_yN_z$  via the NH<sub>4</sub>F mediated HP/ HT synthesis.

The preparation of  $Sc_5P_{12}N_{23}O_3$ ,  $Ti_5P_{12}N_{24}O_2$ , and  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> shows the benefits of NH<sub>4</sub>F over solid-state metathesis. The interstitial *TM* nitrides were believed to be too unreactive as starting materials. Yet, NH<sub>4</sub>F as a mineralizer can activate these compounds. This comes without the drawback of the quantitative use of Li in the starting materials. The formation of the first ternary compound of Ti, P, and N, Ti<sup>IV</sup>P<sub>4</sub>N<sub>8</sub>, also shows that N is a sufficient oxidizing agent for the oxidization of transition metals at the reaction conditions realized by multianvil presses.

Based on the findings of this thesis it is likely that these results can also be projected to other rare earth and transition metal nitrides. This enables systematic access to open shell nitridophosphates incorporating Co, Ni, Fe, Sm, or Mn. If these elements are only used in dopant quantities this could lead to a possible application of nitridophosphates as dilute magnetic semiconductors.<sup>[12]</sup> With band gaps in the visible part of the spectrum, photocatalysis comes to mind as a viable research interest for *TM* nitridophosphates.<sup>[13-15]</sup> Expanding the compositional range even further, syntheses of nitridophosphates containing group 1 or 2 elements and *TM* or *RE* seem feasible from NH<sub>4</sub>F as a mineralizer. In analogy to compounds like La<sub>3-x</sub>Ca<sub>1.5x</sub>Si<sub>6</sub>N<sub>11</sub>:Eu<sup>2+</sup> ( $x \approx 0.77$ )<sup>[16]</sup> nitridophosphates similar in composition could lead to unprecedented host lattices and therefore desired emission properties.

It is reasonable to assume that the activation of refractory nitrides could also be extended to other nitrides, e. g. Be<sub>3</sub>N<sub>2</sub> AlN or Ge<sub>3</sub>N<sub>4</sub>. This would lead to a broad spectrum of novel mixed networks like nitridoalumophosphates with yet unknown structure types and potentially intriguing properties. Not limited to nitridophosphates, the activation of interstitial and refractory nitrides could also lead to the discovery of some proposed nitridic perovskites with interesting electronic properties.

#### 7.2 STEM in Materials Science

Modern aberration-corrected (scanning) transmission electron microscopy enabled the reliable results of crystal structure determinations presented in this thesis. For Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> the preparation of oriented cross-section samples for atomic-resolution STEM-EDX clarified the distribution of Bi and Sn in combination with the results from structure determination from synchrotron diffraction. Both experiments show disorder of Bi and Sn on all corresponding Wyckoff sites. This is especially remarkable given a lattice parameter of 282.64(6) Å which results in 12 Wyckoff sites occupied by Sn/Bi. As the detection efficiency of EDX detectors favors heavy elements and as reports on chemical mapping of Si and P with atomic resolution could not be found in literature, the application of this intriguing for nitridosilicatephosphates technique became and silicon imidonitridophosphates. Considering that both compounds were yielded as powders, the time efficiency of this approach was not limited by the preparation of cross-sections but highly dependent on preferred orientations encountered in CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>. Nonetheless, STEM-EDX mappings could be performed on a crystallite with dimensions of a few hundred nanometers.

The presence of superstructure reflections and their origin could be confirmed by atomicresolution STEM-HAADF and electron diffraction in the case of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>F. In addition, STEM- HAADF showed the presence of imperfect long-range ordering seen as Ti split position. This led to the acquisition of HTSCXRD data and subsequently the discovery of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>.

The results of this thesis showcase the feasibility, necessity, and information gained by electron microscopy for explorative material sciences. Positional and occupational order/ disorder could be readily elucidated by STEM-HAADF and/or STEM-EDX with atomic resolution.

A prospect here is the parallel acquisition of EELS and EDX spectra with atomic resolution which would expand the distinguishable elements to even lighter elements e. g. N, O, F. This will lead to yet unachieved confidence in atom-type assignment in structure determination of (oxo)nitrides.

The advent of iDPC (integrated differential phase contrast) enables the visualization of the projected electronic potential of the sample.<sup>[17]</sup> This eliminates the drawback of STEM-HAADF namely the dependency of the image contrast of  $\sim Z^2$  which leads to an underrepresentation of light elements in the image contrast. With iDPC, image contrast is nearly dependent on Z and therefore direct imaging of light elements like Li with simultaneous presence of heavy elements is feasible. In combination with another emerging field, *in situ* STEM, this could lead to unprecedented insights into the charging and discharging of Li-Ion batteries.

Still unachieved but long sought after is the combination of STEM-HAADF tomography with atomic resolution spectroscopy. This would result in a 3D model of the investigated nanocrystallite with information on the element type of each atom. In addition to 3D electron crystallography methods like ADT (automated electron diffraction tomography)<sup>[18]</sup> or RED (rotation electron diffraction)<sup>[19]</sup>, this could surpass structure elucidation from PXRD data in terms of accuracy and efficiency if the (S)TEM routines become user-friendly.

### 7.3 Final Remarks

A synthetic approach to novel nitrides was found within this thesis that delivered simple and effective access to mixed nitridic networks and transition metal nitride semiconductors. Crystallographic characterizations by atomic-resolution STEM and single-crystal synchrotron diffraction became only possible with techniques that were not available two decades ago.

The structural diversity achievable by mixed networks is undoubtedly immense and, in nature, illustrated by the diversification enabled by going from oxosilicates to oxoalumosilicates. The synthetic pathway presented in this thesis may certainly provide access to further possible host materials for activator ions containing  $SiN_x$  and  $PN_y$  polyhedra simultaneously. Since the synthetic part of this thesis tackled the issue of differing reactivities of nitrides, it is to be expected that the  $NH_4F$  mediated synthesis can be applied to other network-forming refractory nitrides such as  $Be_3N_2$  and AlN.

Although nitridophosphates gained interest as possible components for solid-state lighting, most of these compounds have band gaps at an insulator level. With the preparation of transition metal compounds, the possible range of applications grew by semiconductor applications like photocatalysis and photovoltaic as comparable band gaps to compounds like CdTe were achieved by combining the non-toxic elements Ti, P, and N. Experimental results on transition metal compounds may soon confirm theoretical chemists' claims on the stability and performance of predicted ternary nitrides.<sup>[20,21]</sup>

The electron microscopy presented emphasizes the need for advanced structural characterization in materials science already at the level of explorative and fundamental research. This will be of future interest and need in the emerging field of energy conversion.

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### 8 Summary

### 8.1 General Aspects

The results of this thesis illustrate the synergy of explorative solid-state chemistry and atomic resolution STEM in the preparation and structure elucidation of novel nitrides and other compounds. For instance, the structure elucidation of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> from microfocused synchrotron radiation was complemented by atomic-resolution STEM-HAADF and -EDX mappings. The synergistic approach revealed a complicated stacking sequence resulting in a lattice parameter of 282.64(6) Å and Sn/Bi occupational disorder on 12 Wyckoff sites.

Before this work, only two compounds featured  $SiN_x$  and  $PN_y$  polyhedra simultaneously:  $SiPN_3$  and  $SiP_2N_4NH$ .<sup>[1,2]</sup> By applying NH<sub>4</sub>F in high-pressure high-temperature (HP/HT) synthesis, the first nitridosilicatephosphates  $AESiP_3N_7$  (AE = Sr, Ba) were synthesized from the alkalineearth azides,  $Si_3N_4$  and  $P_3N_5$ .<sup>[3]</sup> The compounds proved suitable as host lattices for activator ions as shown by their luminescence spectra upon Eu<sup>2+</sup> doping. The occupational disorder of Si and P on one Wyckoff site in the simultaneous presence of one Wyckoff site solely occupied by P was investigated by the geometrical considerations of polyhedra volumes and STEM-EDX mapping with atomic resolution.

With the alkaline-earth azides exchanged for the alkaline-earth amides, the new approach resulted in the first nitridic mica-like compounds  $AESi_3P_4N_{10}(NH)_2$  ( $AE = Mg, Mg_{0.94}Ca_{0.06}, Ca, Sr$ ).<sup>[4]</sup> These compounds feature the high-pressure motif of sixfold coordinated Si beside PN<sub>4</sub> tetrahedra as shown by STEM-EDX mappings with atomic resolution in combination with microfocused synchrotron SCXRD experiments. Therefore, the NH<sub>4</sub>F-assisted approach also proved successful in synthesizing imide group-containing compounds.

The preparation of transition metal-containing nitrides with NH<sub>4</sub>F as a mineralizer constitutes another central aspect of this thesis. The oxonitridophosphates  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$  were synthesized from the refractory nitrides ScN or TiN, and  $P_3N_5$  upon the addition of PON or TiO<sub>2</sub>, respectively, as a solid oxygen source.<sup>[5]</sup> The mixed-anion compounds also show mixed-valence for  $Ti_5P_{12}N_{24}O_2$  with the ordering of  $Ti^{3+}$  and  $Ti^{4+}$  ions.

The polymorphs of TiP<sub>4</sub>N<sub>8</sub> showcase the analytical demand ranging from *ex situ* SCXRD, *in situ* HTSCXRD, HPXRD, atomic resolution STEM-HAADF to DFT calculations following the simple reaction of TiN with P<sub>3</sub>N<sub>5</sub>, NH<sub>4</sub>N<sub>3</sub>, and NH<sub>4</sub>F. Additionally, the NH<sub>4</sub>F mediated HP syntheses

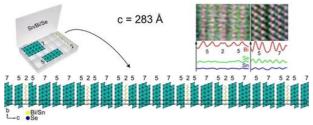
presented themselves as a reasonable approach to new ternary transition metal nitride semiconductors.

# 8.2 A Layered Tin Bismuth Selenide with Three Different Building Blocks that Account for an Extremely Large Lattice Parameter of 283 Å

published in:M. Nentwig, L. Eisenburger, F. Heinke, D. Souchay, O. OecklerChem. Eur. J. 2020, 26, 10676.

**DOI:** 10.1002/chem.202000663

The crystal structure of 150*R*-Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> was elucidated by microfocused synchrotron diffraction. Crystal structure refinements yielded an unusually large unit cell in space group

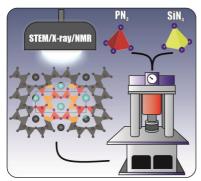


 $R\overline{3}m$  witch a = 4.1819(4) Å and c = 282.64(6) Å. All 12 Sn/Bi Wyckoff sites are occupationally disordered according to SCXRD. A cross-section sample from bulk Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> was analyzed by atomic-resolution STEM-HAADF and STEM-EDX mapping.

 8.3 High-Pressure High-Temperature Synthesis of Mixed Nitridosilicatephosphates and Luminescence of AESiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> (AE = Sr, Ba) published in: L. Eisenburger, O. Oeckler, W. Schnick

*Chem. Eur. J.* **2021**, *27*, 4461. **DOI:** 10.1002/chem.202005495

The first nitridosilicatephosphates were synthesized by highpressure high-temperature synthesis employing a Walkertype multianvil press and NH<sub>4</sub>F as a mineralizer at 8 GPa and 1400–1700 °C. The crystal structures of SrSiP<sub>3</sub>N<sub>7</sub> and BaSiP<sub>3</sub>N<sub>7</sub> were elucidated by single-crystal X-ray diffraction and atomic resolution STEM-EDX mappings. This revealed a disordered Si<sub>0.5</sub>P<sub>0.5</sub> and a fully ordered P site. Additional characterization

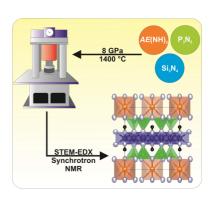


was performed by FTIR, SEM, MAS-NMR, and luminescence spectroscopy on  $\mathrm{Eu}^{2+}$  doped samples.

## 8.4 Nitridic Analogs of Micas $AESi_3P_4N_{10}(NH)_2$ (AE = Mg, Mg<sub>0.94</sub>Ca<sub>0.06</sub>, Ca, Sr)

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#### **DOI:** 10.1002/anie.202114902



High-pressure high-temperature synthesis employing NH<sub>4</sub>F led to the first nitridic analogs of micas. The compounds  $AESi_3P_4N_{10}(NH)_2$  ( $AE = Mg, Mg_{0.94}Ca_{0.06}$ , Ca, Sr) crystallize isotypic to the mineral clintonite. The crystal structure was elucidated on microcrystalline samples of CaSi\_3P\_4N\_{20}(NH)\_2 from microfocused synchrotron radiation experiments and STEM-EDX mappings with atomic resolution revealing PN<sub>4</sub>

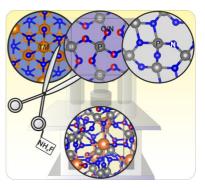
tetrahedra and SiN<sub>4</sub>(NH)<sub>2</sub> octahedra layers. The presence and location of the imide group, -NH, were furthermore confirmed by FTIR and MAS-NMR spectroscopy. Eu2+ doped samples showed narrow-band emission with maxima ranging from 451 to 492 nm.

## 8.5 High-Pressure Synthesis of Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub> and Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub> by Activation of Binary Nitrides ScN and TiN with NH<sub>4</sub>F

published in: L. Eisenburger, V. Weippert, O. Oeckler, W. Schnick,

Chem. Eur. J. **2021**, 27, 14184.

**DOI:** 10.1002/chem.202101858



The compounds  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$  were synthesized by the NH<sub>4</sub>F mediated high-pressure hightemperature approach from the binary nitrides and PON or  $TiO_2$ , respectively, acting as an O source. Both compounds are isotypic to  $Ti_5B_{12}O_{26}$  as elucidated by single-crystal X-ray diffraction. Magnetic measurements revealed the ordering of  $Ti^{3+}$  and  $Ti^{4+}$  ions on Wyckoff positions 32g and 8b respectively. The optical band gaps were determined as direct band gaps of 1.6 eV for  $Ti_5P_{12}N_{24}O_2$  and 3.8 eV for  $Sc_5P_{12}N_{23}O_3$ .

## 8.6 Discovery of Two Polymorphs of TiP<sub>4</sub>N<sub>8</sub> Synthesized from Binary Nitrides

published in:L. Eisenburger, Valentin Weippert, Carsten Paulmann, Dirk Johrendt,Oliver Oeckler, W. Schnick, Angew. Chem. Int. Ed. 2022, DOI 10.1002/anie.202202014.Angew. Chemie 2022, DOI 10.1002/ange.202202014.

The first Ti<sup>IV</sup> nitridophosphate was synthesized via the NH<sub>4</sub>Fassisted high-pressure high-temperature route from the binary nitrides. The combination of STEM-HAADF, in situ single-crystal X-ray diffraction, and annealing of samples, revealed the crystal structures of two polymorphs. While the crystal structure of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> strongly resembles that of orthorhombic svyatoslavite (CaAl<sub>2</sub>Si<sub>2</sub>O<sub>8</sub>) the crystal structure

a-TiP<sub>4</sub>N<sub>8</sub> 8 GPa 1400 °C β-TiP<sub>4</sub>N<sub>8</sub> γ-tiP<sub>4</sub>N<sub>8</sub> γ-tiP<sub>4</sub>N<sub>8</sub>

**DOI:** 10.1002/anie.202202014.

of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is a threefold superstructure along *h* thereof.  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> feature indirect band gaps in the visible part of the spectrum as shown by DFT calculations of 1.78 and 1.5 eV, respectively. DFT calculations also revealed that  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> is a metastable polymorph and the observed phase transition is monotropic.

Appendix

## A Supporting Information for Chapter 2

### **Experimental Procedures**

**Synthesis.**  $(SnSe)_{0.5}Bi_2Se_3 = SnBi_4Se_7$  was synthesized from a stoichiometric mixture of the elements Sn (99.999%, Koch chemicals or Aluterv-FKI), Bi (99.999%, Aldrich or Alfa Aesar), and Se (99.999%, ChemPur). The mixture (total weight ~3 g) was fused in a sealed silica glass ampule (length ~10cm, diameter 15 mm) under dry Ar atmosphere and melted at 900 °C for two days, followed by quenching at air. After annealing the sample at 500 °C for one week, it was quenched in water. The resulting product with a nominal composition of Sn<sub>0.571</sub>Bi<sub>2.286</sub>Se<sub>4</sub> was heated three times up to 500 °C. Theheterogeneous product was examined by scanning electron microscopy (SEM) and energy-dispersive X-ray spectroscopy (EDX) as described below, which revealed regions of different contrast and composition (Fig. S1 and Table S1). Next to a minority phase, which can be identified as Sn<sub>11.49</sub>Bi<sub>12.39</sub>Se<sub>30</sub> (phase 1)<sup>[1]</sup> by EDX, and one main phase of Sn<sub>0.85</sub>Bi<sub>2.15</sub>Se<sub>4</sub> (phase 2),<sup>[1]</sup> there is a slighty Sn-depleted (compared to the nominal composition SnBi<sub>4</sub>Se <sup>1</sup> phase (phase 3), which constitutes the title compound of this contribution.

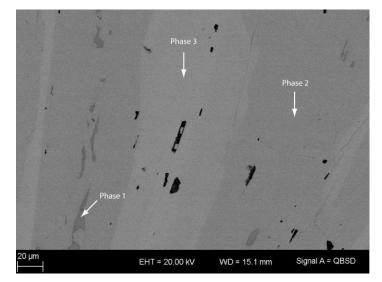
**Conventional transmission electron microscopy.** Selected-area electron diffraction (SAED), high resolution transmission electron microscopy (HRTEM) and further EDX measurements (Table S2) were executed on a Philips CM-200 STEM transmission electron microscope (LaB<sub>6</sub> cathode, 200 kV, super-twin lens, point resolution 0.23 nm) equipped with RTEM 136-5 EDX detector (EDAX, Genesis software).<sup>[2]</sup> A double-tilt low-background sample holder (Gatan) was used. For these TEM measurements, i.e. the selection of crystallites for synchrotron measurements, the powdered sample was suspended in ethanol and deposited on a copper grid coated with a holey carbon film. Evaluation of SAED data was done using the *analySIS* software.<sup>[3]</sup> For HRTEM and SAED simulations, jEMS was used.<sup>[4]</sup> Micrographs and diffraction patterns were evaluated using Digital Micrograph.<sup>[5]</sup>

**Sample preparation for scanning transmission electron microscopy.** Polycrystalline pieces of the sample were embedded in Gatan G1 epoxy resin (Gatan, USA) between two

silicon wafers, which were subsequently glued into a brass tube (inner diameter 2.4 mm, wall thickness 0.3 mm) with the same epoxy resin. Discs of ~500  $\mu$ m thickness were cut from this assembly using a diamond wire saw (well GmbH, Germany). The discs were mechanically thinned to a thickness of ca. 80  $\mu$ m with a disc grinder (Gatan, USA) using SiC abrasive paper with particle sizes of 40  $\mu$ m and subsequently 10  $\mu$ m (Gatan, USA). The sample was further thinned from both sides to a thickness of ~20  $\mu$ m in the center using a dimple grinder (Gatan 656 dimple grinder, Gatan, USA) and diamond paste (5  $\mu$ m and 1  $\mu$ m, Kemet International LTD, Kent, UK). Electron transparency was achieved by Ar-ion milling using a precision ion polishing system (PIPS, Gatan, USA) at incidence angles of 6° and a voltage of 4 keV. After a hole in the center of the disc was visible, the sample was polished at 0.9, 0.6 and 0.3 keV for five minutes each. An image of the prepared sample is shown in Figure A13.

Scanning transmission electron microscopy. The sample was mounted on a double-tilt holder and transferred into a C<sub>S</sub> DCOR probe-corrected Titan Themis 300 (FEI, USA) instrument equipped with X-FEG, post-column filter (Enfinium ER-799, Gatan, USA), US1000XP/FT camera system (Gatan, Germany) and a windowless, 4-quadrant Super-X EDX detector. TEM images were recorded using a 4k × 4k FEI Ceta CMOS camera. The microscope was operated at 300 kV or 200 kV accelerating voltage for SAED and STEM-HAADF (16.6 mrad convergence angle, 50 µm aperture, 63 mrad detector inner half angle for 100 mm camera length). TEM data were evaluated with Digital Micrograph <sup>[5]</sup> for image analysis, ProcessDiffraction7<sup>[6]</sup> for geometric calculations concerning SAED patterns, JEMS<sup>[4]</sup> for SAED simulations and the software suite Velox v2.8 (Thermo Fisher Scientific, USA) for recording and evaluating EDX maps. These were acquired over 120 x 138 pixels with 20 µs dwell time each. The resulting maps consist of 953 (300 keV) / 140 (200 keV) individual frames. Drift corrections were applied. EDX signals were integrated as a projection on a line of 135 pixels (2.441 nm) in length, using an integration width of 100 pixels for 300 keV (Fig. 3c). For 200 keV, they were projected on a line of 285 pixels (2.297 nm) in length, using an integration width of 140 pixels (Fig. 3f). The element maps were filtered using a radial Wiener filter (highest frequency 20.0 software-specific units and edge smoothing 30.0 units for 300 keV; highest frequency 13.0 units and edge smoothing 30.0 units for 200 keV).

Single-crystal X-ray diffraction. Single-crystal data of a crystallite on a copper finder grid (Fig. S2), which was fixed on a glass capillary, were collected at beamline ID11 of the European Synchrotron Radiation Facility (ESRF, Grenoble) using microfocused synchrotron radiation. The crystal was centered in a beam of ca.  $1.5 \times 3 \mu m^2$  with the help of fluorescence scans using the Bi- $K_{\alpha}$  emission lines, which were recorded and processed using a SiriusD Silicon Drift Detector with a XIA Mercury pulse processor (Ketek, Germany). Diffraction data were recorded with a Frelon4k CCD detector.<sup>[7]</sup> Integration and semiempirical absorption correction were performed with CrysalisPro.<sup>[8]</sup> A correction for incomplete absorption of high- energy radiation in the phosphor of the CCD detector was applied.<sup>[9]</sup> Structure solution and refinement as well as the calculation of Fourier maps and simulation of diffraction patterns were performed with JANA2006 (version 08/10/2018).<sup>[10]</sup> For the tentative refinement of the structure model in 3D, the program SHELX <sup>[11]</sup> was used. Crystal structures were visualized with Diamond.<sup>[12]</sup> Further details of the crystal structure investigations may be obtained from the Fachinformationszentrum Karlsruhe, 76344 Eggenstein-Leopoldshafen, Germany (fax: +49-7247-808-666; e-mail: crys-data@fiz- karlsruhe.de), on quoting the depository number 1949450.



Sample characterization by SEM, TEM and EDX

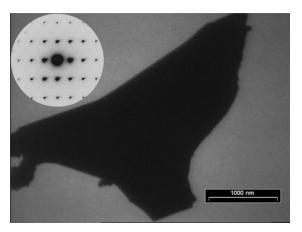
**Figure A1.** SEM-BSE image of a thermally treated sample (3 times cycled between RT and 500 °C) with the nominal composition SnBi<sub>4</sub>Se<sub>7</sub>. The image shows the presence of two majority phases (phase 2 and phase 3) and one minority phase (phase 1) after thermal treatment during the thermoelectric measurements. The compositions of the different phases are displayed in Table S1. Phase 2 (Sn<sub>0.85</sub>Bi<sub>2.15</sub>Se<sub>4</sub>) has been investigated in detail by F. Heinke et al.<sup>[1]</sup>

**Table A1.** SEM-EDX analyses (average of six point measurements for phase 1 and of four point measurements each for phases 2 and 3) for the sample with the nominal composition SnBi<sub>4</sub>Se<sub>7</sub> after thermal treatment (3 times cycled up to 500 °C).<sup>[1]</sup>

Element	Nominal	Phase 1	Calculated for	Phase 2	Calculated for	Phase 3	Calculated for
	compositior	n (minority	Sn11.49Bi12.39Se3(	) (majority	Sn0.85Bi2.15Se	<sup>4</sup> (majority	Sn2.8Bi20.2Se27
	SnBi <sub>4</sub> Se <sub>7</sub>	phase)	ref. [1]	phase)	ref. [1]	phase)	(this work)
Sn / at- %	8.3%	21.9(6)	22.0%	9.6(5)	12.2%	5.6(5)	5.6%
Bi / at- %	33.3%	24.3(4)	22.3%	35.9(3)	30.7%	39.9(4)	40.4%
Se / at- %	58.3%	53.8(9)	55.7%	54.5(10)	57.1%	54.6(10)	54.0 %

Table A2. EDX analysis (average of three point measurements) of the single crystal used for
synchrotron data collection (Fig. S2). Note that these EDX results may be rather imprecise as the
position of the crystal on the TEM grid was not ideal for recording EDX spectra.

Element	Single crystal	Calculated for	Calculated for	
		SnBi <sub>4</sub> Se <sub>7</sub>	Sn2.8(4)Bi20.2(4)Se27	
Sn / at-%	6.9(6)	8.3%	5.6%	
Bi / at-%	35.1(5)	33.3%	40.4%	
Se / at-%	58.0(3)	58.3%	54%	



**Figure A2.** Transmission electron micrograph (Philips CM-200) and a SAED pattern of the crystallite used for synchrotron data collection.

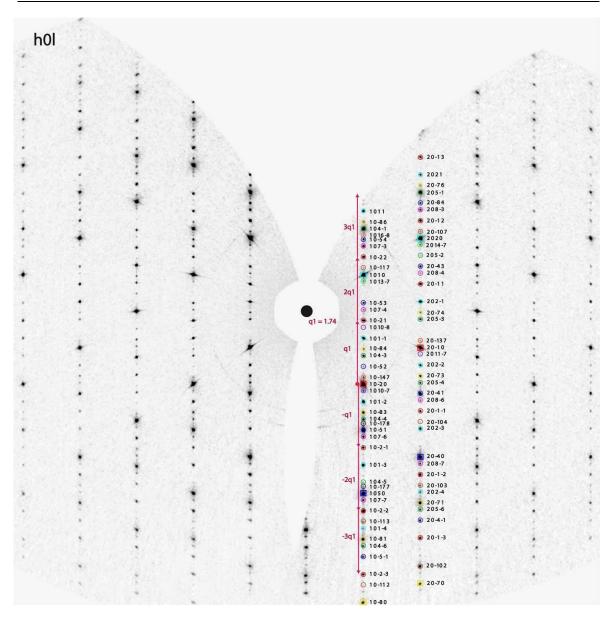
# **Results and Discussion**

**General remarks.** An initial attempt to determine the crystal structure solution from the single-crystal data by "routine data processing" resulted in the space group  $P\bar{3}m1$  with a =4.0955(2) Å c = 22.172(2) Å, which is close to the metrics of BiSe.<sup>[13]</sup> However, ca. 60% of all reflections remained unindexed. Close inspection of the diffraction pattern (cf. following figures) suggested that the diffraction pattern is best described using a (3+1)D superspace approach. A small trigonal unit cell with a = 4.1819(4) Å and c = 5.6528(4) describes the strong main reflections. The other reflections are indexed with a *q*-vector along [001]\* with *q* = 1.74 = 87/50, taking into account satellite reflections up to the 8<sup>th</sup> order. This (3+1)D commensurate modulation explains all reflections on the reciprocal lattice sections (Fig. S3 -S5). Due to the presence of no systematic absences for the satellite reflections, the superspace group  $R\bar{3}m$  (00y)00 was chosen instead of the other possibility - i.e.  $R\bar{3}m(00\gamma)0s$  corresponding to the 3D space group  $R^{\overline{3}}m$ . Structure solution by charge flipping <sup>[14]</sup> confirmed the assumed superspace group  $R\bar{3}m(00\gamma)00$  with a single atom site (0 0 0). This (3+1)- dimensional model with a small trigonal unit cell and a *q*-vector along  $[001]^*$  is consistent with the superspace formalism developed for layered bismuth selenides.<sup>[15]</sup> This straightforward description unifies all these compounds in a single structure model whose modulation vector actually determines which variant is present. This approach can be used for both commensurate and incommensurately modulated structures.

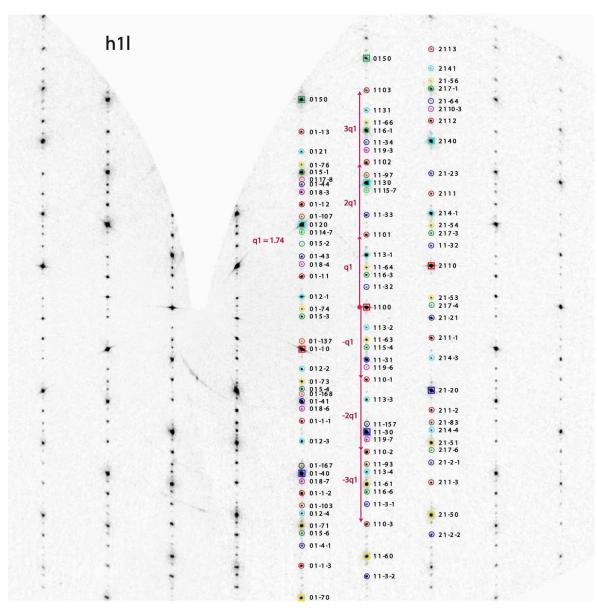
## **Diffraction patterns**

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**Figure A3.** Reciprocal lattice section 0kl with the *q*-vector of 1.74 (red arrows) and identified reflections up to  $8^{th}$  order "satellites". The satellite reflections are highlighted in the same color as the corresponding main reflections.

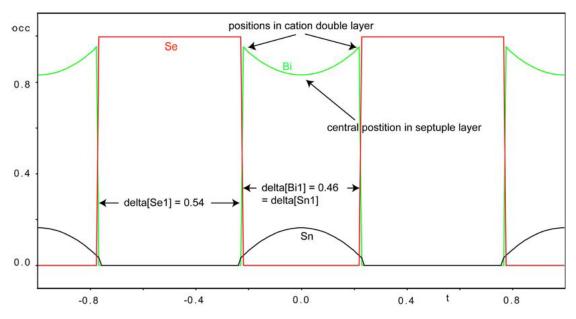


**Figure A4.** Reciprocal lattice section h0l with the *q*-vector of 1.74 (red arrows) and identified reflections up to 8<sup>th</sup> order "satellites". The satellite reflections are highlighted in the same color as the corresponding main reflections.

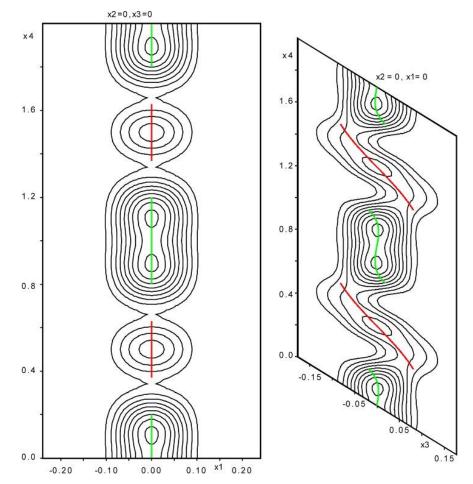


**Figure A5.** Reciprocal lattice section h1l with the *q*-vector of 1.74 (red arrows) and identified reflections up to 8<sup>th</sup> order "satellites". The satellite reflections are highlighted in the same color as the corresponding main reflections.

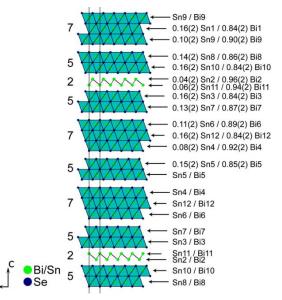
Refinement details. As a starting model, one Bi (Bi1) and Se atom (Se1) were set on the single atom site at the origin (0,0,0), both with formal full occupancy (ai[Bi1] = ai[Se1] =0.08333). Note that in contrast to 3D structure refinements, the additional dimension in (3+1)D descriptions can be used to modulate the occupancy using discontinuous functions, e.g. crenel- or sawtooth-shaped.<sup>[14]</sup> Crenel functions are specified by the width (delta) and the centre (x40) and describe the modulation of occupancy by defining an interval in the additional dimension x4 where the corresponding atom exists. In the refinement for Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub>, crenel functions combined with Legendre polynomials were used for both atoms. The crenel function of Bi1 has its center at x40[Bi1] = 0 and a refined width of delta[Bi1] = 0.4602(12); the function of Se1 atom has its center at x40[Se1] = 0.5 and a refined width of delta[Se1] = 0.5398(2). The sum of the overall width was constrained to 1 to retain full occupancy in x4. As the resulting structure model corresponds to a commensurate modulation, there is a corresponding 3D structure model in space group  $R^{\overline{3}}m$  with lattice parameters a = 4.1819(4) Å and c = 282.64(6) Å. It contains 150 atom layers forming double ("2"), quintuple ("5") and septuple ("7") slabs, which results in a stacking sequence symbolized by 7525757525|7525757525|7525757525. After the refinement of anisotropic atom displacement parameters, additional positional modulation waves for each atom were included. Since the x1 (= x) and x2 (= y) contribution of the q-vector is zero, the modulation functions are restricted to the x3 (z) direction. The positional modulation waves for Se (red) and Bi/Sn (green) thus only show displacements along the x3 direction, which is displayed in the  $F_{obs}$  Fourier maps (Fig. S6). As a next step, *hklm* reflections of m = 5-8 were excluded from the refinement due to their high R1 values compared to the other satellite orders; m =1-3: R1(all) = 0.09-0.11, m = 4: R1(all) = 0.28, m = 5-8: R1(all) = 0.41-1.31. At this stage, the width of the Bi1 wave function refined to delta[Bi1] = 0.4608(10). This value was fixed at delta[Bi1] = 0.46 in order to provide integer numbers of atoms in the unit cell of the commensurate structure. Note that STEM micrographs clearly show the commensurate structure. Therefore, the *q*-vector of q = 1.74 = 87 / 50 can also be written as a fraction of an integer number (87) and the number of atoms in one stacking period 7525757525 (50) (cf. Tab. S9). Sn (Sn1) was introduced into the structure with identical coordinates, modulations and displacement parameters as Bi1, including delta[Sn1] = delta[Bi1]. The occupancies were modulated in a complementary way with site occupancy factors ai[Sn1] = 0.083333 - ai[Bi1], i.e. constrained to full total occupancy. The refined value of ai[Sn1] = 0.014(2) corresponds to a seemingly uniform distribution of 17% Sn and 83% Bi on all cation sites. An additional occupancy wave for Sn1 and Bi1 reveals a varying distribution of 4% (Bi-Bi double layer) up to 16% Sn (central position in septuple layer) on the cation sites. Note that this seemingly yields a chemical composition that is slightly different from the model that assumed uniform distribution, i.e. the uniform distribution is not an "average" of the modulated one; this difference, however, is not significant within standard deviations. Fig. S7 shows the occupancy of the atoms Se (red), Bi (green) and Sn (black) as a function of the t variable in superspace. This graph displays where the corresponding atoms exist in the additional x4 dimension. Note that this occupancy wave can be (rather incorrectly) modelled assuming a modulation of the displacement parameters for a uniform distribution of Sn and Bi on the cation sites. Such a modulation of the displacement parameters is not significant after applying the occupational wave. This corroborates a varying cation distribution, and no modulation of the displacement parameters was applied in the final model. The structure refinement results are displayed in Table S4. Information on the modulation parameters are given in Table S5 and S6.



**Figure A6.** *F*<sub>obs</sub> Fourier maps with the positional modulations waves for Se (red) and Bi/Sn (green), shown along x1 (= x; left) and x3 (= z; right) direction.



**Figure A7.** Graph showing the occupancy of the atoms Se (red), Bi (green) and Sn (black) as a function of the *t* variable in superspace. The Bi/Sn atom has its domain centered at t = 0 and the Se atom centered at t = 0.5.



#### Description of the structure model

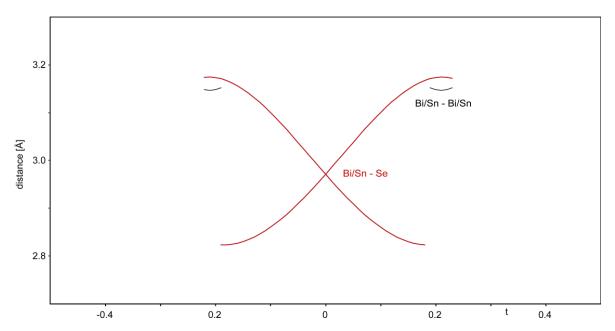
**Figure A8.** Characteristic part of the crystal structure of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> with the stacking sequence 7525757525. The mixed occupancies of Sn and Bi are displayed for the different cation positions listed in Table S8. Site occupancies are given for the asymmetric unit, labelled atoms without s.o.f. values are symmetry-equivalent to the others.

**Table A3.** Bond lengths (in Å) derived from crystal structure analysis of Sn2.8(4)Bi20.2(4)Se27 comparedwith (Sn)/Bi/Se compounds from literature.

Compound	Bi-Bi	Bi-Se (next to van der	Bi-Se	Sn-Se	Se-Se (at van der	Bi <sub>2</sub> -Se
(and references)	DI-DI	Waals gap)	(center) (center)		Waals gap)	[a]
Bi <sub>2</sub> Se <sub>3</sub> <sup>[16]</sup>	-	2.85	3.08	-	3.52	-
BiSe [13]	3.05	2.91	3.03	-	3.51	3.27
Bi <sub>8</sub> Se <sub>9</sub> <sup>[17]</sup>	3.06	2.94	2.95 - 3.07	-	3.45	3.26
Bi <sub>8</sub> Se <sub>7</sub> [16] [b]	2.68 -	2.98 - 3.13	2.98 - 3.15	_	3.05	3.21 -
D18367 1 1111	3.10	2.70 - 5.15	2.70 - 5.15		5.05	3.68
Bi <sub>4</sub> Se <sub>3</sub> <sup>[18]</sup>	3.04	2.88	3.08	-	-	3.47
SnBi <sub>2</sub> Se <sub>4</sub> [19]	-	2.80	3.17	2.95	3.72	-
$Sn_{0.57}Bi_{2.29}Se_4$ [20]	-	2.86	3.12	2.95	3.44	-
Sn <sub>0.57</sub> Bi <sub>2.29</sub> Se <sub>4</sub> [12]	-	2.85	3.11	2.96	3.48	-
Sn <sub>2.8(4)</sub> Bi <sub>20.2(4)</sub> Se <sub>27</sub>	3.15	2.83 - 2.88	3.02 - 3.17	2.95 - 3.00	3.57	3.17

[a] bond length of Se atom to Bi atom (part of cation double layer)

[b] rather unreliable distances from imprecise coordinates derived from X-ray photographic film data, taking into account only 00ℓ reflections.



**Figure A9.** Graph of interatomic distances vs. the *t* variable of superspace. The Bi/Sn – Bi/Sn bonds in the double layers (corresponding to the structure of elemental Bi, cf. main manuscript) are displayed in black, the Bi/Sn – Se bonds within the quintuple and septuple layers are displayed in red.

# Numerical refinement results

Table A4. Crystallographic data of the single-crystal structure refinement of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> in

Sum formula for basic unit cell	Sn0.055(7)Bi0.405(7)Se0.54			
Formula mass / g • mol-1	133.8			
Superspace group	$R\bar{3}m(00\gamma)00$			
Basic cell parameters / Å	a = 4.1819(4), c = 5.6528(5)			
c <sub>supercell</sub> (of commensurate 3D structure) / Å	282.6(6)			
<i>q</i>	1.74 = 87/50			
X-ray density / g ·cm⁻³	7.786			
Temperature / K	295			
Radiation	synchrotron, $\lambda = 0.29470$ Å			
heta range / °	2.32 - 17.42			
μ / mm <sup>-1</sup>	8.703			
Reflections meas. / indep.	6375 / 411			
hklm range	$-6 \le h \le 6$			
	$-8 \le k \le 7$			
	$-11 \le l \le 11$			
	$-3 \le m \le 3$			
Extinction method	type 1 Gaussian isotropic (JANA)			
Extinction coefficient	530 (160)			
Parameters	12			
$R_{ m int}; R_{\sigma}$	0.0310; 0.0054			
$R1 [I>2\sigma(I)] / R1 (all)$	0.446 / 0.451			
$wR$ [I>2 $\sigma$ (I)] / $wR$ (all)	0.0678 / 0.0699			
main reflections: $R1$ [I>2 $\sigma$ (I)]	0.0481 / 0.0600			
1 <sup>st</sup> order satellite reflections: <i>R</i> 1 [I>2 $\sigma$ (I)]	0.0384 / 0.494			
$2^{nd}$ order satellite: <i>R</i> 1 [I>2 $\sigma$ (I)]	0.0600 / 0.1086			
$3^{rd}$ order satellite: <i>R</i> 1 [I>2 $\sigma$ (I)]	0.0347 / 0.0504			
Goof	1.62			
$\Delta \rho$ (min./ max., eÅ <sup>-3</sup> )	-2.56 / 6.10			

Atoms x y z	site occupancy	delta	Resulting	Occupational	Positional	
	x y Z	factor (ai)	uena	occupancy	modulation	modulation
Sn1 /	0 0 0	0.0733(13) /	0.46	0.055(7)/	2 Legendre	2 Legendre
Bi1	Bi1	0.0100(13)	0.40	0.405(7)	polynomials	polynomials
Se1	0 0 0	0.083333	0.54	0.54	1 Legendre	2 Legendre
Se1 0	000	0.0055555	0.34	0.54	polynomial	polynomials

**Table A5.** Atom sites, site occupancy factors (relative to general position),, resulting occupancies andmodulation functions (details cf. Tab. S6) of the structure refinement in superspace.

**Table A6.** Anisotropic displacement parameters (in Å<sup>2</sup>) and modulation parameters of the structure refinement of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> in (3+1)D superspace; $U_{eq} = 1/3[U_{11} + U_{22} + U_{33}]$ .

Atoms	$U_{11} = U_{22} =$	<i>U</i> <sub>33</sub>	U <sub>eq</sub>	oort1	oort2	zort1	zort2	zort3	zort4
	2 <i>U</i> <sub>12</sub>								
Bi1	0.0157(3)	0.0354(4)	0.0223(2)	0	0.0997(13)	-0.0051(2)	0	-0.0136(4)	0
Sn1	0.0157(3)	0.0354(4)	0.0223(2)	0	-0.7293(13)	-0.0051(2)	0	-0.0136(4)	0
Se1	0.0130(4)	0.0285(6)	0.0182(3)	-	-	-0.0855(5)	0	0.0046(7)	0

**3D** structure model. Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> can be described as a long-periodic 3D structure with a stacking sequence of 7525757525|7525757525|7525757525|7525757525 built up from double (2), quintuple (5) and septuple (7) layers. The 3D structure can be derived from the commensurately modulated one. A tentative refinement in 3D has been performed by transforming the  $hk\ell m$  file to a standard SHELX 3D  $hk\ell$  file by calculating  $\ell_{3D}$  values:  $\ell_{3D} = m \cdot (c_{supercell} / c_{basic cell}) + \ell_{4D} \cdot (c_{supercell} / c_{basic cell}) \cdot q$  with  $c_{supercell} = c_{3D} = 50 \cdot c_{basic cell}$ , i.e. the lattice parameter of the 3D model. The structure model derived from the commensurately modulated as starting model. Site occupancies were fixed and the atom coordinates as well as isotropic displacement parameters were refined. The latter were set equal for all cations and all anions, respectively.

The sum formula  $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27}$  corresponds to 23 cations and 27 anions of a stacking sequence of 7525757525. Assuming  $Bi^0$  and  $Sn^0$  in the

double layers (comparable to those in elemental Bi) and Sn<sup>2+</sup>, Bi<sup>3+</sup> and Se<sup>2-</sup> in the NaCl-typelike quintuple and septuple layers, the overall charge of the compound is neutral within standard deviations:

 $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27} = Bi_{3.8}Sn_{0.2}$  (two double layers) +  $Sn_{2.6(4)}Bi_{16.4(4)}Se_{27}$  (five quintuple and three

# septuple layers):

 $\rightarrow$  (3.8 • 0 + 0.2 • 0) + (2.6 • 2 + 16.4 • 3 - 27 • 2) = 0.4 negative charges.

Sum formula	Sn <sub>2.8(4)</sub> Bi <sub>20.2(4)</sub> Se <sub>27</sub>				
Formula mass / g • mol <sup>-1</sup>	6688.66				
Space group	$R\overline{3}m$				
Cell parameters / Å	a = 4.1819(4), c = 282.64(6)				
Cell volume / Å <sup>3</sup>	4280.7(12)				
X-ray density / g ⋅cm <sup>-3</sup>	7.784				
Formula units per cell	3				
F(000)	8207				
Radiation	synchrotron, $\lambda = 0.29470$ Å				
heta range / °	3.32 - 14.18				
Reflections meas. / indep.	896 / 131				
hklm range	$-6 \le h \le 6$				
	$-6 \le k \le 6$				
	$-437 \le l \le 400$				
Parameters	28				
$R_{ m int}$ ; $R_{\sigma}$	0.0251; 0.0081				
<i>R</i> 1 [I>2 $\sigma$ (I)] / <i>R</i> 1 (all)	0.0410 / 0.0411				
$wR$ [I>2 $\sigma$ (I)] / $wR$ (all)	0.1018 / 0.1018				
Goof	1.248				
Δρ (min./ max., eÅ-3)	-2.06/ 1.40				

Table A7. Crystallographic data of the structure refinement of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> in 3D.Sum formula

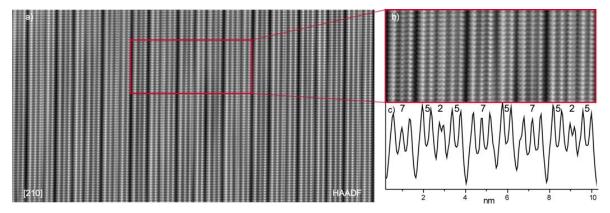
**Table A8.** Wyckoff positions, atomic coordinates, site occupancies and isotropic displacement parameters (in Å<sup>2</sup>) for the commensurable 3D structure model of  $Sn_{2.8(4)}Bi_{20.2(4)}Se_{27}$ . The *z* positions and U<sub>iso</sub> have been refined for each atom. Site occupancy factors were taken from the refinement in (3+1)D and not refined. The z coordinates of the starting model as derived from the refinement in superspace are given as a comparison.

Atom	Wyckoff position	X	у	Ζ	z derived from (3+1)D model	s.o.f.	Uiso
Sn1/Bi1	3 <i>a</i>	0	0	0	0	0.16/0.84	0.0156(4)
Sn2/Bi2	3 <i>a</i>	0	0	0.0609(1)	0.0597	0.04/0.96	0.0156(4)
Sn3/Bi3	6 <i>c</i>	0	0	0.0808(1)	0.0800	0.016/0.84	0.0156(4)
Sn4/Bi4	6 <i>c</i>	0	0	0.1396(5)	0.1399	0.08/0.92	0.0156(4)
Sn5/Bi5	6 <i>c</i>	0	0	0.1593(2)	0.1599	0.15/0.85	0.0156(4)
Sn6/Bi6	6 <i>c</i>	0	0	0.2196(3)	0.2200	0.11/0.89	0.0156(4)
Sn7/Bi7	6 <i>c</i>	0	0	0.2395(3)	0.2399	0.13/0.87	0.0156(4)
Sn8/Bi8	6 <i>c</i>	0	0	0.3004(3)	0.3001	0.14/0.86	0.0156(4)
Sn9/Bi9	6 <i>c</i>	0	0	0.3204(4)	0.3200	0.10/0.90	0.0156(4)
Sn10/Bi10	6 <i>c</i>	0	0	0.3801(4)	0.3801	0.16/0.84	0.0156(4)
Sn11/Bi11	6 <i>c</i>	0	0	0.4005(4)	0.4002	0.06/0.94	0.0156(4)
Sn12/Bi12	6 <i>c</i>	0	0	0.4600(9)	0.4600	0.16/0.84	0.0156(4)
Se1	6 <i>c</i>	0	0	0.0197(8)	0.0185	1	0.0124(6)
Se2	6 <i>c</i>	0	0	0.0401(15)	0.0401	1	0.0124(6)
Se3	6 <i>c</i>	0	0	0.0986(7)	0.0987	1	0.0124(6)
Se4	6 <i>c</i>	0	0	0.1203(7)	0.1204	1	0.0124(6)
Se5	6 <i>c</i>	0	0	0.1792(9)	0.1790	1	0.0124(6)
Se6	6 <i>c</i>	0	0	0.2012(7)	0.2007	1	0.0124(6)
Se7	6 <i>c</i>	0	0	0.2588(10)	0.2592	1	0.0124(6)
Se8	6 <i>c</i>	0	0	0.2812(9)	0.2809	1	0.0124(6)
Se9	6 <i>c</i>	0	0	0.3406(1)	0.3395	1	0.0124(6)
Se10	6 <i>c</i>	0	0	0.3613(7)	0.3612	1	0.0124(6)
Se11	6 <i>c</i>	0	0	0.4198(8)	0.4197	1	0.0124(6)
Se12	6 <i>c</i>	0	0	0.4405(10)	0.4414	1	0.0124(6)
Se13	6 <i>c</i>	0	0	0.4778(2)	0.4784	1	0.0124(6)
Se14	6 <i>c</i>	0	0	1/2	1/2	1	0.0124(6)

#### Comparison with other compounds

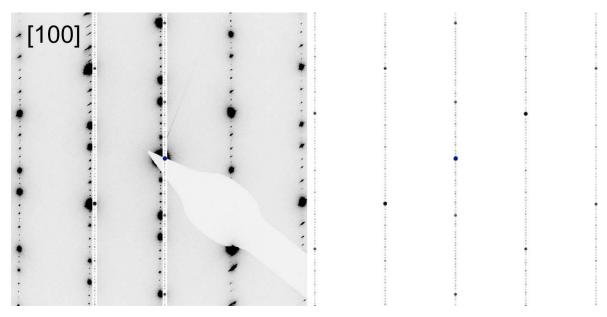
**Table A9.** Comparison of Sn<sub>2.8(4)</sub>Bi<sub>20.2(4)</sub>Se<sub>27</sub> with other Bi-Se compounds that can be described by a (3+1)D superspace formalism, where the commensurate modulation vector determines the actual 3D structure and thus the stacking sequence;<sup>[15]</sup> q = 3 ( $N_{Se}+N_{Bi/Sn-Bi/Sn}$ ) / N with  $N_{Se}$  = number of Se atom layers per unit cell,  $N_{Bi/Sn-Bi/Sn}$  = number of Bi/Sn-Bi/Sn double layers (comparable to those in elemental Bi) per unit cell and N = total number of atom layers per unit cell.

Sum formula	Space group	c / Å	q vector	Number of atom layersper unit cell	Slabs in the unit cell
Bi <sub>2</sub> Se <sub>3</sub>	$R\bar{3}m$	28.6	27/15 = 9/5 = 1.800	15	3x Se-Bi-Se-Bi-Se (5)
Bi <sub>8</sub> Se <sub>9</sub>	R3m	97.1	90/51 = 30/17 = 1.765	51	9x Se-Bi-Se-Bi-Se (5) 3x Bi-Bi (2)
BiSe	$P\bar{3}m1$	22.9	21/12 = 7/4 = 1.750	12	2x Se-Bi-Se-Bi-Se (5)
					Bi-Bi (2)
					15x Se-Bi/Sn-Se-Bi/Sn-Se
Sn2.8(4)Bi20.2(4)Se27	n3	2026	261/150 = 87/50 =	150	(5)
5112.8(4)5120.2(4)5627	Rom	282.0	1.740		9x Se-Bi/Sn-Se-Bi/Sn-Se-
					Bi/Sn-Se (7)
					6x Bi/Sn-Bi/Sn (2)
Bi <sub>8</sub> Se <sub>7</sub>	n31	057	78/45 = 26/15 = 1.733	45	7x Se-Bi-Se-Bi-Se (5)
D18567	Pomi	85.7	/8/43 - 20/13 - 1./33	45	5x Bi-Bi (2)
Bi <sub>4</sub> Se <sub>3</sub>	$R\bar{3}m$	40	36/21 = 12/7 = 1.714	21	3x Se-Bi-Se-Bi-Se (5)
D143C3	кот		30/21 - 12/7 - 1./14	41	3x Bi-Bi (2)



Additional electron diffraction data, STEM images and image of STEM sample

**Figure A10.** a, b) Fourier-filtered STEM-HAADF image along zone axis [210] (Fourier filtered); c) 2D plot of the projected intensity observed in the red box of the left image. The peaks show the stacking sequence 7525757525, which had also been revealed by single-crystal data.



**Figure A11.** Comparison of an experimental SAED pattern (left) along the [100] zone axis with simulated SAED pattern (middle); the model for simulation was taken from the commensurate structure model in  $R\bar{3}m$  with the atom coordinates resulting from the refinement in (3+1)D superspace.

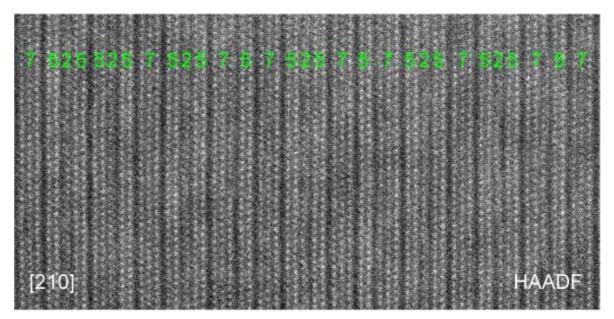


Figure A12. Unfiltered STEM-HAADF image along zone axis [210] showing stacking faults.

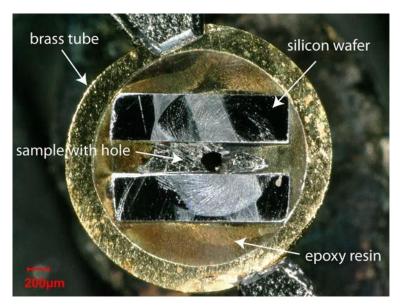


Figure A13. Picture of the prepared sample (cf. experimental section) used for STEM studies.

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# **B** Supporting Information for Chapter 3

# **Experimental Procedures**

## Preparation of starting materials

 $P_3N_5$  was prepared following a procedure by Stock et al. by heating  $P_4S_{10}$  (Sigma-Aldrich, 99,99%) in a constant flow of ammonia gas (5.0, Air Liquide). After drying a fused silica tube and a fused silica boat in a tube furnace (at 1273 K) under dynamic vacuum (<10<sup>-3</sup> mbar) the boat was loaded with  $P_4S_{10}$  in a constant flow of Ar. The apparatus was saturated with ammonia for 4 h and the temperature was increased to 1123 K and maintained for four additional hours. Then, the furnace was cooled to room temperature afterward and heating/cooling ramps were set at 5 K/min each. The product was obtained as an orange powder and purity was confirmed by PXRD as well as CHNS analysis (C 0%, H 0%, N 42.69%, S 0%).<sup>[6]</sup>

The azides  $AE(N_3)_2$  (AE = Sr, Ba) were prepared by reaction of the respective alkaline earth carbonates (SrCO<sub>3</sub>: Sigma Aldrich, 99.995 %; BaCO<sub>3</sub>: Grüssing, 99.8 %) with in situ generated HN<sub>3</sub>.<sup>[7]</sup> An aqueous solution of NaN<sub>3</sub> (Acros Organics, 99 %) was passed through a cation exchanger (Amberlyst 15). Formed HN<sub>3</sub> dropped into a suspension of the alkaline earth carbonate in water until it became clear and the pH of the eluate was neutral. Residual carbonate was filtered off and the water was removed with a rotary evaporator (50 mbar, 40°C). The products were obtained as colorless powders, recrystallized from acetone, and dried under vacuum. The purity of the alkaline earth azides was confirmed by PXRD. Amorphous Si<sub>3</sub>N<sub>4</sub> (UBE, SNA-00) and NH<sub>4</sub>F (Sigma Aldrich,  $\geq$ 98%) were used as purchased.

## High-pressure high-temperature synthesis

The nitridosilicatephosphates *AE*SiP3N7 (*AE* = Sr, Ba) were obtained by high-pressure hightemperature synthesis using a 1000 t press with a modified Walker-type multianvil apparatus.<sup>[8-12]</sup> All products were synthesized from stoichiometric amounts of P<sub>3</sub>N<sub>5</sub>, Si<sub>3</sub>N<sub>4</sub>, the respective alkaline earth azide, and 5 wt% NH<sub>4</sub>F (Table S1). Owing to the air and moisture sensitivity of the azides, the starting materials were handled in an Ar-filled glovebox (Unilab, MBraun, Garching, O<sub>2</sub><1 ppm, H<sub>2</sub>O<0.1 ppm) and ground thoroughly in an agate mortar. The mixture was transferred into a cylindrical crucible made from hexagonal boron nitride (HeBoSint® S100, Henze, Kempten, Germany) and sealed with a boron nitride cap. The crucible was then placed in a specially prepared octahedron (MgO, doped with  $Cr_2O_3$  (5%), edge length 18 mm, Ceramic Substrates & Components, Isle of Wight, UK), which was drilled through and loaded with different components. First, a ZrO<sub>2</sub> sleeve (Cesima Ceramics, Wust-Fischbeck, Germany) was fixed in the borehole and closed on one side by a Mo plate. Then, a long graphite tube (Schunk Kohlenstofftechnik GmbH, Gießen, Germany), a MgO spacer (Cesima Ceramics, Wust-Fischbeck, Germany), and a short graphite tube were put into the octahedron. While ZrO<sub>2</sub> served as a thermal insulator, the two graphite tubes were used as electrical resistance furnace; the combination of two single tubes is necessary to decrease the temperature gradient. After inserting the crucible with the sample, the assembly was completed by adding a second MgO spacer and a further Mo plate to ensure a symmetric setup. In order to generate approximately hydrostatic conditions, the uniaxial force of the hydraulic press was converted by a Walker-type module with inserted steel wedges, forming a cubic cavity. This cavity was filled with eight Co-doped (7%) WC cubes (Hawedia, Marklkofen, Germany). In order to enclose the sample octahedron, the edges of the WC cubes were truncated (11 mm edge length). Furthermore, four of the cubes were equipped with a PTFE film (Vitaflon Technische Produkte GmbH, Bad Kreuznach, Germany) for electrical insulation. The remaining cubes were prepared with pyrophyllite gaskets (Ceramic Substrates & Components, Isle of Wight, UK) to prevent the outflow of the MgO octahedron. Bristol board (369 gm<sup>-2</sup>) was used to fix the gaskets.<sup>[13-15]</sup> Different pressure and temperature settings were used (SrSiP<sub>3</sub>N<sub>7</sub>: 8 GPa, ~1700 °C; BaSiP<sub>3</sub>N<sub>7</sub>: 8 GPa, ~1400 °C).

#### Powder X-ray diffraction

Powder X-ray diffraction was carried out on StadiP diffractometers (STOE & Cie, Darmstadt, Germany) equipped with a MYTHEN 1K detector (Dectris, Baden, Switzerland; angular range  $\Delta 2\theta = 12.5^{\circ}$ ), either a Mo- or Cu-source and a Ge(111) monochromator for Mo-K $\alpha_1$  or Cu-K $\alpha_1$  radiation. Samples were filled into glass capillaries with 0.3 mm diameter and 0.01 mm wall thickness (Hilgenberg GmbH, Malsfeld, Germany). Data were recorded in the angular range between  $2\theta = 5-76^{\circ}$  (BaSiP<sub>3</sub>N<sub>7</sub>) and  $2\theta = 5-120^{\circ}$  (SrSiP<sub>3</sub>N<sub>7</sub>) with a step width of 0.015°/step. Diffraction data were analyzed with the TOPAS-Academic V6 software. Structure models

were refined with the Rietveld method. Estimated standard deviations were calculated by TOPAS-Academic using the matrix inversion method.

#### Scanning electron microscopy and energy-dispersive X-ray spectroscopy (EDX)

Samples were placed on conducting carbon foil and coated with carbon. EDX spectra were collected using a FEI Helios Nanolab G3 Dual Beam UC (FEI, Hillsboro, OR, USA) with an attached X-Max 80 SDD detector (Oxford Instruments, Abingdon, UK).

#### (Scanning) transmission electron microscopy

Samples were ground in absolute ethanol and a drop of the resulting suspension was cast on a TEM grid with holey carbon film (Plano GmbH, Wetzlar, Germany). The grid was mounted on a double-tilt holder and transferred into a Cs DCOR probe-corrected Titan Themis 300 (FEI, USA) TEM equipped with X-FEG, post-column filter (Enfinium ER-799), US1000XP/FT camera system (Gatan, Germany), and a windowless, 4-quadrant Super-X EDX detector. TEM images were recorded using a  $4k \times 4k$  FEI Ceta CMOS camera. The microscope was operated at 300 kV accelerating voltage for SAED and STEM-HAADF and -EDX (convergence angle of 16.6 mrad, 50 µm aperture, detector inner half angle 33 mrad for 245 mm camera length). For evaluation of the TEM data, the following software was used: Digital Micrograph (Fourier filtering of STEM images), ProcessDiffraction7 (geometric calculations for SAED), JEMS (SAED simulations), Velox v2.8 (Thermo Fisher Scientific, USA).

#### Solid-state MAS NMR spectroscopy

Solid-state NMR measurements were performed on an Avance III 500 spectrometer (Bruker, Karlsruhe) equipped with an 11.7 T magnet operating at 500.25 MHz 1H frequency and a commercial double-resonance MAS probe. The sample was ground and loaded in a ZrO2 rotor with an outer diameter of 2.5 mm. NMR spectra were collected at 20 kHz spinning frequency.

#### Single-crystal X-ray diffraction

Single-crystal X-ray diffraction data were collected on a Bruker D8 Venture TXS diffractometer (rotating anode, Mo-K $\alpha$  radiation,  $\lambda = 0.71073$  Å, multilayer monochromator) by combined  $\varphi$ - and  $\omega$ -scans. Indexing, integration, semi-empiric absorption correction

(based on equivalent reflections), were performed by the APEX3 software package. Employing SHELX-2018, the structure was solved by direct methods and refined against  $F^2$  by the full-matrix least-squares method.<sup>[16]</sup>

### Bond-valence sum (BVS) Calculations

BVS calculations were performed using the software VaList v4.0.7.<sup>[17,18]</sup> The structure models of  $AESiP_3N_7$  (AE = Sr, Ba) were taken into account with either all tetrahedra centers occupied by P, occupied by Si, or a mixed occupation of P<sub>0.5</sub>Si<sub>0.5</sub> where it was indicated by BVS values.

## FTIR spectroscopy

FTIR spectra of the title compounds were collected on a Spectrum BX II spectrometer (PerkinElmer, Waltham, MA, USA) between 600 and 4400 cm<sup>-1</sup> with DuraSampler attenuated total reflectance unit (ATR).

### Polyhedra volumes

Comparisons of the different tetrahedra volumes were performed using the software VESTA v3.4.0.<sup>[19]</sup>

## Luminescence measurements

Luminescence spectra of Eu<sup>2+</sup>-doped samples were examined on a HORIBA Fluoromax4 spectrofluorimeter system with an Olympus BX51 microscope. Emission spectra were collected at room temperature with  $\lambda_{exc}$  = 400, in a range from 400 to 800 nm (step width: 2 nm).

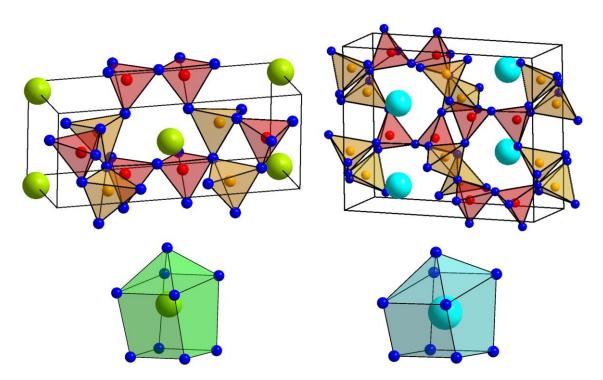
# **Results and Discussion**

## Synthesis

**Table B1** Weighed portions of starting materials for the syntheses of *AE*SiP<sub>3</sub>N<sub>7</sub> (*AE* = Sr, Ba).

Title compounds		Starting mat	erial / mg	
C-C:D N	Sr(N <sub>3</sub> ) <sub>2</sub>	$P_3N_5$	$Si_3N_4$	NH <sub>4</sub> F
SrSiP <sub>3</sub> N <sub>7</sub>	16.79	15.94	4.57	2
DoC:D-N-	Ba(N <sub>3</sub> ) <sub>2</sub>	P <sub>3</sub> N <sub>5</sub>	Si <sub>3</sub> N <sub>4</sub>	NH <sub>4</sub> F
BaSiP <sub>3</sub> N <sub>7</sub>	19.36	19.00	8.18	2

### Structure determination



**Figure B1**. Structure representations of SrSiP<sub>3</sub>N<sub>7</sub> (top left) and BaSiP<sub>3</sub>N<sub>7</sub> (top right) and corresponding *AE*N<sub>9</sub> polyhedra (bottom) (Sr green, Ba light blue, P red, Si<sub>0.5</sub>P<sub>0.5</sub> orange and N blue).

Atom	Wyckoff	X	У	Ζ	$U_{eq}$	s.o.f.
Sr1	2 <i>a</i>	0	0.70251(11)	0.00052(11)	0.01702(15)	1
P1	4b	0.11859(5)	0.17613(14)	0.5034(3)	0.00453(14)	1
P2 Si2	4 <i>b</i>	0.24912(6)	0.31942(15)	0.0235(2)	0.00528(15)	0.5 0.5
N1	4 <i>b</i>	0.3615(2)	0.1391(6)	0.0785(7)	0.0085(5)	1
N2	4 <i>b</i>	0.2811(2)	0.6332(6)	0.1667(6)	0.0062(4)	1
N3	4 <i>b</i>	0.1305(2)	0.2367(7)	0.1673(7)	0.0083(5)	1
N4	2 <i>a</i>	0	0.2788(9)	0.6213(10)	0.0083(7)	1

Table B2 Atomic coordinates, isotropic displacement parameters/ Å<sup>2</sup>, and occupancy of SrSiP<sub>3</sub>N<sub>7</sub>.

**Table B3** Anisotropic displacement parameters for SrSiP<sub>3</sub>N<sub>7</sub> with standard deviations in parentheses.

Atom	<i>U</i> <sub>11</sub>	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	$U_{12}$
Sr1	0.0160(2)	0.0231(2)	0.0119(2)	0.0004(3)	0	0
P1	0.0046(3)	0.0045(3)	0.0045(3)	0.0000(4)	0.0005(3)	0.00006(18)
P2 Si2	0.0059(3)	0.0055(3)	0.0045(3)	0.0004(3)	0.0002(3)	0.0004(2)
N1	0.0088(10)	0.0051(10)	0.0116(12)	-0.0014(8)	-0.0020(8)	0.0001(9)
N2	0.0061(9)	0.0071(10)	0.0053(10)	-0.0002(9)	0.0002(8)	-0.0011(8)
N3	0.0071(10)	0.0136(12)	0.0043(10)	0.0016(9)	-0.0004(8)	-0.0015(9)
N4	0.0028(13)	0.0134(17)	0.0087(17)	-0.0010(13)	0	0

Table B4 Atomic coordinates, isotropic displacement parameters/ Å<sup>2</sup>, and occupancy of BaSiP<sub>3</sub>N<sub>7</sub>.

Atom	Wyckoff	X	у	Ζ	$U_{eq}$	s.o.f.
Ba1	4 <i>c</i>	0.14950(3)	0.25	0.76102(6)	0.01189(8)	1
P1	8 <i>d</i>	0.40923(7)	0.13039(6)	0.29058(14)	0.00225(13)	1
P2 Si2	8 <i>d</i>	0.15603(8)	0.00047(6)	0.28611(15)	0.00256(14)	0.5 0.5
N1	8 <i>d</i>	0.0612(3)	0.1098(2)	0.3097(6)	0.0065(4)	1
N2	4 <i>c</i>	0.3567(4)	0.25	0.2000(8)	0.0068(6)	1
N3	8 <i>d</i>	0.3045(3)	0.0401(2)	0.1254(5)	0.0039(4)	1
N4	8 <i>d</i>	0.1108(3)	0.6112(2)	0.1231(5)	0.0058(4)	1

Atom	$U_{11}$	U <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	$U_{12}$
Ba1	0.01304(13)	0.01160(12)	0.01103(12)	0	0.00102(10)	0
P1	0.0018(3)	0.0028(3)	0.0021(3)	-0.0003(2)	-0.0002(2)	-0.0003(2)
P2 Si2	0.0028(3)	0.0033(3)	0.0016(3)	0.0000(2)	0.0001(3)	0.0000(2)
N1	0.0036(10)	0.0063(10)	0.0097(11)	-0.0013(8)	-0.0017(9)	0.0004(9)
N2	0.0091(15)	0.0030(13)	0.0083(15)	0	-0.0022(13)	0
N3	0.0036(10)	0.0038(10)	0.0043(9)	-0.0007(8)	-0.0001(8)	-0.0016(9)
N4	0.0074(11)	0.0058(11)	0.0042(9)	0.0005(8)	0.0005(8)	0.0023(9)

Table B5 Anisotropic displacement parameters for BaSiP<sub>3</sub>N<sub>7</sub> with standard deviations in parentheses.

Table B6 Interatomic distances (Å) and bond angles (°) in the structure of  $BaSiP_3N_7$ .

Ba1	N1	2x	2.872(3)	N1	-P1-	N4	112.68(14)
	N2	1x	2.921(4)	N1		N2	111.60(17)
	N4	2x	2.987(3)	N1		N3	109.65(13)
	N4	2x	3.132(3)	N4		N2	110.76(17)
	N1	2x	3.230(3)	N4		N3	106.33(13)
P1	N1	1x	1.598(3)	N2		N3	105.43(16)
	N4	1x	1.604(3)	N4	-P2 Si2-	N1	123.79(14)
	N2	1x	1.6062(18)	N4		N3	105.02(13)
	N3	1x	1.703(3)	N4		N3	105.35(13)
P2 Si2	N4	1x	1.627(3)	N1		N3	107.47(13)
	N1	1x	1.635(3)	N1		N3	106.97(13)
	N3	1x	1.724(3)	N3		N3	107.32(9)
	N3	1x	1.726(3)				
-							

Sr1	N1	2x	2.696(3)	N4	-P1-	N3	111.0(2)
	N4	1x	2.735(4)	N4		N1	110.99(19)
	N3	2x	2.876(3)	N4		N2	107.28(18)
	N3	2x	3.149(3)	N3		N1	112.31(18)
	N1	2x	3.270(3)	N3		N2	106.02(15)
P1	N4	1x	1.605(2)	N1		N2	109.01(16)
	N3	1x	1.609(4)	N3	-P2 Si2-	N1	121.41(16)
	N1	1x	1.604(3)	N3		N2	104.84(16)
	N2	1x	1.704(3)	N3		N2	104.59(15)
P2 Si2	N3	1x	1.624(3)	N1		N2	113.53(16)
	N1	1x	1.631(3)	N1		N2	103.82(14)
	N2	1x	1.722(3)	N2		N2	107.70(11)
	N2	1x	1.727(3)				

Table B7 Interatomic distances (Å) and bond angles (°) in the structure of SrSiP<sub>3</sub>N<sub>7</sub>.

#### **EDX measurements**

**Table B8.** EDX measurements (several point analyses) for  $AESiP_3N_7$  (AE = Sr, Ba) compared to valuesderived from sum formulas.

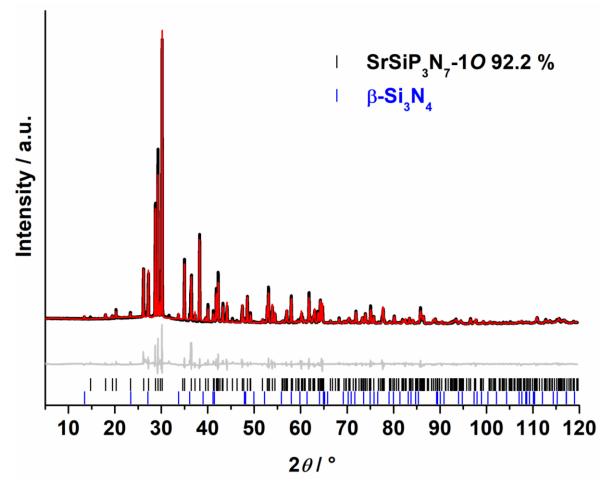
atom %	Sr	Si	Р	Ν	0
	8.3	9.5	25.4	56.8	0
	9.2	7.1	24.4	53.4	5.9
	8.2	9.4	25.1	53.7	3.5
	7.3	7.1	21.6	57.2	6.9
	5.9	6.1	19.8	64.9	3.2
	6.1	6.2	20.3	61.0	6.2
	7.3	7.1	21.6	57.2	6.9
	6.0	6.1	20.1	61.4	6.3
calc.	8.3	8.3	25	58.3	0
	Ba	Si	Р	Ν	0
	9.2	8.4	26.6	51.5	4.3
	8.9	7.3	25.9	50.2	7.7
	8.2	7.9	24.8	54.8	4.3
	8.3	7.5	27.4	48.1	8.7
	6.3	5.7	21.4	56.8	9.9
calc.	8.3	8.3	25	58.3	0

## **Rietveld refinements**

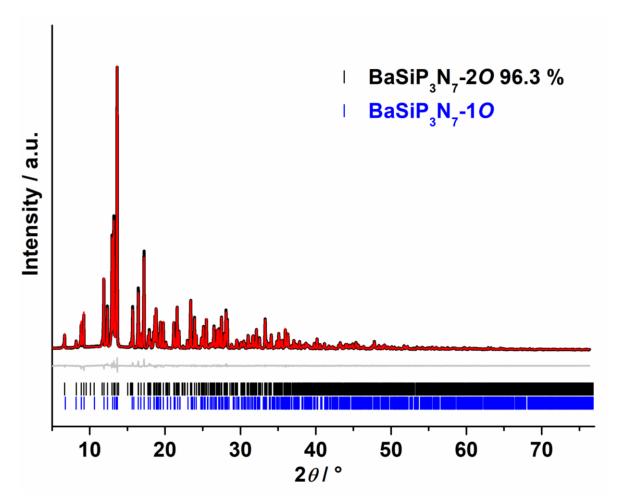
**Table B9.** Crystallographic data from the Rietveld refinements for  $AESiP_3N_7$  (AE = Sr, Ba), standarddeviations are given in parentheses. For BaSiP<sub>3</sub>N<sub>7</sub>, both polytypes were taken into account.

	SrSiP <sub>3</sub> N <sub>7</sub>	BaSiP	3N7	
phase fraction	92.2(9)*)	96.3(1)	3.7(1)	
crystal system	orthorhombic	orthorhombic	orthorhombic	
space group	<i>Pmn</i> 2 <sub>1</sub> (no. 31)	<i>Pnma</i> (no. 62)	<i>Pmn</i> 2 <sub>1</sub> (no. 31)	
formula units/ unit cell	2	4	2	
lattice parameters / Å	a = 11.9975(2)	a = 9.90917(5)	a = 11.9975(2)	
	b = 4.9236(1)	b = 12.19036(8)	b = 4.9236(1)	
	c = 4.70029(9)	c = 4.73862(3)	c = 4.70029(9)	
radiation $\lambda$ / Å	1.54056 (Cu- <i>K</i> α1)	0.70930(N	Λο-Κα1)	
cell volume / ų	277.65(1)	572.409(6)	283.28(6)	
$2\theta$ – range/°	5 <2 <i>θ</i> < 120	5 <2 <i>0 &lt; 1</i>	76.35	
refined parameters	29	32		
(incl. side phases)				
background function	Shifted Cl	nebyshev (12 parameters)		
<i>R</i> -values	$R_{\rm p} = 0.0609$	$R_{\rm p} = 0.0296$ ,		
	$wR_{\rm p} = 0.0929$	$wR_{\rm p}=0.$	0396	
	$R_{\mathrm{Bragg}} = 0.0353$	$R_{\mathrm{Bragg}} = 0.0165$	$R_{\mathrm{Bragg}} = 0.0191$	

\*) Si<sub>3</sub>N<sub>4</sub> as a side phase



**Figure B2**. Rietveld refinement for SrSiP<sub>3</sub>N<sub>7</sub>-1*O*; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of SrSiP<sub>3</sub>N<sub>7</sub>-1O (vertical black bars),  $\beta$ -Si<sub>3</sub>N<sub>4</sub> (vertical blue bars), and difference profile (gray line).



**Figure B3**. Rietveld refinement for BaSiP<sub>3</sub>N<sub>7</sub>-2*O*; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of BaSiP<sub>3</sub>N<sub>7</sub>-2O (vertical black bars), BaSiP<sub>3</sub>N<sub>7</sub>-1*O* (vertical blue bars), and difference profile (gray line).

## Polyhedra volumes

**Table B10** Tetrahedra volumes of several known nitridosilicates and nitridophosphates.

$T\mathrm{N}_4$ volumes in Å $^3$								
	SiN <sub>4</sub>		PN <sub>4</sub>					
$Sr_2Si_5N_8^{[20]}$	2.5450, 2.7127	$BaP_2N_4^{[21]}$	2.2816					
$La_3Si_6N_{11}^{[22]}$	2.6125, 2.6964	$BaP_8N_{14}^{[23]}$	2.1911, 2.2289					
$Ce_5Si_3N_9^{[24]}$	2.5182, 2.7600, 2,5840	$Ce_2P_3N_7^{[25]}$	2.2511, 2.2805					
$SrSi_7N_{10}^{[26]}$	2.6844, 2.6067, 2.6183	$Cs_3P_6N_{11}^{[27]}$	2.2299					
$CeSi_3N_5^{\left[28\right]}$	2.6041, 2.6677, 2.6199	LiPN2 <sup>[29]</sup>	2.2687					
$LiSi_2N_3^{[30]}$	2.6870	BP <sub>3</sub> N <sub>6</sub> [31]	2.2223, 2.1882, 2.1313					
$LaSi_3N_5^{\left[28\right]}$	2.6907, 2.6425, 2.6265	$NaP_4N_7^{[32]}$	2.2465, 2.2310					
Ø	2.6398		2.2292					

Tetrahedra volumes/ ų	$Pnm2_1$	<i>P</i> 2 <sub>1</sub>	Pn	Pm
1	2.2423	2.2177	2.2225	2.2212
2	2.4024	2.2183	2.2162	2.2124
3	-	2.3779	2.3758	2.3703
4	-	2.3753	2.3754	2.3860

**Table B11.** Tetrahedra volumes for  $SrSiP_3N_7$  in space groups  $Pnm2_1$ ,  $P2_1$ , Pn, and Pm. Red values indicate deviations from expected volumes for  $PN_4$  tetrahedra (V ~2.2 Å<sup>3</sup>).

**Table B12.** Tetrahedra volumes for BaSiP<sub>3</sub>N<sub>7</sub> in space groups Pnma, Pna2<sub>1</sub>, P2<sub>1</sub>/c, and P2<sub>1</sub>ma. Red values indicate deviations from expected volumes for PN<sub>4</sub> tetrahedra (V  $\sim$ 2.2 Å<sup>3</sup>).

Tetrahedra volumes/ Å <sup>3</sup>	Pnma	$Pna2_1$	$P2_{1}/c$	$P2_1ma$
1	2.2105	2.2076	2.2128	2.1965
2	2.3839	2.3915	2.2084	2.4043
3	-	2.2118	2.3825	2.3639
4	-	2.3718	2.3850	2.2300

#### **BVS calculations**

**Table B13**. BVS for SrSiP<sub>3</sub>N<sub>7</sub> in space group Pnm2<sub>1</sub> and subgroups **P2<sub>1</sub>**, **Pn**, and **Pm**. The structures were refined against SCXRD with all tetrahedra center positions occupied by P. Red values indicate mixed occupancy of P/Si (BVS  $P_{0.5}Si_{0.5} = 4.5$ ).

Atom	$Pnm2_1$	<i>P</i> 2 <sub>1</sub>	Pn	Pm
Sr1	1.63	1.67	1.67	1.67
Sr2	-	-	-	1.67
P1	4.92	5.00	4.99	4.99
P2	4.40	4.99	5.01	5.02
P3	-	4.47	4.47	4.49
P4	-	4.48	4.47	4.45

<b>Table B14.</b> BVS for SrSiP <sub>3</sub> N <sub>7</sub> in space group Pnm2 <sub>1</sub> and subgroups <b>P2<sub>1</sub></b> , <b>Pn</b> , and <b>Pm</b> . The structures
were refined against SCXRD with all tetrahedra center positions occupied by Si. Red values indicate
mixed occupancy of P/Si (BVS $P_{0.5}Si_{0.5} = 4.5$ ).

$Pnm2_1$	<i>P</i> 2 <sub>1</sub>	Pn	P <i>m</i>
1.63	1.68	1.67	1.67
-	-	-	1.68
5.22	5.31	5.32	5.30
4.62	5.31	5.32	5.33
-	4.69	4.70	4.70
-	4.70	4.69	4.67
	1.63 - 5.22 4.62 -	1.63       1.68         -       -         5.22       5.31         4.62       5.31         -       4.69	1.63       1.68       1.67         -       -       -         5.22       5.31       5.32         4.62       5.31       5.32         -       4.69       4.70

**Table B15**. BVS for SrSiP<sub>3</sub>N<sub>7</sub> in space group Pnm2<sub>1</sub> and subgroups **P2<sub>1</sub>**, **Pn**, **and Pm**. The structure models were refined against SCXRD with tetrahedra center positions occupied by P (positions with a BVS value of ~5) and tetrahedra center positions occupied by P<sub>0.5</sub>Si<sub>0.5</sub> (positions with a BVS value of ~4.5). Red values show a mixed occupancy of P/Si with their respective occupation factor indicated by BVS.

Atom	$Pnm2_1$	P21	P1n1	Pm
Sr1	1.63	1.68	1.68	1.68
Sr2	-	-	-	1.68
Si1/P1	4.91	5.00	4.99	4.99
Si2/P2	4.65/4.40	4.99	5.01	5.01
	48/52			
Si3/P3	-	4.71/4.46	4.71/4.47	4.74/4.49
		43/57	43/57	41/59
Si4/P4	-	4.72/4.47	4.71/4.46	4.69/4.44
		42/58	43/57	45/55

Atom	Pnma	Pna2 <sub>1</sub>	P21/c	$P2_1ma$
Ba1	2.46	2.47	2.46	2.43
Ba2	-	-	-	2.54
P1	5.01	5.01	5.01	5.06
P2	4.42	4.44	5.02	4.38
Р3	-	5.03	4.44	4.45
P4	-	4.43	4.43	4.93

**Table B16.** BVS for BaSiP<sub>3</sub>N<sub>7</sub> in space group Pnma and subgroups **Pna2**<sub>1</sub>, **P2**<sub>1</sub>**ma**, and **P2**<sub>1</sub>/c. The structure models were refined against SCXRD with all tetrahedra center positions occupied by P. Red values indicate mixed occupancy of P/Si (BVS  $P_{0.5}Si_{0.5} = 4.5$ ).

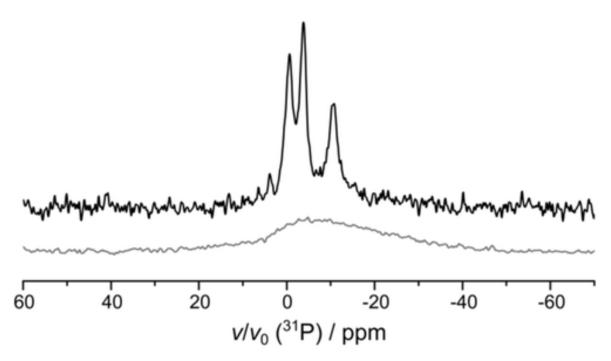
**Table B17.** BVS for BaSiP<sub>3</sub>N<sub>7</sub> in space group Pnma and subgroups Pna2<sub>1</sub>, P2<sub>1</sub>ma, and P2<sub>1</sub>/c. The structure models were refined against SCXRD with all tetrahedra center positions occupied by Si. Red values indicate mixed occupancy of P/Si (BVS  $P_{0.5}Si_{0.5} = 4.5$ ).

Atom	Pnma	$Pn2_1a$	P2 <sub>1</sub> /c	$P2_1ma$
Ba1	2.46	2.46	2.46	2.43
Ba2	-	-	-	2.53
Si1	5.30	5.27	5.30	5.30
Si2	4.68	4.70	5.30	4.59
Si3	-	5.35	4.68	4.73
Si4	-	4.66	4.67	5.25

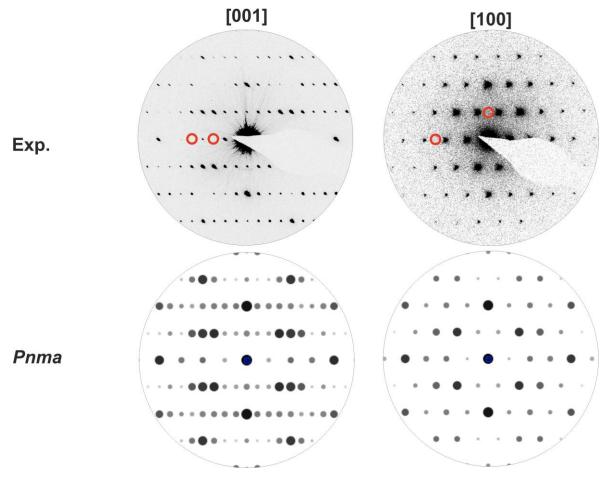
**Table B18.** BVS for BaSiP<sub>3</sub>N<sub>7</sub> in space group Pnma and subgroups Pna2<sub>1</sub>, P2<sub>1</sub>ma, and P2<sub>1</sub>/c. The structure models were refined against SCXRD with tetrahedra center positions occupied by P (positions with a BVS value of ~5) and tetrahedra center positions occupied by P<sub>0.5</sub>Si<sub>0.5</sub> (positions with a BVS value of ~4.5). Red values show a mixed occupancy of P/Si with their respective occupation factor indicated by BVS.

Atom	Pnma	$Pn2_1a$	$P2_{1}/c$	$P2_1ma$
Ba1	2.46	2.47	2.46	2.57
Ba2	-	-	-	2.40
Si1/P1	5.01	4.99	5.01	4.94
Si2/P2	4.68/4.43	4.69/4.44	5.02	4.69/4.45
	46/54	45/55		44/66
Si3/P3	-	5.04	4.68/4.44	4.63/4.39
			45/55	49/51
Si4/P4	-	4.68/4.43	4.67/4.43	5.05
		45/55	46/54	

Solid-state NMR spectra

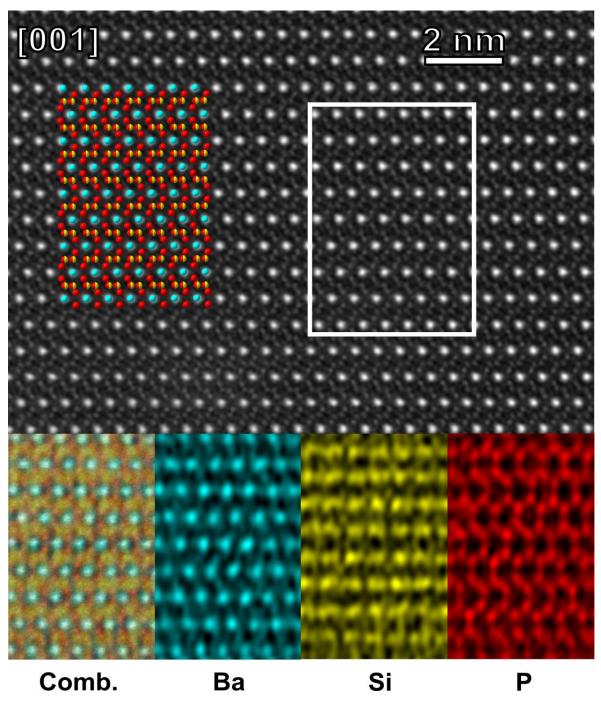


**Figure B4.** <sup>31</sup>P solid-state NMR spectrum (MAS) of BP<sub>3</sub>N<sub>6</sub> showing three narrow signals (fwhm < 5 ppm).<sup>[31]</sup>



(Scanning) transmission electron microscopy

**Figure B5.** Experimental electron diffraction patterns of BaSiP<sub>3</sub>N<sub>7</sub> (top) and simulations for the final structure model in *Pnma*. Reflections, affected by systematic absences, are marked by a red circle.



**Figure B6**. STEM and EDX mapping of BaSiP<sub>3</sub>N<sub>7</sub> along [001]. STEM HAADF image (top) with structure overlay (Ba cyan, P red, Si yellow). The inset shows the corresponding area for EDX maps (bottom, atomic resolution) with a combined color map, Ba map (light blue), Si map (yellow), and P map (red).

# FTIR spectra

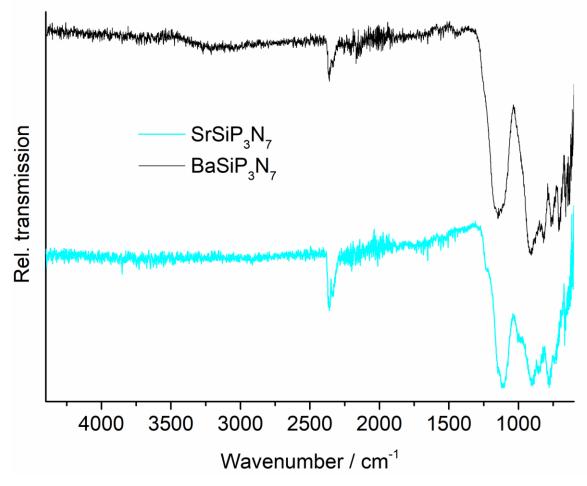


Figure B7. FTIR spectra for SrSiP<sub>3</sub>N<sub>7</sub> (blue) and BaSiP<sub>3</sub>N<sub>7</sub> (black).

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# C Supporting Information for Chapter 4

# **Experimental Procedures**

#### Preparation of starting materials

 $P_3N_5$  was prepared following a procedure by Stock et al.<sup>[6]</sup> The product was obtained as an orange powder. Purity was confirmed by PXRD as well as CHNS analysis (C 0%, H 0%, N 42.69%, S 0% expected C 0%, H 0%, N 42.98%, S 0%).

The amides  $AE(NH_2)_2$  (AE = Sr, Ba) were prepared by reaction of the respective alkaline earth metal and supercritical ammonia according to the literature.<sup>[33,34]</sup>

Amorphous Si<sub>3</sub>N<sub>4</sub> (UBE, SNA-00) and NH<sub>4</sub>F (Sigma Aldrich,  $\geq$ 98%) were used as purchased.

#### High-pressure high-temperature synthesis

The silicon imidonitridophosphates  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg,  $Mg_{0.94}Ca_{0.06}$ , Ca, Sr) were obtained by high-pressure high-temperature synthesis using a 1000 t press with a modified Walker-type multianvil apparatus.<sup>[8-10,35,36]</sup> All products were synthesized from stoichiometric amounts of P<sub>3</sub>N<sub>5</sub>, Si<sub>3</sub>N<sub>4</sub>, the respective alkaline earth amide, and 5 wt% NH<sub>4</sub>F (Table S1). Owing to the air and moisture sensitivity of the amides, starting materials were handled in an Ar-filled glovebox (Unilab, MBraun, Garching, O<sub>2</sub><1 ppm, H<sub>2</sub>O<0.1 ppm) and ground thoroughly in an agate mortar. Details on the preparation and handling of the 1000 t Walker-type multianvil press are described in the literature.<sup>[14,15,37]</sup> All reactions were carried out at 8 GPa and 1400°C.

#### Powder X-ray diffraction

Powder X-ray diffraction was carried out on a StadiP diffractometer (STOE & Cie, Darmstadt, Germany) equipped with an MYTHEN 1K detector (Dectris, Baden, Switzerland; angular range  $\Delta 2\theta = 12.5^{\circ}$ ), a Cu source and a Ge(111) monochromator for Cu-K $\alpha_1$  radiation. Samples were filled into glass capillaries with 0.3 mm diameter and 0.01 mm wall thickness (Hilgenberg GmbH, Malsfeld, Germany). Data were recorded in the angular range 5° < 2 $\theta$  < 120°. Diffraction data were analyzed with the TOPAS-Academic V6 software, where phase

compositions were analyzed with the Rietveld method. Rietveld refinements were performed with atomic and lattice parameters from  $CaSi_3P_4N_{10}(NH)_2$  as starting values. H positions and isotropic displacement parameters were not refined.

Unit cell volumes of  $AESi_{3}P_{4}N_{10}(NH)_{2}$  (AE = Mg,  $Mg_{0.94}Ca_{0.06}$ , Ca) were extracted from Pawley Fits with Si (a = 5.43054 Å) as an internal standard. The Ca content of  $Mg_{0.94}Ca_{0.06}Si_{3}P_{4}N_{10}(NH)_{2}$  was determined from the unit cell volume of  $Mg_{0.94}Ca_{0.06}Si_{3}P_{4}N_{10}(NH)_{2}$  in relation to the end members  $MgSi_{3}P_{4}N_{10}(NH)_{2}$  and  $CaSi_{3}P_{4}N_{10}(NH)_{2}$ .

High-temperature PXRD patterns were recorded on a StadiP diffractometer (STOE & Cie, Darmstadt, Germany) in the range  $2^{\circ} < 2\theta < 75^{\circ}$  and Mo-K $\alpha_1$  radiation ( $\lambda = 0.709300$  Å). An image-plate position-sensitive detector was used for X-ray detection. Diffraction patterns were collected in 50°C steps up to 950 °C. Samples were filled into silica capillaries with 0.4 mm diameter and 0.01 mm wall thickness (Hilgenberg GmbH, Malsfeld, Germany). Lattice parameters were extracted from Pawley-Fits.

#### Scanning electron microscopy and energy-dispersive X-ray spectroscopy (EDX)

Samples were placed on conducting carbon foil and coated with carbon. EDX spectra were collected using an FEI Helios Nanolab G3 Dual Beam UC (FEI, Hillsboro, OR, USA) with an attached X-Max 80 SDD detector (Oxford Instruments, Abingdon, UK).

#### (Scanning) transmission electron microscopy

Samples were ground in absolute ethanol and a drop of the resulting suspension was cast on a TEM grid with lacey carbon film (Plano GmbH, Wetzlar, Germany). The grid was mounted on a double-tilt low background holder and transferred into a Cs DCOR probe-corrected Titan Themis 300 (FEI, USA) TEM equipped with X-FEG, post-column filter (Enfinium ER-799), US1000XP/FT camera system (Gatan, Germany), and a windowless, 4-quadrant Super-X EDX detector (FEI, USA) TEM images were recorded using a  $4k \times 4k$  FEI Ceta CMOS camera. The microscope was operated at 300 kV accelerating voltage for SAED, STEM-HAADF and EDX (semiconvergence angle 16.6 mrad, 50  $\mu$ m aperture, HAADF detector inner half angle 33 mrad for 245 mm camera length). The element maps were filtered using a radial Wiener filter (highest frequency 20.0 software-specific units and edge smoothing 30.0 units). Samples for data collection with microfocused synchrotron radiation were prepared as described above but TEM finder grids were used. Suitable crystallites were identified by EDX and electron diffraction. Their positions on the grid were mapped using bright-field images at different magnifications. This allowed the correct positioning of the grid in the synchrotron beam. For evaluation of the TEM data, the following software was used: Digital Micrograph (Fourier filtering of STEM images), ProcessDiffraction7 (geometric calculations for SAED), JEMS (SAED simulations), Velox v2.8 (Thermo Fisher Scientific, USA for EDX mappings).

#### Solid-state MAS NMR spectroscopy

Solid-state NMR measurements were performed on an Avance III 500 spectrometer (Bruker, Karlsruhe) equipped with an 11.7 T magnet operating at 500.25 MHz  $^{1}$ H frequency and a commercial double-resonance MAS probe. The sample was ground and loaded in a ZrO<sub>2</sub> rotor with an outer diameter of 2.5 mm. NMR spectra were collected at 20 kHz spinning frequency and a relaxation delay of 12 s.

#### Single-crystal X-ray diffraction

The crystallites on TEM finder grids were optically centered in the synchrotron beam with a high magnification telescope at ESRF beamline ID11. Centering was precisely optimized by fluorescence and diffraction scans utilizing a hexapod setup (Symétrie Hexapods Nanopos and PI-MARS P563 piezo stage).<sup>[38-40]</sup> The datasets of two crystallites were evaluated and combined for the sake of better completeness using CrysalisPro (Agilent Technologies). Absorption correction was performed using SADABS.<sup>[41]</sup> Employing SHELX-2018, the structure was solved by direct methods and refined against  $F^2$  by the full-matrix least-squares method.<sup>[16]</sup> Raw diffraction data are archived at the ESRF.<sup>[42]</sup>

#### Bond-valence sum (BVS) Calculations

BVS calculations were performed using the software VaList v4.0.7.<sup>[17,18]</sup>

#### Luminescence measurements

Luminescence spectra of Eu<sup>2+</sup>-doped samples were examined on a HORIBA Fluoromax4 spectrofluorimeter system with an Olympus BX51 microscope. Emission spectra were collected at room temperature with  $\lambda_{exc}$  = 400, in a range from 400 to 800 nm (step width:

2 nm). Cryo spectra between 6 K and 300 K were recorded in an evacuated, He-cooled sample chamber with an Ocean Optics HR2000 + ES spectrometer (2.048 pixels, grating UA (200-1.100 nm) slit 50).

#### IR spectroscopy

FTIR spectra of the title compounds were collected on a Spectrum BX II spectrometer (PerkinElmer, Waltham, MA, USA) between 600 and 4400 cm<sup>-1</sup> with DuraSampler attenuated total reflectance unit (ATR).

#### UV-Vis spectroscopy

A V-650 UV-Vis spectrophotometer (JASCO, Gross-Umstadt, Germany) equipped with a photomultiplier tube detector and a single monochromator with 1200 lines/mm was used for recording UV-Vis spectra in the range of 240 to 800 nm. A deuterium (240–330 nm) and a halogen lamp (330–880 nm) were used with a 2 nm resolution and a scan speed of 400 nm/min controlled by the Spectra Manager II software. Samples were fixed between a fused silica glass slide and a BaSO<sub>4</sub>-coated stamp. Pseudo-absorption spectra were calculated by the Kubelka-Munk function  $F(R) = (1-R)^2/2R$  where *R* is reflectance. In the Tauc plots, *hv* was plotted against  $[F(R)hv]^{1/n}$  with n = 1/2 for direct bandgaps.<sup>[43]</sup>

# **Results and Discussion**

#### Synthesis

**Table C1** Weighed portions of starting materials for the syntheses of  $AESi_3P_4N_{10}(NH)_2$  ( $AE = Mg_{0.94}Ca_{0.06}$ , Mg, Ca, Sr).

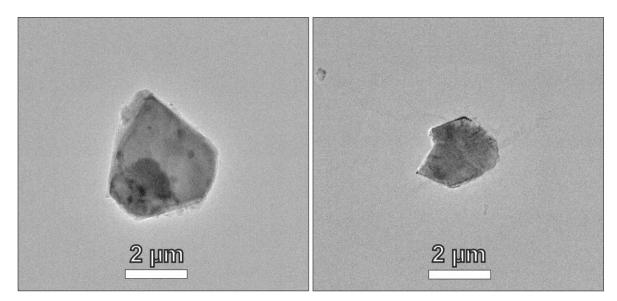
Compound		Starting	gmaterial	
$Mg_{0.94}Ca_{0.06}Si_3P_4N_{10}(NH)_2$	$Mg(NH_2)_2$	$P_3N_5$	$Si_3N_4$	NH <sub>4</sub> F
	14.9 mg,	64.1 mg,	41.4 mg,	6 mg
	0.27 mmol	0.39 mmol	0.30 mmol	
	$Ca(NH_2)_2$			
	2.1 mg,			
	0.03 mmol			
$MgSi_3P_4N_{10}(NH)_2$	$Mg(NH_2)_2$	$P_3N_5$	$Si_3N_4$	NH <sub>4</sub> F
	5.6 mg,	21.59 mg,	13.94 mg,	2 mg
	0.10 mmol	0.13 mmol	0.10 mmol	
CaSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	Ca(NH <sub>2</sub> ) <sub>2</sub>	P <sub>3</sub> N <sub>5</sub>	Si <sub>3</sub> N <sub>4</sub>	NH <sub>4</sub> F
	6.93 mg,	20.88 mg,	13.48 mg,	2 mg
	0.10 mmol	0.13 mmol	0.10 mmol	
SrSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	Sr(NH <sub>2</sub> ) <sub>2</sub>	$P_3N_5$	$Si_3N_4$	NH <sub>4</sub> F
	10.27,	18.66,	12.04,	2 mg
	0.09 mmol	0.11 mmol	0.09 mmol	

### Structure determination

As diffraction data of both crystallites showed threefold twinning, reflection files of each crystallite were first transformed to SHELX HKLF 5 files taking into account the *hkl* transformation matrix shown below.<sup>[44]</sup> The two crystallites showed slightly different twin volume fractions of the threefold twinning, which was taken into account by refinements with six BASF where BASF 1–3 represent the first crystal and BASF 4–6 represent the second crystal. This allowed merging of the two datasets and refinement of the structure model with increased data completeness.

formula	CaSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>		
molar mass /g·mol⁻¹	418.37		
crystal system	monoclinic		
space group	<i>C</i> 2/ <i>m</i> (no. 12)		
lattice parameters / Å, °	a = 4.91790(10)		
	<i>b</i> = 8.5159(2)		
	c = 9.8267(2)		
	$\beta = 99.476(3)^{\circ}$		
cell volume / ų	405.930(15)		
formula units/ unit cell	2		
density / g·cm <sup>3</sup>	3.423		
$\mu$ / mm <sup>-1</sup>	0.247		
temperature / K	293(2)		
absorption correction	semiempirical		
radiation	synchrotron, ESRF ID11 ( $\lambda$ = 0.3252 Å)		
F(000)	416		
$\theta$ range / °	$2.2 \le \theta \le 21.9$		
total no. of reflections	16003		
Independent reflections $[I \ge 2\sigma(I) / all]$	4198/ 4895		
$R_{\sigma}$	0.0383		
refined parameters / restraints	61 / 1		
Goodness of fit	1.087		
<i>R</i> -values [ $I \ge 2\sigma(I)$ ]	$R_1 = 0.0805; wR_2 = 0.2262$		
<i>R</i> -values (all data)	$R_1 = 0.0862; wR_2 = 0.2398$		
$\Delta ho_{ m max}$ , $\Delta ho_{ m min}$ / e·Å <sup>-3</sup>	1.563, -1.558		

**Table C2** Crystallographic data of the single-crystal structure refinements of  $CaSi_3P_4N_{10}(NH)_2$ .Standard deviations are given in parentheses.



**Figure C1**. TEM bright-field images of the crystallites used for data collection with microfocused synchrotron radiation.

Table C3 Atomic coordinates, isotropic displacement parameters (/Å<sup>2</sup>), and site occupancies of  $CaSi_3P_4N_{10}(NH)_2$ .

Atom	Wyckoff	X	у	Z	$U_{eq}$	<i>0CC.</i>
	position					
Ca1	2 <i>a</i>	0	0	0	0.0124(5)	1
P1	8 <i>j</i>	0.0738(3)	0.33298(16)	0.22228(16)	0.0126(4)	1
Si1	4h	0	0.1643(2)	1/2	0.0132(6)	1
Si2	2 <i>d</i>	0	1/2	1/2	0.0127(7)	1
N1	4 <i>i</i>	0.1419(17)	0	0.4097(7)	0.0109(11)	1
N2	8 <i>j</i>	0.1313(15)	0.3318(5)	0.3978(5)	0.0122(8)	1
N3	4 <i>i</i>	0.6701(12)	0	0.1680(8)	0.0129(14)	1
N4	8 <i>j</i>	0.2496(10)	0.1927(6)	0.1663(6)	0.0131(9)	1
H1	4 <i>i</i>	0.12(3)	0	0.319(5)	0.04(4)	1

**Table C4** Anisotropic displacement parameters (*Uij* in Å<sup>2</sup>) for CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> with standard deviations in parentheses. The anisotropic displacement factor is expressed as  $\exp[-2\pi^2 (U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}l^2c^{*2} + U_{12}hka^*b^* + U_{13}hla^*c^* + U_{23}klb^*c^*)$ .

Atom	$U_{11}$	$U_{22}$	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	$U_{12}$
Ca1	0.0100(8)	0.0132(10)	0.0156(11)	0	0.0064(7)	0
P1	0.0089(7)	0.0121(8)	0.0184(9)	0.0000(5)	0.0074(7)	0.0000(3)
Si1	0.0089(12)	0.0145(16)	0.0187(12)	0	0.0096(13)	0
Si2	0.0084(14)	0.0145(18)	0.0173(17)	0	0.0085(14)	0
N1	0.009(3)	$U_{11}$	0.016(3)	0	0.005(3)	0
N2	0.011(2)	0.014(3)	<i>U</i> <sub>22</sub>	-0.0006(16)	0.008(2)	0.0003(12)
N3	0.012(3)	0.010(3)	0.019(4)	0	0.008(2)	0
N4	0.0095(19)	0.0117(18)	0.020(3)	-0.0015(18)	0.0072(17)	0.0009(14)

Table C5 Interatomic distances (Å) and bond angles (°) in the structure of  $CaSi_3P_4N_{10}(NH)_2$ .

Ca1	N4	4x	2.492(5)	N4	P1	N3	109.5(3)
Ca1	N3	2x	2.498(6)	N4	P1	N4	109.7(3)
P1	N4	1x	1.614(6)	N3	P1	N4	109.5(3)
P1	N3	1x	1.617(4)	N4	P1	N2	109.5(3)
P1	N4	1x	1.622(6)	N3	P1	N2	109.3(3)
P1	N2	1x	1.702(5)	N4	P1	N2	109.3(3)
Si1	N1	1x	1.854(5)	N1	Si1	N1	82.0(4)
Si1	N2	1x	1.916(5)	N1	Si1	N2	176.8(3)
Si1	N2	1x	1.923(7)	N1	Si1	N2	97.2(3)
Si2	N1	2x	1.837(8)	N2	Si1	N2	83.8(3)
Si2	N2	4x	1.921(5)	N1	Si1	N2	82.3(3)
N1	H1	1x	0.88(5)	N1	Si1	N2	99.2(3)
				N2	Si1	N2	94.8(2)
				N2	Si1	N2	178.0(3)
				N1	Si2	N1	180
				N1	Si2	N2	82.8(2)
				N1	Si2	N2	97.2(2)
				N2	Si2	N2	180
				N2	Si2	N2	96.4(3)
				N2	Si2	N2	83.6(3)
				N1	Si2	N2	82.8(2)
				Si2	N1	H1	117(9)
				Si1	N1	H1	119(4)

# EDX measurements

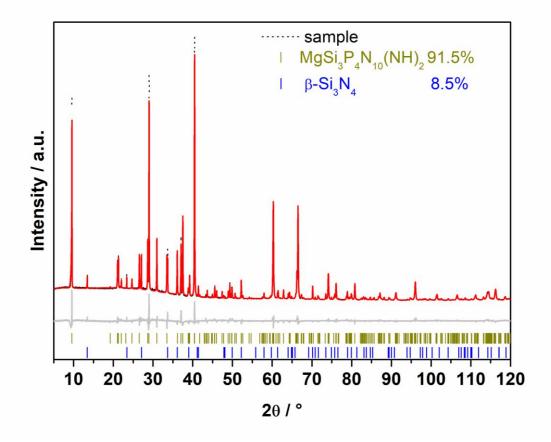
**Table C6.** EDX measurements (several point analyses) for  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg, Ca, Sr) comparedto values derived from sum formulas.

EDX point / atom %	Mg	Si	Р	Ν
1	7.0	17.1	19.1	56.7
2	7.0	17.0	20.3	55.5
3	7.5	18.4	20.2	53.7
4	8.1	16.9	22.0	52.9
average	7.4(4)	17.4(6)	20(1)	54(1)
	Ca	Si	Р	Ν
1	4.8	13.8	18.9	62.4
2	5.0	15.6	20.7	58.6
3	5.1	15.4	20.9	58.5
4	4.1	15.6	20.2	60.1
average	4.8(4)	15.1(7)	20.2(8)	60(2)
	Sr	Si	Р	Ν
1	5.3	15.4	21.4	57.8
2	5.6	16.5	20.6	57.3
3	5.6	16.5	22.6	55.3
4	5.6	15.8	19.8	58.8
average	5.5(1)	16.0(5)	21(1)	57(1)
Calc.	5	15	20	60

# **Rietveld refinements**

**Table C7.** Crystallographic data from the Rietveld refinements for  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg, Ca, Sr, Mg<sub>0.94</sub>Ca<sub>0.06</sub>) standard deviations are given in parentheses.

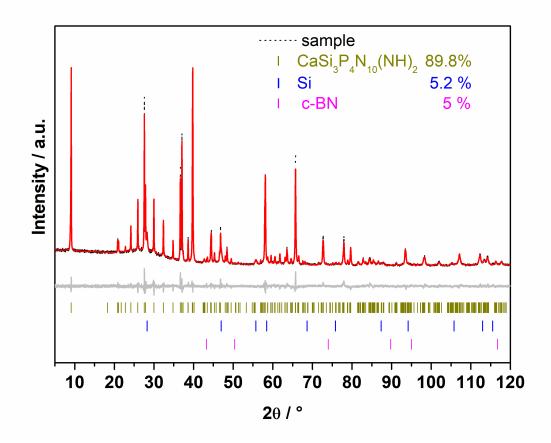
	MgSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	CaSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	SrSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	Mg0.94Ca0.06Si3P4N10(NH)2
crystal system		n	nonoclinic	
space group		<i>C</i> 2	/ <i>m</i> (no. 12)	
lattice	a = 4.86919(6)	a = 4.9184(2)	a = 4.93846(8)	a = 4.8716(1)
parameters /	b = 8.4260(1)	b = 8.5137(3)	b = 8.5496(1)	<i>b</i> = 8.4339(2)
Å,°	c = 9.36667(9)	c = 9.8294(3)	c = 10.1206(1)	c = 9.3928(2)
	$\beta = 99.901(1)$	$\beta=99.447(1)$	$\beta=99.197(1)$	$\beta = 99.873(3)$
radiation $\lambda$ / Å		1.54	-056 (Cu-K <sub>α1</sub> )	
cell volume /	378.560(9)	406.02(3)	421.82(1)	380.02(2)
Å <sup>3</sup>				
$2\theta$ – range/°		5	<2 <i>θ</i> < 120	
refined	54	94	74	54
parameters				
(incl. side				
phases)				
thereof	12	24	24	18
background				
<i>R</i> -values	$R_{\rm p} = 0.0520$	$R_{\rm p} = 0.0480$	$R_{\rm p} = 0.0727$	$R_{\rm p} = 0.0560$
	$wR_{\rm p} = 0.0702$	$wR_{\rm p} = 0.0624$	$wR_{\rm p} = 0.0983$	$wR_{\rm p} = 0.0743$
	$R_{\rm Bragg} = 0.0377$	$R_{\rm Bragg} = 0.0230$	$R_{\rm Bragg} = 0.0568$	$R_{\mathrm{Bragg}} = 0.0337$



**Figure C2.** Rietveld refinement for MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) Xray powder diffraction patterns, positions of Bragg reflections of MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars),  $\beta$ -Si<sub>3</sub>N<sub>4</sub> (vertical blue bars), and difference profile (gray line).

**Table C8** Atomic coordinates, isotropic displacement parameters (/Ų), and occupancy ofMgSi $_3P_4N_{10}(NH)_2$ .

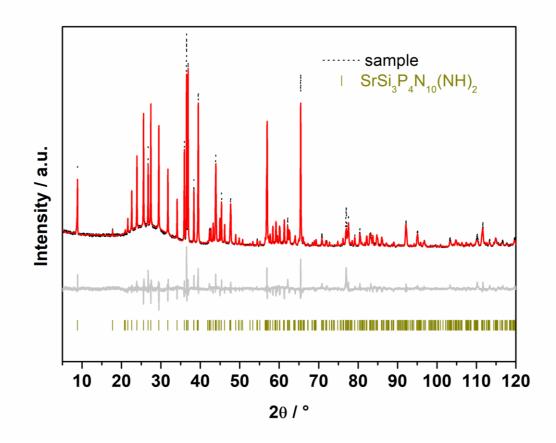
Atom	Wyckoff	X	У	Ζ	Uiso	осс.
Mg1	2 <i>a</i>	0	0	0	0.021(1)	1
P1	8 <i>j</i>	0.0643(6)	0.3332(2)	0.2073(1)	0.0070(9)	1
Si1	4h	0	0.1647(3)	1/2	0.0036(9)	1
Si2	2 <i>d</i>	0	1/2	1/2	0.0036(9)	1
N1	4 <i>i</i>	0.158(1)	0	0.4041(7)	0.002(1)	1
N2	8 <i>j</i>	0.137(1)	0.3304(7)	0.3908(5)	0.002(1)	1
N3	4 <i>i</i>	0.696(2)	0	0.1531(7)	0.002(1)	1
N4	8 <i>j</i>	0.228(1)	0.1812(6)	0.1523(4)	0.002(1)	1
H1	4 <i>i</i>	0.12	0	0.319	0.04	1



**Figure C3.** Rietveld refinement for CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) Xray powder diffraction patterns, positions of Bragg reflections of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars), Si (vertical blue bars), c-BN (vertical violet bars) and difference profile (gray line).

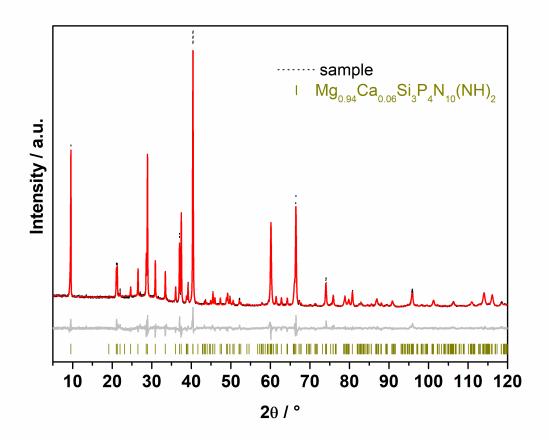
Table C9 Atomic coordinates, isotropic displacement parameters/  ${\rm \AA}^2$ , and occupancy of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>.

Atom	Wyckoff	X	у	Ζ	Uiso	<i>0CC.</i>
Ca1	2 <i>a</i>	0	0	0	0.021(1)	1
P1	8 <i>j</i>	0.0680(6)	0.3322(3)	0.2206(2)	0.013(1)	1
Si1	4h	0	0.1664(2)	1/2	0.009(1)	1
Si2	2 <i>d</i>	0	1/2	1/2	0.009(1)	1
N1	4 <i>i</i>	0.123(2)	0	0.4056(9)	0.010(1)	1
N2	8 <i>j</i>	0.116(1)	0.3312(7)	0.3982(5)	0.010(1)	1
N3	4 <i>i</i>	0.690(2)	0	0.169(1)	0.010(1)	1
N4	8 <i>j</i>	0.252(1)	0.1920(8)	0.1685(7)	0.010(1)	1
H1	4 <i>i</i>	0.12	0	0.319	0.04	1



**Figure C4.** Rietveld refinement for SrSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) Xray powder diffraction patterns, positions of Bragg reflections of SrSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars) and difference profile (gray line).

Atom	Wyckoff	X	у	Ζ	Uiso	OCC.
Sr1	2 <i>a</i>	0	0	0	0.030(1)	1
P1	8 <i>j</i>	0.0758(7)	0.3328(4)	0.2309(1)	0.018(1)	1
Si1	4h	0	0.1659(5)	1/2	0.009(1)	1
Si2	2 <i>d</i>	0	1/2	1/2	0.009(1)	1
N1	4 <i>i</i>	0.130(3)	0	0.404(1)	0.015(1)	1
N2	8 <i>j</i>	0.128(2)	0.333(1)	0.4035(7)	0.015(1)	1
N3	4 <i>i</i>	0.676(2)	0	0.176(1)	0.015(1)	1
N4	8 <i>j</i>	0.268(1)	0.186(1)	0.1779(7)	0.015(1)	1
H1	4 <i>i</i>	0.12	0	0.319	0.04	1



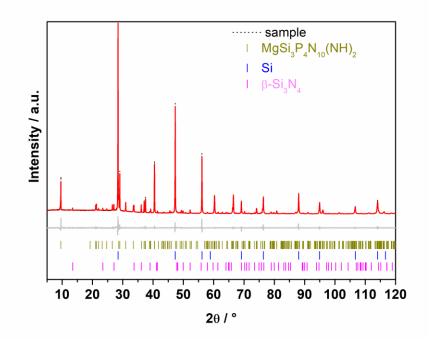
**Figure C5.** Rietveld refinement for Mg<sub>0.94</sub>Ca<sub>0.06</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of Mg<sub>0.94</sub>Ca<sub>0.06</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars) and difference profile (gray line).

Table C11Atomic coordinates, isotropic displacement parameters/ Ų, and occupancy of $Mg_{0.94}Ca_{0.06}Si_3P_4N_{10}(NH)_2$ .

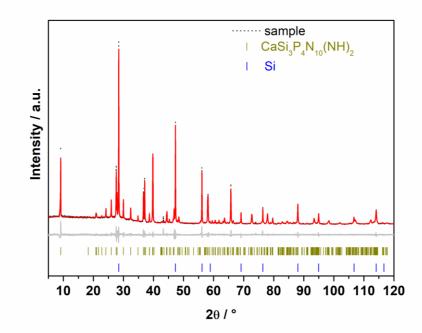
Atom	Wyckoff	X	у	Ζ	Uiso	осс.
Mg1	2 <i>a</i>	0	0	0	0.0264(2)	0.94
Ca1	2 <i>a</i>	0	0	0	0.0264(2)	0.06
P1	8 <i>j</i>	0.0631(7)	0.3330(2)	0.2084(1)	0.012(1)	1
Si1	4 <i>h</i>	0	0.1640(4)	1/2	0.012(1)	1
Si2	2 <i>d</i>	0	1/2	1/2	0.008(1)	1
N1	4 <i>i</i>	0.175(1)	0	0.4031(9)	0.003(1)	1
N2	8 <i>j</i>	0.147(1)	0.3300(8)	0.3931(6)	0.003(1)	1
N3	4 <i>i</i>	0.690(2)	0	0.1509(9)	0.003(1)	1
N4	8 <i>j</i>	0.224(1)	0.1826(7)	0.1543(6)	0.003(1)	1
H1	4 <i>i</i>	0.12	0	0.319	0.04	1

	MgSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	CaSi <sub>3</sub> P <sub>4</sub> N <sub>10</sub> (NH) <sub>2</sub>	$Mg_{0.94}Ca_{0.06}Si_3P_4N_{10}(NH)_2$		
crystal system	monoclinic				
space group	<i>C</i> 2/ <i>m</i> (no. 12)				
lattice parameters / Å,°	a = 4.8682(1)	a = 4.9167(2)	a = 4.8709(2)		
	b = 8.4228(3)	b = 8.5136(3)	b = 8.4296(4)		
	c = 9.3639(2)	c = 9.8198(2)	c = 9.3895(2)		
	$\beta=99.927(3)$	$\beta=99.429(3)$	$\beta = 99.872(4)$		
radiation $\lambda$ / Å	1.54056 (Cu-Kα1)				
cell volume / ų	378.21(2)	405.49(3)	379.83(3)		
$2\theta$ – range/°		5 <2 <i>θ</i> < 120	)		
refined parameters	363	352	352		
(incl. side phases)					
thereof	18	18	18		
background					
<i>R</i> -values	$R_{\rm p} = 0.0517$	$R_{\rm p} = 0.0536$	$R_{\rm p} = 0.0727$		
	$wR_{\rm p} = 0.0713$	$wR_{\rm p} = 0.0751$	$wR_{\rm p} = 0.0983$		

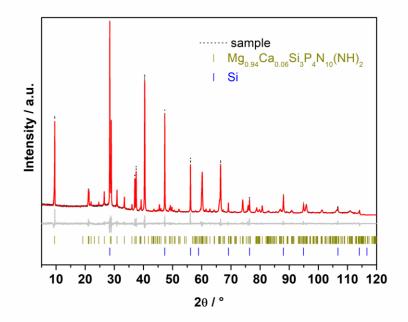
**Table C12.** Crystallographic data from the Pawley Fits for  $AESi_3P_4N_{10}(NH)_2$  (AE = Mg, Ca,  $Mg_{0.94}Ca_{0.06}$ ) standard deviations are given in parentheses.



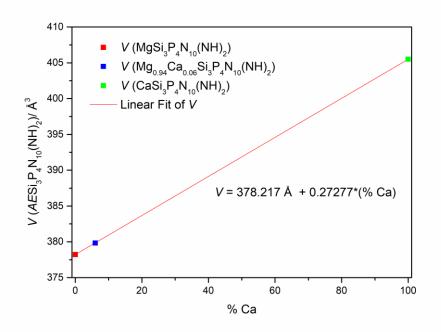
**Figure C6.** Pawley Fit for MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars), Si (vertical blue bars),  $\beta$ -Si3N4 (vertical magenta bars) and difference profile (gray line).



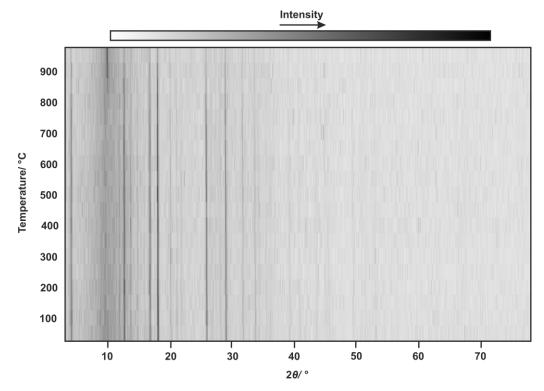
**Figure C7.** Pawley Fit for CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars), Si (vertical blue bars) and difference profile (gray line).



**Figure C8.** Pawley Fit for Mg<sub>0.94</sub>Ca<sub>0.06</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>; observed (black line) and calculated (red line) Xray powder diffraction patterns, positions of Bragg reflections of Mg<sub>0.94</sub>Ca<sub>0.06</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (vertical yellow bars), Si (vertical blue bars) and difference profile (gray line).



**Figure C9**. Linear Fit of the unit cell volumes of MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (red square), Mg<sub>0.94</sub>Ca<sub>0.06</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (blue square) and CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (green square) used for Ca-content estimation of Mg<sub>0.94</sub>Ca<sub>0.06</sub>Si<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>.



#### High-temperature PXRD

**Figure C10.** Temperature-dependent powder X-ray diffraction patterns (Mo-K<sub> $\alpha$ 1</sub> radiation,  $\lambda$  = 0.71073 Å) of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> measured under Ar atmosphere showing a minor expansion of lattice parameters with increasing temperature, and decomposition above 900 °C.

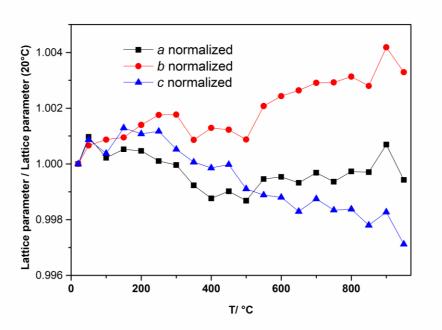


Figure C11. Change of lattice parameters of  $CaSi_3P_4N_{10}(NH)_2$  during heating. Lattice parameters were normalized with respect to their values at 20 °C.

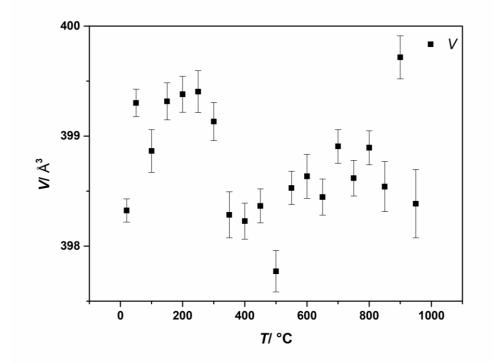


Figure C12. Change of unit cell volume of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> during heating with error bars.

### Solid-state NMR spectra

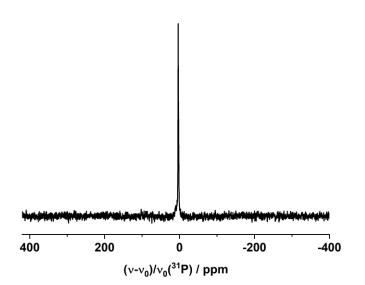


Figure C13. <sup>31</sup>P solid-state MAS NMR spectrum of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>.

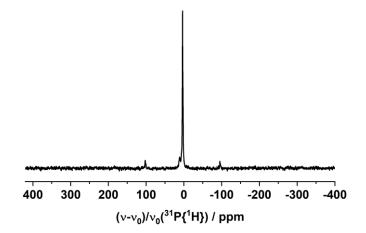


Figure C14.  ${}^{31}P{}^{1}H$  solid-state MAS NMR spectrum of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>. Spinning sidebands are marked with asterisks.

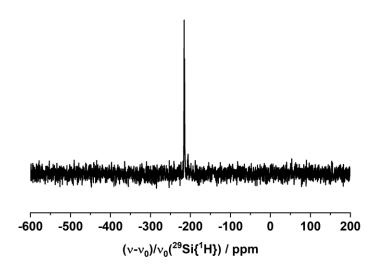


Figure C15. <sup>29</sup>Si{<sup>1</sup>H} solid-state MAS NMR spectrum of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>.

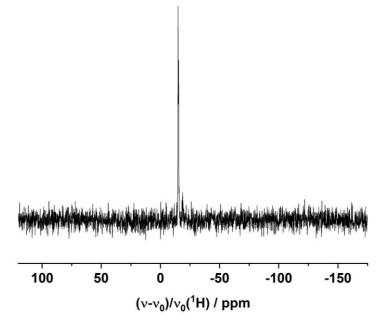


Figure C16. <sup>1</sup>H solid-state MAS NMR spectrum of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>.

Atoms	Wyckoff position	distance/ Å	mean value/ Å
Si1-H1		2.4114	
Si1-H1	4h	2.4114	3.06
Si1-H1		3.7060	
Si1-H1		3.7060	
Si2-H1		2.3583	
Si2-H1	2 <i>d</i>	2.3583	3.07
Si2-H1		3.7744	
Si2-H1		3.7744	

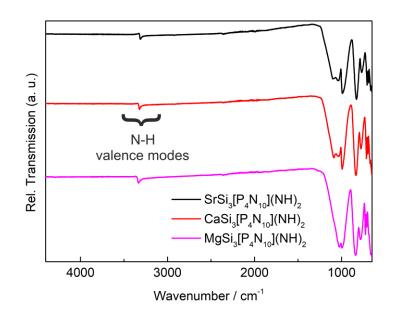
Table C13 Distances (Å) of Si and H in CaSi $_3P_4N_{10}(NH)_2$  from SCXRD Data and the arithmetic mean (Å)

# Bond-valence sum (BVS) Calculations

site	bond-valence sum	expected charge
Ca1	2.50	+2
P1	4.89	+5
Si1	3.86	+4
Si2	3.92	+4
H1	1.57	+1
N1	3.53	-3
N2	2.88	-3
N3	3.01	-3
N4	3.01	-3

**Table C14** Calculated bond valence sums and expected charges for the atom sites in the crystalstructure of  $CaSi_3P_4N_{10}(NH)_2$ .

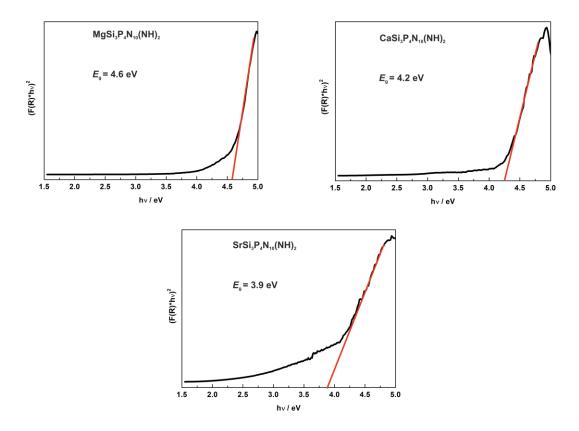
# FTIR spectra



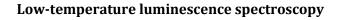
**Figure C17.** FTIR spectra of MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (violet), CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (red) and SrSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (black).

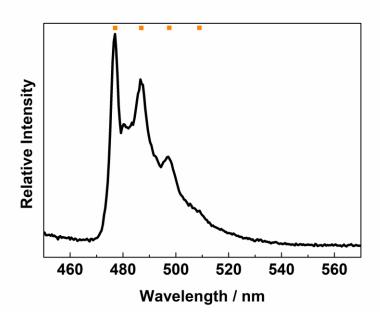
#### UV-Vis spectra

Figure C18: Tauc plots (for UV-Vis spectra of MgSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (top left), CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (top right)



and  $SrSi_3P_4N_{10}(NH)_2$  (bottom). Red lines represent linear regressions from the steep regions to estimate the direct optical band gaps from the inflection points with the abscissa. Data between 4.6 and 4.9 eV (MgSi\_3P\_4N\_{10}(NH)\_2), 4.3 to 4.8 eV (CaSi\_3P\_4N\_{10}(NH)\_2) and 4.15 to 4.85 eV (SrSi\_3P\_4N\_{10}(NH)\_2) were used for linear regressions.





**Figure C19:** Low-temperature emission spectrum of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>, obtained at 6 K (blue line). Red squares indicate the energetic distance of the main phonon mode. The zero-phonon-line is found at ca. 20970 cm<sup>-1</sup>. The Stokes shift was estimated to a value of approx. 258cm<sup>-1</sup> following the equation  $E_{Stokes}$  = (2S–1)ħ $\omega$  with Huang-Rhys-factor S ~0.8 and ħ $\omega$  ~430 cm<sup>-1</sup>.

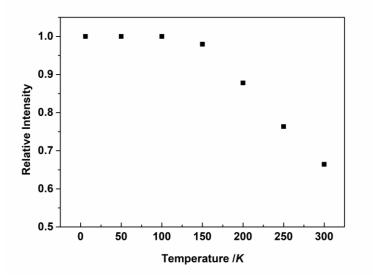


Figure C20: Thermal quenching behavior of CaSi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub>.

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# **D Supporting Information for Chapter 5**

# **Experimental Procedures**

#### Preparation of starting materials

 $P_3N_5$  was prepared following a procedure by Stock et al. by heating  $P_4S_{10}$  (Sigma-Aldrich, 99,99%) in a constant flow of ammonia gas (5.0, Air Liquide). After drying a fused silica tube and a fused silica boat in a tube furnace (at 1273 K) under dynamic vacuum (<10<sup>-3</sup> mbar) the boat was loaded with  $P_4S_{10}$  in a constant flow of Ar. The apparatus was saturated with ammonia for 4 h and the temperature was increased to 1123 K and maintained for four additional hours. Then, the furnace was cooled to room temperature afterwards and heating/cooling ramps were set at 5 K/min each. The product was obtained as an orange powder and purity was confirmed by PXRD as well as CHNS analysis (C 0%, H 0%, N 42.69%, S 0%, expected C 0%, H 0%, N 42.98%, S 0% .<sup>[6]</sup> Cristoballite-type PON was synthesized according to the literature and phase purity was confirmed by PXRD.<sup>[45]</sup> ScN was prepared by nitridation of Sc in flowing N<sub>2</sub> utilizing a radio-frequency furnace according to the literature.<sup>[46]</sup> TiN (Alfa Aesar, 99.7%), TiO<sub>2</sub> (Alfa Aesar, anatase, 99.9%) and NH<sub>4</sub>F (Sigma Aldrich, ≥98%) were used as purchased.

#### High-pressure high-temperature synthesis

The oxonitridophosphates  $TM_5P_{12}N_{26-x}O_x$  (TM = Sc, Ti x = 2, 3) were obtained by highpressure high-temperature synthesis using a 1000 t press with a modified Walker-type multianvil apparatus.  $Sc_5P_{12}N_{23}O_3$  was synthesized from stoichiometric amounts of  $P_3N_5$ , ScN, and PON.  $Ti_5P_{12}N_{24}O_2$  was synthesized from 0.8 equiv.  $TiO_2$  and stoichiometric amounts of TiN and  $P_3N_5$ . 8 wt% NH<sub>4</sub>F were added to the starting mixtures as a mineralizing agent (Table S1). The starting materials were handled in an Ar-filled glovebox (Unilab, MBraun, Garching,  $O_2 < 1$  ppm,  $H_2O < 0.1$  ppm) and ground thoroughly in an agate mortar to minimize external oxygen contamination. The mixture was transferred into a cylindrical crucible made from hexagonal boron nitride (HeBoSint® S100, Henze, Kempten, Germany) closed with a boron nitride cap. The crucible was then placed in an octahedron (MgO with 5 % Cr<sub>2</sub>O<sub>3</sub>, edge length 18 mm, Ceramic Substrates & Components, Isle of Wight, UK), which was drilled through and loaded with different components. First, a ZrO<sub>2</sub> cylinder (Cesima Ceramics, Wust-Fischbeck, Germany) was fixed in the borehole and closed on one side by a Mo plate. Then, a graphite tube (Schunk Kohlenstofftechnik GmbH, Gießen, Germany), a MgO spacer (Cesima Ceramics, Wust-Fischbeck, Germany), and a shorter graphite tube were put into the octahedron. While ZrO<sub>2</sub> served as a thermal insulator, the two graphite tubes were used as electrical resistance furnaces; the combination of two single tubes is necessary to decrease the temperature gradient. After inserting the crucible with the sample, the assembly was completed by adding a second MgO spacer and a further Mo plate to ensure a symmetric setup. To generate approximately hydrostatic conditions, the uniaxial force of the hydraulic press was converted by a Walker-type module with inserted steel wedges, forming a cubic cavity. This cavity was filled with eight Co-doped (7%) WC cubes (Hawedia, Marklkofen, Germany). To enclose the sample octahedron, the edges of the WC cubes were truncated (11 mm edge length). Furthermore, four of the cubes were equipped with a PTFE film (Vitaflon Technische Produkte GmbH, Bad Kreuznach, Germany) for electrical insulation. The remaining cubes were prepared with pyrophyllite gaskets (Ceramic Substrates & Components, Isle of Wight, UK) to prevent the outflow of the MgO octahedron. Bristol board (369 g m<sup>-2</sup>) was used to fix the gaskets. The reactions were carried out under 8 GPa of pressure and 1400 °C.[8,11,12,23]

#### **Powder X-ray diffraction**

Powder X-ray diffraction was carried out on a StadiP diffractometer (STOE & Cie, Darmstadt, Germany) equipped with a MYTHEN 1K detector (Dectris, Baden, Switzerland; angular range  $\Delta 2\theta = 12.5^{\circ}$ ), a Cu source and a Ge(111) monochromator for Cu-K $\alpha_1$  radiation. Samples were filled into glass capillaries with 0.3 mm diameter and 0.01 mm wall thickness (Hilgenberg GmbH, Malsfeld, Germany). Data were recorded in the angular range between  $2\theta = 5-120^{\circ}$  with a step width of 0.015°/step. Diffraction data were analyzed with the TOPAS-Academic V6 software. Structure models were refined with the Rietveld method.

#### Scanning electron microscopy and energy-dispersive X-ray spectroscopy (EDX)

Samples were placed on conducting carbon foil and coated with carbon. EDX spectra were collected using a FEI Helios Nanolab G3 Dual Beam UC (FEI, Hillsboro, OR, USA) with an attached X-Max 80 SDD detector (Oxford Instruments, Abingdon, UK).

# (Scanning) transmission electron microscopy

Samples were ground in absolute ethanol and a drop of the resulting suspension was cast on a TEM grid with holey carbon film (Plano GmbH, Wetzlar, Germany). The grid was mounted on a double-tilt holder and transferred into a Cs DCOR probe-corrected Titan Themis 300 (FEI, USA) TEM equipped with X-FEG, post-column filter (Enfinium ER-799), US1000XP/FT camera system (Gatan, Germany), and a windowless, 4-quadrant Super-X EDX detector. TEM images were recorded using a  $4k \times 4k$  FEI Ceta CMOS camera. The microscope was operated at 300 kV accelerating voltage for SAED and STEM-HAADF and -EDX (convergence angle of 16.6 mrad, 50 µm aperture, detector inner half angle 33 mrad for 245 mm camera length). For evaluation of the TEM data, the following software was used: Digital Micrograph (Fourier filtering of STEM images), ProcessDiffraction7 (geometric calculations for SAED), JEMS (SAED simulations), Velox v2.8 (Thermo Fisher Scientific, USA).

# Single-crystal X-ray diffraction

Single-crystal X-ray diffraction data were collected on a Bruker D8 Venture TXS diffractometer (rotating anode, Mo-K $\alpha$  radiation,  $\lambda = 0.71073$  Å, multilayer monochromator) by combined  $\varphi$ - and  $\omega$ -scans. Indexing, integration, semi-empirical absorption correction (based on equivalent reflections), were performed by the APEX3 software package. Employing SHELX-2018, the structure was solved by direct methods and refined against  $F^2$  by the full-matrix least-squares method. <sup>[16]</sup>

# Bond-valence sum (BVS) Calculations

BVS calculations were performed using the software VaList v4.0.7.<sup>[18]</sup> The structure models of  $TM_5P_{12}N_{26-x}O_x$  (TM = Sc, Ti x = 2, 3) were taken into account with either all anion positions occupied by N or O to determine mixed occupied sites where it was indicated by BVS values. BVS values of cations were determined for the final structure models by taking mixed occupied anionic sites into account.<sup>[17]</sup>

### FTIR spectroscopy

FTIR spectra of the title compounds were collected on a Spectrum BX II spectrometer (PerkinElmer, Waltham, MA, USA) between 600 and 4400 cm<sup>-1</sup> with DuraSampler attenuated total reflectance unit (ATR).

### **Magnetization Measurements**

Magnetization isotherms and susceptibility measurements of a powdered sample was performed with a Quantum Design Inc. Physical Property Measurements System (PPMS) equipped with a vibrating sample magnetometer (VSM) option. Data were collected at temperatures of 2-300 K with field strengths of ±50 kOe with the PPMS MultiVu software package.<sup>[11]</sup> Inverse molar susceptibility data were fitted with a modified Curie-Weiss equation:

$$\frac{1}{\chi_m(T)} = \frac{1}{\frac{1}{\chi_0 + \frac{\mu_B^2}{\chi_0 + \frac{n \cdot (N_A \cdot \frac{\mu_B^2}{3k_B}) \cdot \mu_{eff}}{T - \theta}}} = \frac{1}{\chi_0 + \frac{n \cdot 0.12503776 \cdot \mu_{eff}}{T - \theta}}$$

# UV-Vis spectroscopy

A V-650 UV-Vis spectrophotometer (JASCO, Gross-Umstadt, Germany) equipped with a photomultipler tube detector and a single monochromator with 1200 lines/mm was used for recording UV-Vis spectra in the range of 240 to 800 nm. A deuterium (240–330 nm) and a halogen lamp (330–880 nm) was used with a 2 nm resolution and a scan speed of 400 nm/min controlled by the Spectra Manager II software. Samples were affixed between a fused silica glass slide and a BaSO<sub>4</sub>-coated stamp.

# **Results and Discussion**

# Synthesis

Table D1 Weighed portions of starting materials for the syntheses of  $Sc_5P_{12}N_{23}O_3$  and  $Ti_5P_{12}N_{24}O_2$ .

Title compounds		Starting ma	iterial / mg	
Sc <sub>5</sub> P <sub>12</sub> N <sub>23</sub> O <sub>3</sub>	ScN	P <sub>3</sub> N <sub>5</sub>	PON	NH <sub>4</sub> F
	22.87	37.93	14.19	6
Ti5P12N24O2	TiN	P <sub>3</sub> N <sub>5</sub>	TiO <sub>2</sub>	NH <sub>4</sub> F
	20.43	53.79	5.27	5

#### Structure determination

**Table D2** Atomic coordinates, isotropic displacement parameters/ $Å^2$ , and occupancy of  $Sc_5P_{12}N_{23}O_3$ .

Atom	Wyckoff	X	У	Ζ	$U_{eq}$	Occ.
Sc1	32 <i>g</i>	0.35363(2)	0.06101(2)	0.04356(2)	0.00594(5)	1
Sc2	8 <i>b</i>	0	1/4	1/8	0.02492(13)	1
P1	32 <i>g</i>	0.26625(2)	0.14843(2)	0.16907(2)	0.00410(6)	1
P2	32 <i>g</i>	0.25578(2)	0.30895(2)	0.08368(2)	0.00340(5)	1
Р3	32 <i>g</i>	0.06818(2)	0.34112(2)	0.25058(2)	0.00323(5)	1
N1	32 <i>g</i>	0.34971(7)	0.07463(7)	0.13246(4)	0.00485(14)	1
N2	32 <i>g</i>	0.02045(7)	0.05360(7)	0.12108(4)	0.00525(14)	1
N3	32 <i>g</i>	0.00848(8)	0.08744(7)	0.29230(4)	0.00497(14)	1
N4	32 <i>g</i>	0	1/4	0.03295(5)	0.00475(19)	1
N5/05	32 <i>g</i>	0.17606(7)	0.05919(7)	0.04477(4)	0.00651(14)	0.25/0.75
N6	16 <i>b</i>	0.16331(7)	0.37506(8)	0.04840(4)	0.00554(15)	1
N7	32 <i>g</i>	0.16249(7)	0.28013(7)	0.28644(4)	0.00552(15)	1

Atom	$U_{11}$	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	<i>U</i> <sub>12</sub>
Sc1	0.00665(8)	0.00607(8)	0.00512(8)	-0.00010(6)	-0.00009(6)	-0.00129(6)
Sc2	0.0351(2)	$U_{11}$	0.00456(18)	0	0	0.0213(3)
P1	0.00476(11)	0.00383(11)	0.00371(11)	-0.00010(8)	-0.00038(8)	0.00029(8)
P2	0.00357(11)	0.00353(11)	0.00309(10)	0.00018(8)	0.00012(8)	-0.00029(8)
Р3	0.00342(10)	0.00306(10)	0.00322(10)	-0.00003(8)	-0.00001(8)	-0.00027(8)
N1	0.0054(3)	0.0052(3)	0.0040(3)	0.0000(3)	0.0000(3)	0.0017(3)
N2	0.0053(3)	$U_{11}$	0.0051(3)	0.0007(3)	0.0020(3)	0.0005(3)
N3	0.0043(3)	0.0050(3)	0.0056(3)	0.0018(3)	-0.0003(3)	-0.0002(3)
N4	0.0044(5)	0.0057(5)	0.0042(4)	0	0	-0.0019(4)
N5/05	0.0059(3)	0.0073(3)	0.0062(3)	0.0012(3)	0.0001(3)	-0.0006(3)
N6	0.0052(3)	0.0056(4)	0.0058(4)	-0.0011(3)	-0.0021(3)	0.0008(3)
N7	0.0057(4)	0.0044(3)	0.0065(4)	-0.0006(3)	-0.0020(3)	0.0004(3)

Table D3 Anisotropic displacement parameters for  $Sc_5P_{12}N_{23}O_3$  with standard deviations in parentheses.

**Table D4** Atomic coordinates, isotropic displacement parameters/ $Å^2$ , and occupancy of  $Ti_5P_{12}N_{24}O_2$ .

Atom	Wyckoff	X	у	Ζ	$U_{eq}$	Осс.
Ti1	32 <i>g</i>	0.33980(2)	0.07346(2)	0.04256(2)	0.00686(8)	1
Ti2	8 <i>b</i>	0	1/4	1/8	0.01103(11)	1
P1	32 <i>g</i>	0.26168(3)	0.15620(3)	0.16788(2)	0.00303(8)	1
P2	32 <i>g</i>	0.25231(3)	0.31833(3)	0.08339(2)	0.00280(8)	1
Р3	32 <i>g</i>	0.07509(3)	0.33688(3)	0.25077(2)	0.00279(8)	1
N1	32 <i>g</i>	0.34379(10)	0.07642(9)	0.13085(5)	0.00412(19)	1
N2/02	32 <i>g</i>	0.01550(9)	0.06164(9)	0.12382(4)	0.00480(19)	0.85/0.15
N3	32 <i>g</i>	0.00218(10)	0.08563(9)	0.29203(4)	0.00405(19)	1
N4	16 <i>d</i>	0	1/4	0.03403(7)	0.0041(3)	1
N5/05	32 <i>g</i>	0.17064(9)	0.06275(9)	0.04280(5)	0.00465(19)	0.8/0.2
N6/06	32 <i>g</i>	0.16239(9)	0.38895(9)	0.04614(5)	0.00509(19)	0.85/0.15
N7	32 <i>g</i>	0.16260(9)	0.26818(9)	0.28992(4)	0.00400(18)	1

Atom	$U_{11}$	$U_{22}$	$U_{33}$	$U_{23}$	$U_{13}$	$U_{12}$
Ti1	0.00887(12)	0.00499(12)	0.00673(12)	-0.00132(7)	0.00305(8)	-0.00158(8)
Ti2	0.01242(16)	$U_{11}$	0.0082(2)	0	0	-0.00250(18)
P1	0.00336(14)	0.00297(14)	0.00276(14)	0.00027(10)	-0.00014(10)	0.00033(10)
P2	0.00259(14)	0.00308(15)	0.00274(14)	0.00044(10)	-0.00004(10)	-0.00005(10)
Р3	0.00292(14)	0.00275(14)	0.00270(14)	-0.00001(10)	-0.00004(10)	-0.00030(10)
N1	0.0047(4)	0.0044(4)	0.0033(4)	-0.0003(3)	-0.0001(3)	0.0012(3)
N2/02	0.0050(4)	0.0044(4)	0.0050(4)	-0.0001(3)	0.0010(3)	0.0004(3)
N3	0.0040(4)	0.0043(4)	0.0038(4)	0.0013(3)	-0.0002(3)	-0.0002(3)
N4	0.0050(6)	0.0049(6)	0.0025(6)	0	0	-0.0023(5)
N5/05	0.0048(4)	0.0046(5)	0.0045(4)	0.0004(3)	-0.0002(3)	-0.0003(3)
N6/06	0.0057(4)	0.0045(4)	0.0050(4)	0.0000(3)	-0.0014(3)	0.0002(3)
N7	0.0053(4)	0.0034(4)	0.0033(4)	-0.0001(3)	-0.0017(3)	0.0000(3)

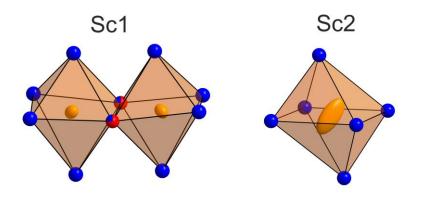
 $\label{eq:table_displacement} \textbf{Table D5} \mbox{ Anisotropic displacement parameters for } Ti_5P_{12}N_{24}O_2 \mbox{ with standard deviations in } parentheses.$ 

Table D6 Interatomic distances (Å) and bond angles (°) in the structure of  $Sc_5P_{12}N_{23}O_3$ .

Sc1	N1	1x	2.1421(9)	N5/05	-P1-	N1	110.18(5)
	N5/05	1x	2.1526(9)	N5/05		N2	111.79(5)
	N5/05	1x	2.1950(9)	N1		N2	110.97(5)
	N7	1x	2.2125(9)	N5/05		N7	105.87(5)
	N3	1x	2.2402(10)	N1		N7	108.16(5)
	N6	1x	2.3108(10)	N2		N7	109.68(5)
Sc2	N4	2x	2.2106(13)	N1	-P2-	N2	108.53(5)
	N2	4x	2.4424(9)	N1		N6	106.04(5)
P1	N5/05	1x	1.5722(9)	N2		N6	109.20(5)
	N1	1x	1.6338(9)	N1		N3	109.40(5)
	N2	1x	1.6372(9)	N2		N3	111.87(5)
	N7	1x	1.6423(9)	N6		N3	111.60(5)
P2	N1	1x	1.6310(9)	N4	-P3-	N7	108.20(4)
	N2	1x	1.6378(9)	N4		N6	104.20(5)
	N6	1x	1.6407(9)	N7		N6	108.92(5)
	N3	1x	1.6434(9)	N4		N3	112.19(4)
Р3	N4	1x	1.6207(7)	N7		N3	109.82(5)
	N7	1x	1.6336(10)	N6		N3	113.24(5)
	N6	1x	1.6351(10)				
	N3	1x	1.6376(9)				

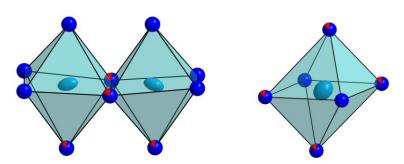
Ti1	N5/05	1x	2.0435(11)	N5/05	-P1-	N2/02	112.79(6)
	N5/05	1x	2.0546(11)	N5/05		N7	106.31(6)
	N1	1x	2.1064(11)	N2/02		N7	110.62(6)
	N3	1x	2.1141(13)	N5/05		N1	108.27(6)
	N7	1x	2.1697(11)	N2/02		N1	110.09(6)
	N6/06	1x	2.2384(11)	N7		N1	108.60(6)
Ti2	N4	1x	2.1691(16)	N1	-P2-	N2/02	109.88(6)
	N4	1x	2.1692(16)	N1		N6/06	107.86(6)
	N2/02	4x	2.2911(11)	N2/02		N6/06	110.24(6)
P1	N5/05	1x	1.6002(12)	N1		N3	108.25(6)
	N2/02	1x	1.6368(11)	N2/02		N3	109.95(6)
	N7	1x	1.6418(11)	N6/06		N3	110.61(6)
	N1	1x	1.6447(12)	N4	-P3-	N6/06	106.18(6)
P2	N1	1x	1.6336(12)	N4		N7	108.96(5)
	N1	1x	1.6354(11)	N6/06		N7	111.50(6)
	N2/02	1x	1.6462(12)	N4		N3	110.86(5)
	N6/06	1x	1.6502(12)	N6/06		N3	110.90(6)
P3	N4	1x	1.6206(9)	N7		N3	108.44(6)
	N6/06	1x	1.6395(12)				
	N7	1x	1.6402(12)				
	N3	1x	1.6516(12)				

Table D7 Interatomic distances (Å) and bond angles (°) in the structure of  $Ti_5P_{12}N_{24}O_2$ .









**Figure D1.** Thermal displacement ellipsoids for the *TM* sites of Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub> (top) and Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub> (bottom). Ellipsoids are displayed at a probability level of 99% (Sc orange, Ti light blue, N dark blue and mixed O/N sites in red/blue.

# EDX measurements

**Table D8**. EDX measurements (several point analyses) for  $TM_5P_{12}N_{26-x}O_x$  (TM = Sc, Ti x = 2, 3) compared to values derived from sum formulas.

atom %	Sc	Р	Ν	0
Crystal no.				
1	11.2	26.3	52.1	10.2
2	12.5	29.0	50.3	8.01
3	11.3	26.4	51.7	10.5
4	11.3	26.5	43.5	18.6
5	12.8	28.3	42.6	16.2
6	9.6	22.9	53.7	13.7
7	11.5	26.9	52.0	9.4
calc.	11.6	27.9	53.5	6.9
	Ti	Р	N	0
1	11.7	33.1	48.8	6.5
2	11.1	31.4	51.5	6.0
3	11.9	31.3	46.3	8.3
4	9.9	26.9	56.7	6.5
5	11.4	30.3	52.2	6.2
Calc.	11.6	27.9	55.8	4.7

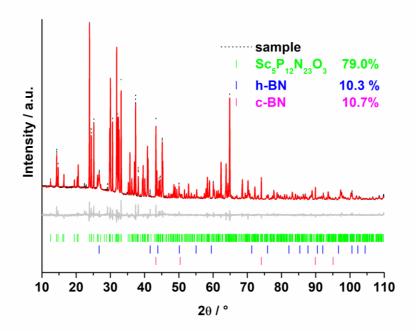
# **Rietveld refinements**

**Table D9**. Crystallographic data from the Rietveld refinements for  $TM_5P_{12}N_{26-x}O_x$  (TM = Sc, Ti x = 2, 3), standard deviations are given in parentheses.

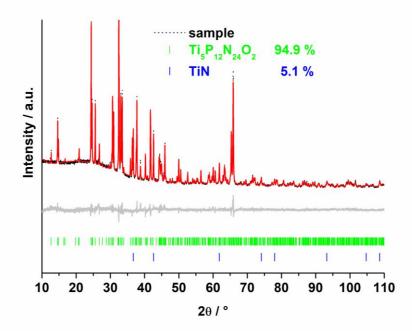
	$Sc_5P_{12}N_{23}O_3^*$	$Ti_5P_{12}N_{24}O_2^{**}$
phase fraction	79(1)	94.9(3)
crystal system	tetrag	gonal
space group	<i>I</i> 4 <sub>1</sub> / <i>acd</i> (1	no. 142)
formula units/ unit cell	8	
lattice parameters / Å	a = 12.36955(7)	a = 12.1133(1)
	c = 24.0304(1)	c = 23.8656(4)
radiation λ / Å	1.54056	$(Cu-K_{\alpha 1})$
cell volume / ų	3676.80(5)	3501.8(1)
2 $\theta$ - range/°	5 <2 <i>0</i> ·	< 120
refined parameters (incl. side phases)	29	26
background function	Shifted Chebyshev	(12 parameters)
<i>R</i> -values	$R_{\rm p} = 0.0440$	$R_{\rm p} = 0.0465$
	$wR_{\rm p} = 0.0612$	$wR_{\rm p} = 0.0662$
	$R_{\rm Bragg} = 0.0273$	$R_{\mathrm{Bragg}} = 0.0294$

\*) BN from crucible as side phases

\*\*) TiN (starting material) as side phase



**Figure D2** Rietveld refinement for Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub>; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub> (vertical green bars), h-BN (vertical blue bars), c-BN (vertical violet bars), and difference profile (gray line).



**Figure D3.** Rietveld refinement for Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub> observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of Ti<sub>5</sub>P<sub>12</sub>N<sub>24</sub>O<sub>2</sub> (vertical green bars), TiN (vertical blue bars), and difference profile (gray line).

### **BVS calculations**

 $\label{eq:constraint} \textbf{Table D10.} Anion BVS for Sc_5P_{12}N_{23}O_3. Assumed mixed occupied sites are highlighted in green.$ 

N1	N2	N3	N4	N5	N6	N7
3.14	2.77	2.93	3.09	2.68	2.90	3.03
01	02	03	04	05	06	07
2.42	2.15	2.27	2.40	2.01	2.25	2.34

**Table D11.** Cation BVS for  $Sc_5P_{12}N_{23}O_3$ . Assumed mixed occupied anion sites were considered for calculations.

	Sc1	Sc2	P1	P2	P3
calc.	3.17	2.42	4.89	4.89	4.96
expected	3	3	5	5	5

Table D12. Anion BVS for  $Ti_5P_{12}N_{24}O_2$ . Assumed mixed occupied sites are highlighted in green.

N1	N2	N3	N4	N5	N6	N7
3.05	2.86	2.99	3.09	2.85	2.89	2.98
01	02	03	04	05	06	07
2.39	2.20	2.33	2.39	2.16	2.25	2.32

 $\label{eq:table_$ 

calculations.

	Ti1	Ti2	P1	P2	Р3
calc.	3.65	2.64	4.91	4.79	4.87
expected	3	4	5	5	5

# FTIR spectra

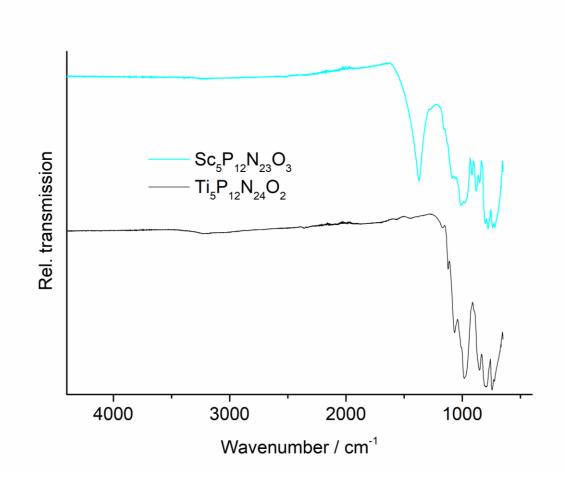
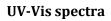
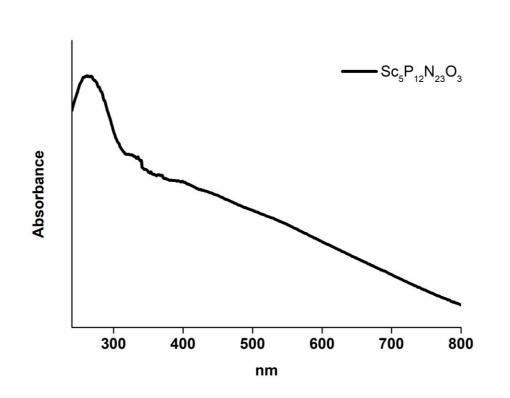


Figure D4. FTIR spectra of  $Sc_5P_{12}N_{23}O_3$  (blue) and  $Ti_5P_{12}N_{24}O_2$  (black).





**Figure D5:** UV-Vis spectrum of Sc<sub>5</sub>P<sub>12</sub>N<sub>23</sub>O<sub>3</sub>.

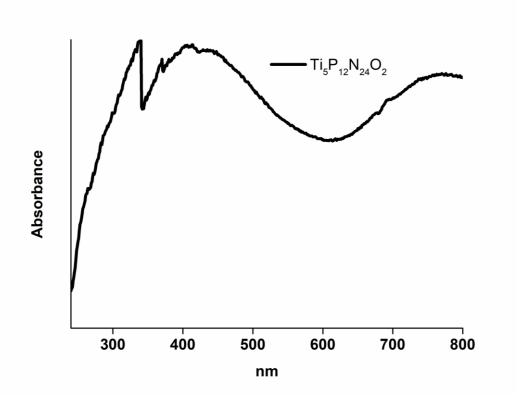
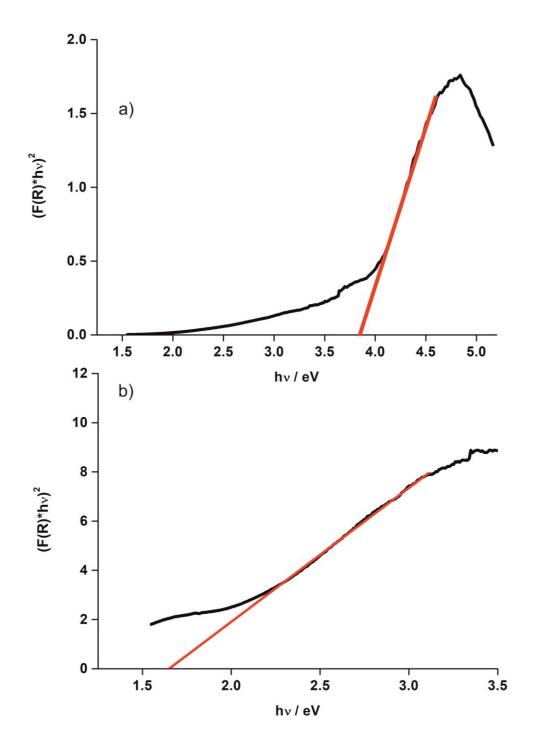


Figure D6: UV-Vis spectrum of  $Ti_5P_{12}N_{24}O_2$ .



**Figure D7:** Tauc plots for  $Sc_5P_{12}N_{23}O_3$  (top) and  $Ti_5P_{12}N_{24}O_2$  (bottom). Red lines represent linear regressions from the steep regions to estimate the optical band gaps from the inflection points with the abscissa. Data between 2.1 and 3.1 eV ( $Ti_5P_{12}N_{24}O_2$ ) and 4.0 to 4.5 eV ( $Sc_5P_{12}N_{23}O_3$ ) were used for linear regressions.

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# **E Supporting Information for Chapter 6**

# **Experimental Procedures**

# Preparation of starting materials

 $P_3N_5$  was prepared following a procedure by Stock et al.<sup>[1]</sup> The product was obtained as an orange powder. Purity was confirmed by PXRD as well as CHNS analysis (C 0%, H 0%, N 42.69%, S 0% expected C 0%, H 0%, N 42.98%, S 0%). The azide NH<sub>4</sub>N<sub>3</sub> was prepared according to literature.<sup>[2]</sup> TiN (Alfa Aesar, 99.7%) and NH<sub>4</sub>F (Sigma Aldrich, ≥98%) were used as purchased.

# High-pressure high-temperature synthesis

The nitridophosphate  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> was obtained by high-pressure high-temperature synthesis using a 1000 t press with a modified Walker-type multianvil apparatus.<sup>[3-7]</sup> All products were synthesized from stoichiometric amounts of P<sub>3</sub>N<sub>5</sub>, TiN, NH<sub>4</sub>N<sub>3</sub>, and additional 10 wt% NH<sub>4</sub>F (Table E1). To prevent oxygen contamination, starting materials were handled in an Ar-filled glovebox (Unilab, MBraun, Garching, O<sub>2</sub><1 ppm, H<sub>2</sub>O<0.1 ppm) and ground thoroughly in an agate mortar. Details on the preparation and handling of the 1000 t Walker-type multianvil press are described in the literature.<sup>[8-10]</sup> The reaction was carried out at 8 GPa and 1400°C.

# Annealing of TiP<sub>4</sub>N<sub>8</sub>

 $\beta$ -TiP<sub>4</sub>N<sub>8</sub> was annealed in sealed silica ampoules in a tube furnace at 850 °C for 20 h with heating and cooling rates of 5 °C/ min to obtain bulk  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>.

# Powder X-ray diffraction

Powder X-ray diffraction was carried out on a StadiP diffractometer (STOE & Cie, Darmstadt, Germany) equipped with an MYTHEN 1K detector (Dectris, Baden, Switzerland; angular range  $\Delta 2\theta = 12.5^{\circ}$ ), a Cu source and a Ge(111) monochromator for Cu-K $\alpha_1$  radiation. Samples were filled into glass capillaries with 0.3 mm diameter and 0.01 mm wall thickness (Hilgenberg GmbH, Malsfeld, Germany). Data were recorded in the angular range  $5^{\circ} < 2\theta <$ 

120°. Diffraction patterns were analyzed with the TOPAS-Academic V6 software, where phase fractions were approximately quantified with the Rietveld method.

Temperature-dependent PXRD patterns were recorded on a StadiP diffractometer (STOE & Cie, Darmstadt, Germany) in the range  $2^{\circ} < 2\theta < 75^{\circ}$  and Mo-K $\alpha_1$  radiation ( $\lambda = 0.709300$  Å). An image-plate position-sensitive detector was used for X-ray detection. Diffraction patterns were collected in steps of 50°C up to 700 °C for two cycles. Samples were filled into silica capillaries with 0.4 mm diameter and 0.01 mm wall thickness (Hilgenberg GmbH, Malsfeld, Germany). Lattice parameters were extracted by means of Pawley fits.

### Scanning electron microscopy and energy-dispersive X-ray spectroscopy (EDX)

Samples were placed on conducting carbon foil and coated with carbon. EDX spectra were collected using an FEI Helios Nanolab G3 Dual Beam UC (FEI, Hillsboro, OR, USA) with an attached X-Max 80 SDD detector (Oxford Instruments, Abingdon, UK).

### (Scanning) transmission electron microscopy (TEM, STEM)

Samples were ground in absolute ethanol and a drop of the resulting suspension was cast on a TEM grid with lacey carbon film (Plano GmbH, Wetzlar, Germany). The grid was mounted on a double-tilt low background holder and transferred into a Cs DCOR probe-corrected Titan Themis 300 (FEI, USA) TEM equipped with X-FEG, post-column filter (Enfinium ER-799), US1000XP/FT camera system (Gatan, Germany), and a windowless, 4-quadrant Super-X EDX detector (FEI, USA) TEM images were recorded using a  $4k \times 4k$  FEI Ceta CMOS camera. The microscope was operated at 300 kV accelerating voltage for SAED, STEM-HAADF and EDX (semiconvergence angle 16.6 mrad, 50  $\mu$ m aperture, HAADF detector inner half angle 33 mrad for 245 mm camera length). For evaluation of the TEM data, the following software was used: Digital Micrograph (Fourier filtering of STEM images), ProcessDiffraction7 (geometric calculations for SAED), JEMS (SAED simulations).

# Single-crystal X-ray diffraction

Single-crystal X-ray diffraction data at ambient conditions were collected on a Bruker D8 Venture TXS diffractometer (rotating anode, Mo-K<sub> $\alpha$ </sub> radiation,  $\lambda = 0.71073$  Å, multilayer

monochromator) by combined  $\varphi$ - and  $\omega$ -scans. Indexing, integration, semi-empirical absorption correction (based on equivalent reflections), were performed by the APEX3 software package. Employing SHELX-2018, the structure was solved by direct methods and refined against F<sup>2</sup> by the full-matrix least-squares method.<sup>[11]</sup>

Temperature-dependent high-intensity synchrotron diffraction data of a TiP<sub>4</sub>N<sub>8</sub> single crystal were obtained at beamline P24, DESY, Hamburg, ( $\lambda = 0.49920$  Å) with a Pilatus 1M CdTe. Single crystals of disordered  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> were mounted and sealed inside silica glass capillaries and held in place with silica glass fibers. Integration was done with CrysAlis Pro and semiempirical absorption correction with SADABS.<sup>[12,13]</sup> Solution and refinement of the average structure were done with SHELX-2018.

In some cases, separate scale factors for reflections with h = 3n and  $h \neq 3n$  were used (HKLF 5 format).

### Bond-valence sum (BVS) Calculations

BVS calculations were performed using the software VaList v4.0.7.<sup>[14,15]</sup>

#### IR spectroscopy

FTIR spectra of the title compounds were collected on a Spectrum BX II spectrometer (PerkinElmer, Waltham, MA, USA) between 600 and 4400 cm<sup>-1</sup> with DuraSampler attenuated total reflectance unit (ATR).

#### **UV-Vis spectroscopy**

A V-650 UV-Vis spectrophotometer (JASCO, Gross-Umstadt, Germany) equipped with a photomultiplier tube detector and a single monochromator with 1200 lines/mm was used for recording UV-Vis spectra in the range of 240 to 800 nm. A deuterium (240–330 nm) and a halogen lamp (330–880 nm) were used with a 2 nm resolution and a scan speed of 400 nm/min controlled by the Spectra Manager II software. Samples were fixed between a fused silica glass slide and a BaSO<sub>4</sub>-coated stamp. Pseudo-absorption spectra were calculated by the Kubelka-Munk function  $F(R) = (1-R)^2/2R$  where *R* is reflectance. In the Tauc plots, *hv* was plotted against  $[F(R)hv]^{1/n}$  with n = -1/2 for indirect bandgaps.<sup>[16]</sup>

### **DFT calculations**

The Vienna ab initio simulation package (VASP) based on density functional theory (DFT) and plane wave basis sets<sup>[17,18]</sup> was used for the electronic structure calculations The projectoraugmented waves (PAW)<sup>[19]</sup> approach was applied, and contributions of correlation and exchange were treated in the generalized-gradient approximation (GGA)<sup>[20]</sup> using PBE<sup>[20]</sup> and SCAN<sup>[21]</sup> potentials. Dispersion corrections with the Grimme D2-Method were applied.<sup>[22]</sup> The Brillouin zones were sampled with either a 4x20x11 ( $\beta$ -TiP<sub>4</sub>N<sub>8</sub>) or 12x20x11 ( $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>)  $\Gamma$ -centered Monkhorst Pack *k*-point mesh and a 500 eV plane wave energy cutoff. Both structures were optimized until forces between atoms are below 5×10-5 eVÅ-1 and total energy changes below  $10^{-8}$  eV. For E(V) calculations the cell volumes were varied  $\pm 5\%$  around the equilibrium and the results evaluated with the universal Vinet equation of state (EOS). Band structures and electronic density of states are shown in Figure E12.Phonon dispersions and phonon density-of states (Figure E13) were calculated from forces acting on displaced atoms in 2x2x2 supercells using PHONOPY.<sup>[23]</sup> Thermodynamic properties were obtained within the quasi-harmonic approximation (QHA)<sup>[24]</sup> by calculating the phonon states at nine volumes ( $\pm$  5% around the equilibrium volume) and findig the Gibbs free energy G(T,p) at the minima of the function  $[U(V) + F_{phonon}(V,T) + pV]$  at eeach temperature as shown exemplarily for zero pressure in Figure E14. The energy barrier between  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> was calculated using the climbing nudged elasting band (CI-NEB) method<sup>[25]</sup> with 16 images between the initial and final states. Linear extrapolations between the initial and final structures were used as starting parameters, respectively.

# **Results and Discussion**

# Synthesis

Table E11 Weighed portions of starting materials for the syntheses of  $TiP_4N_8$ .

Compound	d	Sta	rting material	
TiP4N8	TiN	P <sub>3</sub> N <sub>5</sub>	NH <sub>4</sub> N <sub>3</sub>	NH <sub>4</sub> F
	16.4 mg	57.4 mg	2.1 mg	6 mg
	0.26 mmol	0.35 mmol	0.03 mmol	0.16 mmol

### Structure representations

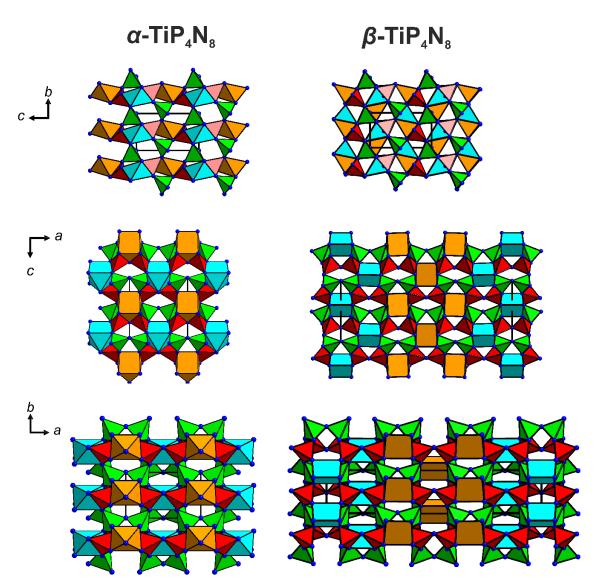


Figure E1. Structure projetions of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> along [100], [010] and [001]

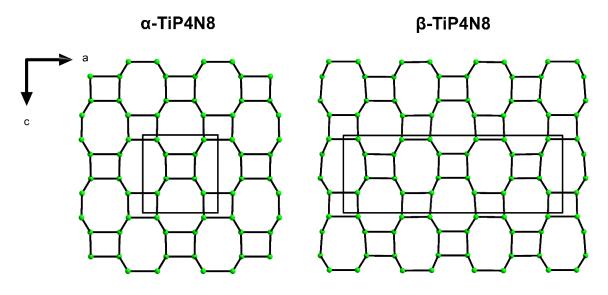


Figure E2. Nodal representations of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> along [010]

# Structure determination

 $\label{eq:table_$ 

compound	α-TiP <sub>4</sub> N <sub>8</sub>	β-TiP <sub>4</sub> N <sub>8</sub>	disordered	HT-β-TiP <sub>4</sub> N <sub>8</sub>		
			TiP <sub>4</sub> N <sub>8</sub>			
temperature/K	296(2)	298(2)	293	873		
molar mass /g·mol <sup>-1</sup>		283.	.86			
crystal system		orthorh	ombic			
space group		<i>Pmn</i> 2 <sub>1</sub> (2	no. 31)			
lattice parameters / Å,	a = 7.6065(2)	a = 22.9196(5)	a = 22.907(5)	a = 22.962(5)		
0						
	b = 4.63320(10)	b = 4.58800(10)	b = 4.5827(9)	b = 4.6045(9)		
	c = 7.8601(3)	c = 8.0970(2)	c = 8.0988(16)	c = 8.0982(16)		
cell volume / ų	277.009(14)	851.44(3)	850.2(3)	856.2(3)		
formula units/ per	2	6	6	6		
unit cell						
density / g⋅cm <sup>3</sup>	3.403	3.322	3.327	3.303		
$\mu$ / mm <sup>-1</sup>	2.640	2.577	0.939	0.939		
temperature / K	296(2)	298(2)	293(3)	873(5)		
absorption correction	semiempirical					
radiation	Mo- $K_{\alpha}$ ( $\lambda$ =	0.71073 Å)	synchrotron, DE	SY P24		
			$(\lambda=0.4992$	Å)		
F(000)	276	828	828	828		
heta range / °	$3.7 \leq \theta \leq 36.3$	$1.7 \leq \theta \leq 34.9$	$1.9 \leq \theta \leq 23.3$	$1.9 \leq \theta \leq 23.3$		
$\sin \theta \cdot \lambda^{-1} / \text{\AA}^{-1}$	0.83	0.80	0.79	0.79		
total no. of reflections	7500	9236	11025	33096		
Independent	1357/1416	6268/6574	3055/3136	23542/23542		
reflections [I $\geq 2\sigma$ (I) /						
all]						
$R_\sigma/R_{ m int}$	0.0317 / 0.0428	0.0133/ -*	0.0217/0.0182	0.0342/-*		
refined parameters /	65/1	179/1	193/2	184/1		
restraints	1 001	1 1 2 6	1 1 0 0	1		
Goodness of fit	1.091	1.126	1.102	1.081		
<i>R</i> -values [I $\geq 2\sigma$ (I)]	R1 = 0.0255;	R1 = 0.0220;	R1 = 0.0244	R1 = 0.0395		
	wR2 = 0.0617	wR2 = 0.0684	wR2 = 0.0721	wR2 = 0.1217		
<i>R</i> -values (all data)	R1 = 0.0277;	R1 = 0.0232;	R1 = 0.0247	R1 = 0.0419		
• • • • • • • • •	wR2 = 0.0631	wR2 = 0.0692	wR2 = 0.0730	wR2 = 0.1256		
$\Delta ho_{ m max}$ , $\Delta ho_{ m min}$ / e·Å <sup>-3</sup>	0.494, -0.643	0.720, -0.467	0.706/ -1.482	1.725/ -1.075		

standard deviations are given in parentheses.

\*separate scale factors were used (HKLF 5 format)

Atom	Wyckoff position	X	у	Ζ	$U_{eq}$	<i>0CC.</i>
Ti1	2a	0	0.61255(12)	0.42990(9)	0.00623(12)	1
P1	4b	0.30454(8)	0.41321(13)	0.24765(9)	0.00409(12)	1
P2	4b	0.31171(8)	0.09790(13)	0.56012(8)	0.00397(12)	1
N1	4b	0.3147(3)	0.1809(5)	0.0931(3)	0.0059(4)	1
N2	4 <i>b</i>	0.1736(3)	0.6911(5)	0.2278(3)	0.0050(3)	1
N3	4 <i>b</i>	0.2060(3)	0.2858(4)	0.4157(3)	0.0058(3)	1
N4	2 <i>a</i>	0	0.0007(7)	0.0000(4)	0.0075(5)	1
N5	2 <i>a</i>	0	0.4548(7)	0.7394(5)	0.0092(5)	1

Table E3 Atomic coordinates, isotropic displacement parameters (/Å<sup>2</sup>), and site occupancies of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>.

**Table E4** Anisotropic displacement parameters ( $U_{ij}$  in Å<sup>2</sup>) for  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> with estimated standard deviations in parentheses. The anisotropic displacement factor is expressed as  $\exp[-2\pi^2 (U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}l^2c^{*2} + U_{12}hka^*b^* + U_{13}hla^*c^* + U_{23}klb^*c^*)$ .

Atom	<i>U</i> <sub>11</sub>	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	<i>U</i> <sub>12</sub>
Ti1	0.0056(3)	0.0073(2)	0.0058(2)	0.0002(2)	0	0
P1	0.0036(2)	0.0045(2)	0.0041(2)	-0.0005(2)	0.0003(2)	0.00020(18)
P2	0.0041(3)	0.0038(2)	0.0040(2)	-0.0001(2)	0.00008(19)	-
						0.00024(18)
N1	0.0047(9)	0.0065(8)	0.0065(9)	-0.0011(6)	0.0005(6)	-0.0015(7)
N2	0.0049(8)	0.0052(8)	0.0049(9)	0.0018(7)	0.0006(7)	0.0013(6)
N3	0.0065(8)	0.0061(8)	0.0048(8)	0.0012(7)	0.0009(7)	0.0007(6)
N4	0.0037(11)	0.0096(13)	0.0093(12)	0.0039(10)	0	0
N5	0.0055(12)	0.0054(11)	0.0167(14)	0.0024(11)	0	0

Ti1	N2	2x	2.098(2)	N5	P1	N1	100.23(1
Ti1	N1	2x	2.132(2)	N5	P1	N3	126.54(1
Ti1	N3	2x	2.182(2)	N1	P1	N3	112.83(1
T1	N5	1x	2.540(4)	N5	P1	N2	105.06(1
P1	N5		1.6089(14)	N1	P1	N2	118.62(1
P1	N1		1.625(2)	N3	P1	N2	94.74(1)
P1	N3		1.629(2)	N4	P2	N1	110.73(1
P1	N2		1.636(2)	N4	P2	N3	113.05(1
P2	N4		1.5758(15)	N1	P2	N3	103.95(1
P2	N1		1.631(2)	N4	P2	N2	110.52(1
P2	N3		1.641(2)	N1	P2	N2	112.56(1
P2	N2		1.644(2)	N3	P2	N2	105.81(1

Table E6 Atomic coordinate	s, isotropic displacement parameters	(/Å <sup>2</sup> ), and site occupancies of $\beta$ -
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TiP <sub>4</sub> N <sub>8</sub> .						
Atom	Wyckoff	X	У	Ζ	Ueq	<i>0CC.</i>
	position					
Ti1	2 <i>a</i>	0	0.88420(17)	0.39418(13)	0.0161(2)	1
Ti2	4 <i>b</i>	0.33332(3)	0.62120(15)	0.40469(12)	0.0206(2)	1
P1	4 <i>b</i>	0.43551(4)	0.39932(19)	0.23748(12)	0.0110(2)	1
P2	4 <i>b</i>	0.43389(4)	0.09922(19)	0.54809(13)	0.0116(2)	1
Р3	4 <i>b</i>	0.09807(4)	0.40398(18)	0.24156(12)	0.0111(2)	1
P4	4 <i>b</i>	0.10473(4)	0.10293(19)	0.55210(12)	0.0119(2)	1
P5	4 <i>b</i>	0.22996(4)	0.40090(19)	0.24674(13)	0.0124(2)	1
P6	4 <i>b</i>	0.23394(4)	0.09939(19)	0.55727(13)	0.0126(2)	1
N1	4 <i>b</i>	0.43858(14)	0.1694(8)	0.0878(6)	0.0172(6)	1
N2	4 <i>b</i>	0.05888(14)	0.6935(7)	0.2108(5)	0.0148(6)	1
N3	4 <i>b</i>	0.10349(14)	0.1949(7)	0.0813(5)	0.0152(6)	1
N4	4 <i>b</i>	0.10919(14)	0.3320(7)	0.7053(5)	0.0152(6)	1
N5	4 <i>b</i>	0.22510(14)	0.1803(7)	0.0925(6)	0.0167(6)	1
N6	4 <i>b</i>	0.27091(14)	0.6830(7)	0.2172(5)	0.0153(6)	1
N7	4 <i>b</i>	0.33060(13)	0.0094(9)	0.0123(5)	0.0232(8)	1
N8	2 <i>a</i>	0	0.4769(9)	0.7734(6)	0.0164(8)	1
N9	4 <i>b</i>	0.16395(11)	0.5218(8)	0.2821(6)	0.0216(7)	1
N10	4 <i>b</i>	0.06771(12)	0.2294(7)	0.3937(4)	0.0144(6)	1
N11	4 <i>b</i>	0.40100(14)	0.2628(7)	0.3960(4)	0.0172(6)	1
N12	4 <i>b</i>	0.26554(13)	0.2686(8)	0.4033(4)	0.0179(6)	1
N13	2 <i>a</i>	0	0.0044(11)	0.0017(7)	0.0218(10)	1

**Table E7** Anisotropic displacement parameters ( $U_{ij}$  in Å<sup>2</sup>) for β-TiP<sub>4</sub>N<sub>8</sub> with estimated standard deviations in parentheses. The anisotropic displacement factor is expressed as exp[ $-2\pi^2$  ( $U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}l^2c^{*2} + U_{12}hka^*b^* + U_{13}hla^*c^* + U_{23}klb^*c^*$ ).

Atom	<i>U</i> <sub>11</sub>	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	<i>U</i> <sub>12</sub>
Ti1	0.0150(4)	0.0217(4)	0.0116(4)	-0.0008(4)	0	0
Ti2	0.0136(3)	0.0264(4)	0.0217(4)	-0.0029(3)	-0.00007(16)	-0.00033(17)
P1	0.0107(4)	0.0105(4)	0.0119(4)	-0.0018(3)	0.0010(3)	0.0008(3)
P2	0.0132(4)	0.0106(4)	0.0111(5)	0.0013(3)	-0.0014(3)	0.0000(3)
Р3	0.0095(3)	0.0099(4)	0.0138(4)	-0.0012(3)	-0.0004(4)	0.0008(3)
P4	0.0114(4)	0.0115(4)	0.0130(5)	0.0009(3)	0.0006(4)	-0.0002(3)
P5	0.0110(4)	0.0108(4)	0.0153(5)	-0.0013(3)	-0.0012(4)	-0.0006(3)
P6	0.0126(4)	0.0107(4)	0.0144(5)	0.0012(4)	0.0005(4)	0.0003(3)
N1	0.0127(11)	0.0178(13)	0.0210(16)	-0.0054(12)	0.0001(12)	-0.0060(10)
N2	0.0176(12)	0.0123(12)	0.0147(13)	0.0039(11)	0.0054(12)	0.0041(10)
N3	0.0136(10)	0.0143(12)	0.0177(13)	-0.0021(11)	-0.0012(12)	-0.0036(11)
N4	0.0167(13)	0.0143(11)	0.0147(13)	-0.0044(10)	-0.0056(12)	0.0025(11)
N5	0.0152(11)	0.0156(12)	0.0194(16)	-0.0044(11)	0.0005(12)	0.0053(11)
N6	0.0168(13)	0.0137(12)	0.0154(14)	0.0053(11)	-0.0056(12)	-0.0039(11)
N7	0.0115(12)	0.0233(17)	0.035(2)	0.0106(16)	-0.0005(13)	-0.0002(10)
N8	0.0096(15)	0.0162(17)	0.023(2)	0.0074(16)	0	0
N9	0.0103(12)	0.0164(14)	0.038(2)	-0.0095(15)	-0.0010(13)	0.0001(9)
N10	0.0142(13)	0.0149(11)	0.0141(12)	0.0047(16)	0.0035(10)	-0.0014(9)
N11	0.0146(12)	0.0207(13)	0.0163(13)	0.0012(17)	0.0058(10)	0.0018(9)
N12	0.0149(13)	0.0230(14)	0.0159(15)	0.0046(17)	-0.0060(10)	-0.0028(10)
N13	0.0111(17)	0.022(2)	0.032(3)	0.0103(19)	0	0

Ti1	N1	2x	2.105(3)	N8	P1	N1	109.75(14)
Ti1	N2	2x	2.216(3)	N8	P1	N4	110.37(15)
Ti1	N10	2x	2.219(2)	N1	P1	N4	113.19(16)
Ti2	N4	1x	2.086(3)	N8	P1	N11	115.55(17)
Ti2	N6	1x	2.108(3)	N1	P1	N11	111.69(12)
Ti2	N3	1x	2.200(3)	N4	P1	N11	95.75(13)
Ti2	N5	1x	2.201(3)	N11	P2	N2	112.40(12)
Ti2	N12	1x	2.229(2)	N11	P2	N3	106.08(13)
Ti2	N11	1x	2.241(2)	N2	P2	N3	114.16(17)
P1	N8	1x	1.5998(15)	N11	P2	N13	113.16(17)
P1	N1	1x	1.617(3)	N2	P2	N13	105.28(15)
P1	N4	1x	1.630(2)	N3	P2	N13	105.74(15)
P1	N11	1x	1.644(3)	N10	Р3	N2	107.35(12)
P2	N11	1x	1.617(3)	N10	Р3	N3	112.57(12)
P2	N2	1x	1.618(3)	N2	Р3	N3	114.08(17)
P2	N3	1x	1.625(2)	N10	Р3	N9	112.78(13)
P2	N13	1x	1.6608(16)	N2	Р3	N9	104.78(12)
Р3	N10	1x	1.622(2)	N3	Р3	N9	105.13(12)
Р3	N2	1x	1.623(2)	N7	P4	N1	110.98(13)
Р3	N3	1x	1.625(3)	N7	P4	N4	109.01(12)
Р3	N9	1x	1.650(2)	N1	P4	N4	113.24(17)
P4	N7	1x	1.603(2)	N7	P4	N10	115.92(13)
P4	N1	1x	1.623(2)	N1	P4	N10	95.29(12)
P4	N4	1x	1.624(3)	N4	P4	N10	111.95(12)
P4	N10	1x	1.642(2)	N5	P5	N6	116.33(14)
P5	N5	1x	1.617(3)	N5	P5	N12	116.08(12)
P5	N6	1x	1.623(2)	N6	P5	N12	97.20(12)
P5	N12	1x	1.629(3)	N5	P5	N9	105.03(12)
P5	N9	1x	1.651(2)	N6	P5	N9	106.24(12)
P6	N7	1x	1.608(2)	N12	P5	N9	115.90(13)
P6	N5	1x	1.630(2)	N7	P6	N5	109.81(13)
P6	N12	1x	1.634(3)	N7	P6	N12	113.08(13)
P6	N6	1x	1.634(3)	N5	P6	N12	104.35(12)
				N7	P6	N6	109.08(12)
				N5	P6	N6	111.55(13)
				N12	P6	N6	108.93(12)

**Table E8** Interatomic distances (/Å) and bond angles (/°) in the structure of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>.

Atom	Wyckoff	X	у	Ζ	$U_{eq}$	осс.
	position					
Ti1	2 <i>a</i>	0	0.88211(9)	0.39689(12)	0.00513(10)	1
Ti2	4 <i>b</i>	0.33322(2)	0.62344(7)	0.40121(10)	0.00587(9)	0.9077(14)
Ti2a	4 <i>b</i>	0.33324(13)	0.8723(8)	0.3901(6)	0.00587(9)	0.0923(14)
P1	4 <i>b</i>	0.43596(2)	0.40265(10)	0.23892(8)	0.00383(11)	1
P2	4 <i>b</i>	0.43308(2)	0.10273(10)	0.54869(8)	0.00416(13)	1
Р3	4 <i>b</i>	0.09742(2)	0.39922(9)	0.24256(7)	0.00387(11)	1
P4	4 <i>b</i>	0.10499(2)	0.09899(10)	0.55237(7)	0.00424(12)	1
P5	4 <i>b</i>	0.23046(2)	0.39854(10)	0.24582(8)	0.00399(12)	1
P6	4 <i>b</i>	0.23377(2)	0.09840(10)	0.55549(8)	0.00440(12)	1
N1	4 <i>b</i>	0.43890(6)	0.1729(4)	0.0873(3)	0.0055(3)	1
N2	4 <i>b</i>	0.05901(7)	0.6923(4)	0.2099(3)	0.0058(3)	1
N3	4 <i>b</i>	0.10446(7)	0.1916(4)	0.0819(3)	0.0052(3)	1
N4	4 <i>b</i>	0.10868(6)	0.3269(4)	0.7051(3)	0.0061(3)	1
N5	4 <i>b</i>	0.22376(6)	0.1797(4)	0.0910(3)	0.0055(3)	1
N6	4 <i>b</i>	0.27019(7)	0.6859(4)	0.2152(3)	0.0062(3)	1
N7	4 <i>b</i>	0.33052(6)	0.0165(4)	0.0100(3)	0.0078(3)	1
N8	2 <i>a</i>	0	0.4786(5)	0.7778(3)	0.0066(4)	1
N9	4b	0.16386(5)	0.5155(4)	0.2848(3)	0.0064(3)	1
N10	4b	0.06739(5)	0.2302(3)	0.3972(4)	0.0052(3)	1
N11	4b	0.40067(6)	0.2674(3)	0.3987(4)	0.0065(2)	1
N12	4b	0.26601(6)	0.2692(3)	0.4040(3)	0.0065(3)	1
N13	2 <i>a</i>	0	0.0092(5)	0.0020(3)	0.0077(4)	1

Table E9 Atomic coordinates, isotropic displacement parameters (/Å<sup>2</sup>), and site occupancies of disordered  $TiP_4N_8$ .

**Table E10** Anisotropic displacement parameters ( $U_{ij}$  in Å<sup>2</sup>) for disordered TiP<sub>4</sub>N<sub>8</sub> with estimated standard deviations in parentheses. The anisotropic displacement factor is expressed as exp[ $-2\pi^2$  ( $U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}l^2c^{*2} + U_{12}hka^*b^* + U_{13}hla^*c^* + U_{23}klb^*c^*$ ).

Atom	<i>U</i> <sub>11</sub>	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	<i>U</i> <sub>12</sub>
Ti1	0.00450(16)	0.00693(16)	0.0040(2)	0.0002(2)	0	0
Ti2	0.00440(13)	0.00832(15)	0.00490(19)	-0.00061(18)	-0.00001(7)	-0.00008(7)
Ti2a	0.00440(13)	0.00832(15)	0.00490(19)	-0.00061(18)	-0.00001(7)	-0.00008(7)
P1	0.00414(18)	0.00435(18)	0.0030(3)	-0.0005(2)	0.00030(17)	0.00025(13)
P2	0.00463(18)	0.00417(19)	0.0037(3)	0.0004(2)	-0.00035(18)	-0.00012(13)
Р3	0.00405(17)	0.00414(18)	0.0034(2)	-0.00040(19)	0.0000(2)	0.00003(13)
P4	0.00413(18)	0.00432(18)	0.0043(3)	0.0001(2)	0.0002(2)	-0.00013(13)
P5	0.00392(17)	0.00440(19)	0.0037(3)	-0.0004(2)	-0.00031(19)	-0.00025(13)
P6	0.00466(18)	0.00418(18)	0.0043(3)	0.0001(2)	-0.00006(18)	0.00031(13)
N1	0.0048(5)	0.0056(6)	0.0060(9)	-0.0015(6)	0.0007(5)	-0.0023(5)
N2	0.0068(5)	0.0058(6)	0.0048(8)	0.0012(6)	0.0014(6)	0.0013(5)
N3	0.0051(5)	0.0052(6)	0.0052(7)	-0.0014(6)	0.0005(6)	-0.0020(5)
N4	0.0064(6)	0.0068(6)	0.0050(9)	-0.0017(7)	-0.0018(6)	0.0010(5)
N5	0.0050(5)	0.0056(6)	0.0057(8)	-0.0020(6)	-0.0009(6)	0.0018(5)
N6	0.0065(6)	0.0068(6)	0.0053(8)	0.0018(6)	-0.0020(6)	-0.0009(5)
N7	0.0044(5)	0.0084(6)	0.0107(9)	0.0035(8)	0.0004(5)	0.0006(4)
N8	0.0040(7)	0.0067(8)	0.0090(11)	0.0028(9)	0	0
N9	0.0047(5)	0.0063(6)	0.0082(8)	-0.0025(8)	-0.0002(5)	-0.0006(4)
N10	0.0053(5)	0.0058(5)	0.0044(7)	0.0017(10)	0.0009(8)	-0.0005(4)
N11	0.0054(5)	0.0100(5)	0.0042(6)	0.0020(9)	0.0017(7)	0.0008(4)
N12	0.0051(5)	0.0104(6)	0.0041(7)	0.0014(9)	-0.0017(7)	-0.0009(4)
N13	0.0028(7)	0.0088(8)	0.0114(11)	0.0036(9)	0	0

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Ti1	N1	2x	2.098(2)	N8	P1	N1	109.45(10)
Ti1	N2	2x	2.208(2)	N8	P1	N4	110.52(10)
Ti1	N10	2x	2.2196(14)	N1	P1	N4	113.15(13)
Ti2	Ti2a	1x	1.144(3)	N8	P1	N11	115.13(12)
Ti2	N4	1x	2.085(2)	N1	P1	N11	111.81(11)
Ti2	N6	1x	2.106(2)	N4	P1	N11	96.33(9)
Ti2	N5	1x	2.209(2)	N11	P2	N2	112.83(10)
Ti2	N3	1x	2.213(2)	N11	P2	N3	105.67(9)
Ti2	N12	1x	2.2374(15)	N2	P2	N3	114.08(14)
Ti2	N11	1x	2.2470(16)	N11	P2	N13	113.49(13)
Ti2a	N5	1x	2.099(5)	N2	P2	N13	105.05(10)
Ti2a	N3	1x	2.129(4)	N3	P2	N13	105.68(10)
Ti2a	N6	1x	2.196(4)	N3	Р3	N10	112.41(11)
Ti2a	N4	1x	2.202(4)	N3	Р3	N2	114.09(13)
Ti2a	N11	1x	2.381(4)	N10	Р3	N2	106.83(9)
Ti2a	N12	1x	2.386(4)	N3	Р3	N9	105.34(9)
P1	N8	1x	1.5959(10)	N10	Р3	N9	112.65(10)
P1	N1	1x	1.619(2)	N2	Р3	N9	105.43(9)
P1	N4	1x	1.6299(18)	N7	P4	N4	109.11(9)
P1	N11	1x	1.647(2)	N7	P4	N1	110.70(9)
P2	N11	1x	1.611(3)	N4	P4	N1	113.16(14)
P2	N2	1x	1.619(2)	N7	P4	N10	116.14(11)
P2	N3	1x	1.6221(19)	N4	P4	N10	112.10(10)
P2	N13	1x	1.6601(11)	N1	P4	N10	95.15(9)
Р3	N3	1x	1.620(2)	N5	P5	N6	116.12(12)
Р3	N10	1x	1.626(2)	N5	P5	N12	115.63(10)
Р3	N2	1x	1.6273(19)	N6	P5	N12	97.76(9)
Р3	N9	1x	1.6483(14)	N5	P5	N9	105.25(9)
P4	N7	1x	1.6064(15)	N6	P5	N9	106.56(9)
P4	N4	1x	1.621(2)	N12	P5	N9	115.50(11)
P4	N1	1x	1.6260(18)	N7	P6	N5	109.37(9)
P4	N10	1x	1.638(2)	N7	P6	N6	109.24(9)
P5	N5	1x	1.613(2)	N5	P6	N6	111.56(12)
P5	N6	1x	1.6200(18)	N7	P6	N12	113.55(11)
P5	N12	1x	1.630(2)	N5	P6	N12	103.75(9)
P5	N9	1x	1.6474(14)	N6	P6	N12	109.32(11)
P6	N7	1x	1.6069(15)				
P6	N5	1x	1.6288(18)				

Table E11 Interatomic distances (/Å) and bond angles (/°) in the structure of disordered  $TiP_4N_8$ .

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P6	N6	1x	1.631(2)
P6	N12	1x	1.632(2)

Table E12 Atomic coordinates, isotropic displacement parameters (/Å<sup>2</sup>), and site occupancies of  $\beta$ -

TiP4N8 at HT.

Atom	Wyckoff	X	У	Ζ	$U_{eq}$	осс.
	position					
Ti1	2 <i>a</i>	0	0.88420(17)	0.39418(13)	0.0161(2)	1
Ti2	4b	0.33332(3)	0.62120(15)	0.40469(12)	0.0206(2)	1
P1	4b	0.43551(4)	0.39932(19)	0.23748(12)	0.0110(2)	1
P2	4 <i>b</i>	0.43389(4)	0.09922(19)	0.54809(13)	0.0116(2)	1
Р3	4 <i>b</i>	0.09807(4)	0.40398(18)	0.24156(12)	0.0111(2)	1
P4	4 <i>b</i>	0.10473(4)	0.10293(19)	0.55210(12)	0.0119(2)	1
P5	4 <i>b</i>	0.22996(4)	0.40090(19)	0.24674(13)	0.0124(2)	1
P6	4 <i>b</i>	0.23394(4)	0.09939(19)	0.55727(13)	0.0126(2)	1
N1	4 <i>b</i>	0.43858(14)	0.1694(8)	0.0878(6)	0.0172(6)	1
N2	4 <i>b</i>	0.05888(14)	0.6935(7)	0.2108(5)	0.0148(6)	1
N3	4 <i>b</i>	0.10349(14)	0.1949(7)	0.0813(5)	0.0152(6)	1
N4	4 <i>b</i>	0.10919(14)	0.3320(7)	0.7053(5)	0.0152(6)	1
N5	4 <i>b</i>	0.22510(14)	0.1803(7)	0.0925(6)	0.0167(6)	1
N6	4 <i>b</i>	0.27091(14)	0.6830(7)	0.2172(5)	0.0153(6)	1
N7	4 <i>b</i>	0.33060(13)	0.0094(9)	0.0123(5)	0.0232(8)	1
N8	2 <i>a</i>	0	0.4769(9)	0.7734(6)	0.0164(8)	1
N9	4 <i>b</i>	0.16395(11)	0.5218(8)	0.2821(6)	0.0216(7)	1
N10	4 <i>b</i>	0.06771(12)	0.2294(7)	0.3937(4)	0.0144(6)	1
N11	4 <i>b</i>	0.40100(14)	0.2628(7)	0.3960(4)	0.0172(6)	1
N12	4 <i>b</i>	0.26554(13)	0.2686(8)	0.4033(4)	0.0179(6)	1
N13	2a	0	0.0044(11)	0.0017(7)	0.0218(10)	1

**Table E13** Anisotropic displacement parameters ( $U_{ij}$  in Å<sup>2</sup>) for  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> at HT with estimated standard deviations in parentheses. The anisotropic displacement factor is expressed as exp[ $-2\pi^2$  ( $U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}l^2c^{*2} + U_{12}hka^*b^* + U_{13}hla^*c^* + U_{23}klb^*c^*$ ).

Atom	$U_{11}$	<i>U</i> <sub>22</sub>	<i>U</i> <sub>33</sub>	<i>U</i> <sub>23</sub>	<i>U</i> <sub>13</sub>	$U_{12}$
Ti1	0.0150(4)	0.0217(4)	0.0116(4)	-0.0008(4)	0	0
Ti2	0.0136(3)	0.0264(4)	0.0217(4)	-0.0029(3)	-0.00007(16)	-0.00033(17)
P1	0.0107(4)	0.0105(4)	0.0119(4)	-0.0018(3)	0.0010(3)	0.0008(3)
P2	0.0132(4)	0.0106(4)	0.0111(5)	0.0013(3)	-0.0014(3)	0.0000(3)
Р3	0.0095(3)	0.0099(4)	0.0138(4)	-0.0012(3)	-0.0004(4)	0.0008(3)
P4	0.0114(4)	0.0115(4)	0.0130(5)	0.0009(3)	0.0006(4)	-0.0002(3)
P5	0.0110(4)	0.0108(4)	0.0153(5)	-0.0013(3)	-0.0012(4)	-0.0006(3)
P6	0.0126(4)	0.0107(4)	0.0144(5)	0.0012(4)	0.0005(4)	0.0003(3)
N1	0.0127(11)	0.0178(13)	0.0210(16)	-0.0054(12)	0.0001(12)	-0.0060(10)
N2	0.0176(12)	0.0123(12)	0.0147(13)	0.0039(11)	0.0054(12)	0.0041(10)
N3	0.0136(10)	0.0143(12)	0.0177(13)	-0.0021(11)	-0.0012(12)	-0.0036(11)
N4	0.0167(13)	0.0143(11)	0.0147(13)	-0.0044(10)	-0.0056(12)	0.0025(11)
N5	0.0152(11)	0.0156(12)	0.0194(16)	-0.0044(11)	0.0005(12)	0.0053(11)
N6	0.0168(13)	0.0137(12)	0.0154(14)	0.0053(11)	-0.0056(12)	-0.0039(11)
N7	0.0115(12)	0.0233(17)	0.035(2)	0.0106(16)	-0.0005(13)	-0.0002(10)
N8	0.0096(15)	0.0162(17)	0.023(2)	0.0074(16)	0	0
N9	0.0103(12)	0.0164(14)	0.038(2)	-0.0095(15)	-0.0010(13)	0.0001(9)
N10	0.0142(13)	0.0149(11)	0.0141(12)	0.0047(16)	0.0035(10)	-0.0014(9)
N11	0.0146(12)	0.0207(13)	0.0163(13)	0.0012(17)	0.0058(10)	0.0018(9)
N12	0.0149(13)	0.0230(14)	0.0159(15)	0.0046(17)	-0.0060(10)	-0.0028(10)
N13	0.0111(17)	0.022(2)	0.032(3)	0.0103(19)	0	0

Ti1         N1         2x           Ti1         N2         2x           Ti1         N10         2x	2.123(4) 2.192(4) 2.224(3)	N1 N1	P1	N8	109.1(2)
			P1	N4	114.0(3)
111 N10 2X		N8	P1	N4 N4	109.8(2)
Ti2 N4 1x	2.097(4)	NO N1	P1	N4 N11	111.1(2)
Ti2 N4 1x		N1 N8	P1	N11 N11	
	2.107(4)				116.1(2)
	2.206(4)	N4	P1	N11	96.46(18)
Ti2 N5 1x	2.225(4)	N3	P2	N11	105.39(18)
Ti2 N12 1x	2.249(3)	N3	P2	N13	106.6(2)
Ti2 N11 1x	2.268(3)	N11	P2	N13	113.0(2)
P1 N1 1x	1.611(4)	N3	P2	N2	114.0(2)
P1 N8 1x	1.613(2)	N11	P2	N2	112.74(18)
P1 N4 1x	1.628(3)	N13	P2	N2	105.1(2)
P1 N11 1x	1.634(4)	N3	Р3	N2	114.0(2)
P2 N3 1x	1.626(3)	N3	Р3	N10	110.2(2)
P2 N11 1x	1.629(4)	N2	Р3	N10	106.47(17)
P2 N13 1x	1.635(2)	N3	Р3	N9	106.63(19)
P2 N2 1x	1.635(4)	N2	Р3	N9	105.65(19)
P3 N3 1x	1.621(4)	N10	Р3	N9	114.0(2)
P3 N2 1x	1.628(3)	N7	P4	N1	110.7(2)
P3 N10 1x	1.628(3)	N7	P4	N4	107.65(18)
P3 N9 1x	1.640(3)	N1	P4	N4	113.7(3)
P4 N7 1x	1.605(3)	N7	P4	N10	115.8(2)
P4 N1 1x	1.626(3)	N1	P4	N10	95.49(18)
P4 N4 1x	1.632(4)	N4	P4	N10	113.34(18)
P4 N10 1x	1.645(3)	N5	P5	N6	115.5(2)
P5 N5 1x	1.614(4)	N5	P5	N12	113.7(2)
P5 N6 1x	1.621(3)	N6	Р5	N12	97.11(18)
P5 N12 1x	1.627(3)	N5	Р5	N9	106.6(2)
P5 N9 1x	1.640(3)	N6	Р5	N9	106.85(19)
P6 N7 1x	1.606(3)	N12	Р5	N9	117.1(2)
P6 N5 1x	1.620(3)	N7	P6	N5	109.1(2)
P6 N12 1x	1.639(4)	N7	P6	N12	112.6(2)
P6 N6 1x	1.641(4)	N5	P6	N12	104.76(18)
		N7	P6	N6	107.88(18)
		N5	P6	N6	112.7(2)
		N12	P6	N6	109.87(19)

Table E14 Interatomic distances (/Å) and bond angles (/°) in the structure of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> at HT.

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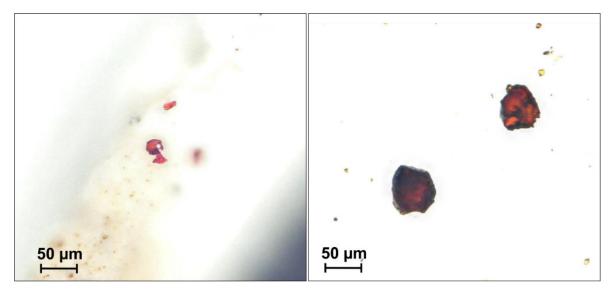
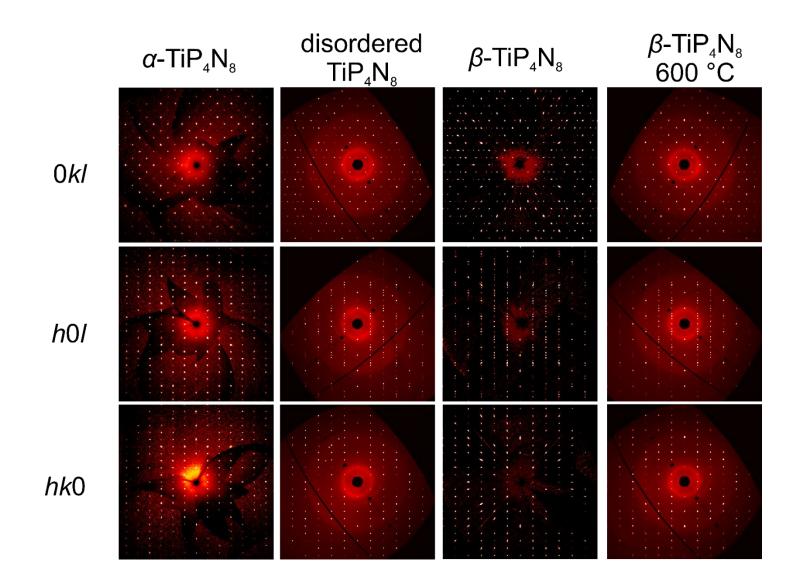


Figure E3. Photographs of representative crystals of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (left) and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (right).



**Figure E4**. Reciprocal lattice sections of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>, disordered  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>,  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> and HT  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>.

#### **EDX measurements**

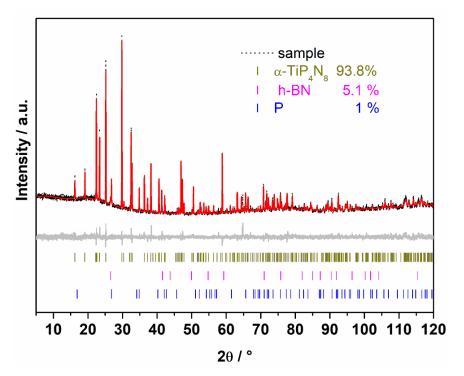
EDX point / atom %	Ti	Р	Ν
1	6.2	28.8	63.9
2	6.4	29.3	63.2
3	6.8	30.5	61.7
4	6.6	33.2	60.2
5	7.5	33.8	58.8
6	7.9	35.8	56.3
average	6.9(6)	32(2)	60(2)
calculated	7.7	30.8	61.5

 Table E15. EDX measurements for TiP<sub>4</sub>N<sub>8</sub> compared to values derived from sum formula.

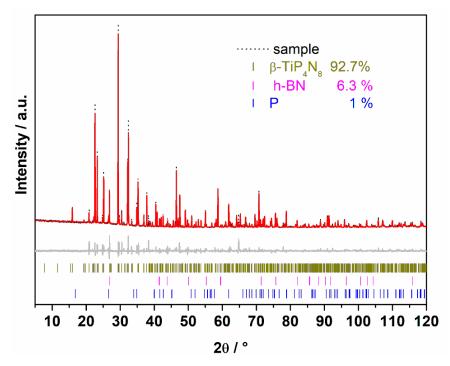
#### **Rietveld refinements**

**Table E16.** Crystallographic data from the Rietveld refinements for  $\alpha$ - and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> standard deviations are given in parentheses.

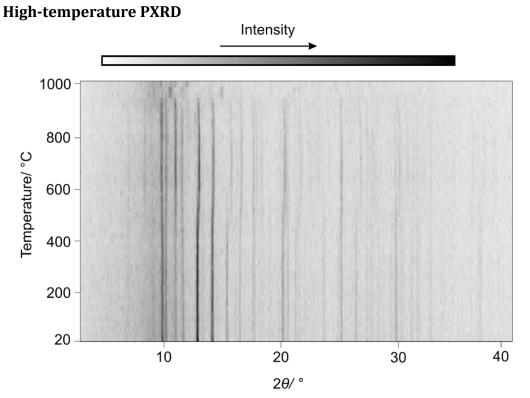
	α-TiP <sub>4</sub> N <sub>8</sub>	β-TiP <sub>4</sub> N <sub>8</sub>
crystal system	orthorhombic	
space group	<i>Pmn</i> 2 <sub>1</sub> (no. 31)	
lattice parameters / Å	a = 7.6063(1)	a = 22.9196(5)
	b = 4.6354(1)	b = 4.58800(10)
	c = 7.8418(2)	c = 8.0970(2)
radiation $\lambda$ / Å	1.54056	6 (Cu-K <sub>α1</sub> )
cell volume / ų	276.49(1)	852.27(1)
$2\theta$ – range/°	5 <2	<i>θ</i> < 120
refined parameters (incl. side phases)	53	119
thereof	18	12
background		
<i>R</i> -values	$R_{\rm p} = 0.0625$	$R_{\rm p} = 0.0880$
	$wR_{\rm p} = 0.0821$	$wR_{\rm p} = 0.1383$
	$R_{\rm Bragg} = 0.0334$	$R_{\mathrm{Bragg}} = 0.0400$



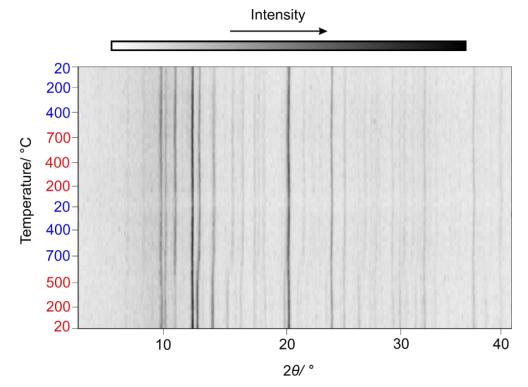
**Figure E5.** Rietveld refinement for  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (vertical yellow bars), h-BN (vertical magenta bars), black P (vertical blue bars) and difference profile (gray line).



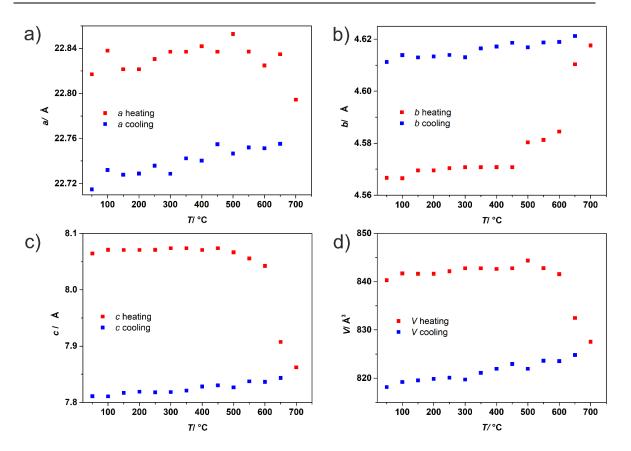
**Figure E6.** Rietveld refinement for  $\beta$ -TiP<sub>4</sub>N<sub>8</sub>; observed (black line) and calculated (red line) X-ray powder diffraction patterns, positions of Bragg reflections of  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (vertical yellow bars), h-BN (vertical magenta bars), black P (vertical blue bars) and difference profile (gray line).



**Figure E7.** Temperature-dependent powder X-ray diffraction patterns (Mo-K<sub> $\alpha$ 1</sub> radiation,  $\lambda$  = 0.71073 Å) of TiP<sub>4</sub>N<sub>8</sub> measured under Ar atmosphere showing thermal stability of TiP<sub>4</sub>N<sub>8</sub> until 950 °C.



**Figure E8.** Temperature-dependent powder X-ray diffraction patterns (Mo-K<sub> $\alpha$ 1</sub> radiation,  $\lambda$  = 0.71073 Å) of TiP<sub>4</sub>N<sub>8</sub> measured under Ar atmosphere showing a change of lattice parameters at 700°C of the first heating ramp. Heating in red and cooling in blue.



**Figure E9.** Change of lattice parameters and unit cell volume of  $TiP_4N_8$  from the first HTPXRD cycle. Heating curves in red and cooling curves in blue.

#### Bond-valence sum (BVS) Calculations

Table E17 Bond valence sums and expected charges for the atom sites in the crystal structure of  $\alpha$ -

site	bond-valence sum	expected charge
Ti1	3.69	+4
P1	5.06	+5
P2	5.06	+5
N1	3.07	-3
N2	3.07	-3
N3	2.98	-3
N4	2.89	-3
N5	2.82	-3

Table E18 Bond valence sums and expected charges for the atom sites in the crystal structure of  $\beta$ -

site	bond-valence sum	expected charge
Ti1	3.18	+4
Ti2	3.18	+4
P1	5.05	+5
P2	4.98	+5
Р3	4.95	+5
P4	5.03	+5
P5	5.01	+5
P6	5.03	+5
N1	3.14	-3
N2	2.98	-3
N3	3.00	-3
N4	3.12	-3
N5	3.01	-3
N6	3.10	-3
N7	2.67	-3
N8	2.64	-3
N9	2.47	-3
N10	2.89	-3
N11	2.89	-3
N12	2.90	-3
N13	2.51	-3

TiP4N8.

TiP4N8.

#### IR spectroscopy

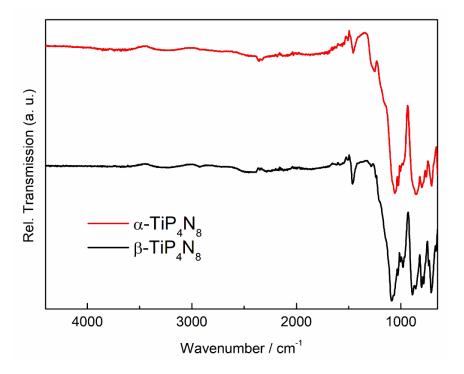
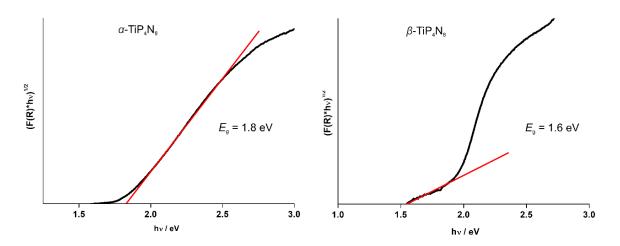
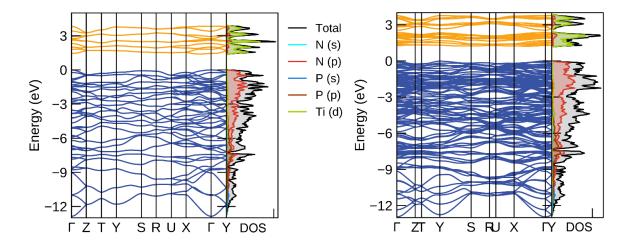


Figure E10. IR spectra of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (red) and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (black).

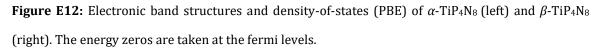
#### **UV-Vis spectra**

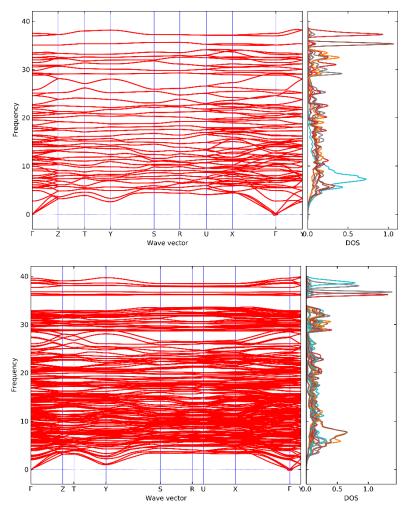


**Figure E11:** Tauc plots for UV-Vis spectra of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (left) and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (right). Red lines represent linear regressions from the steep regions to estimate the direct optical band gaps from the inflection points with the abscissa. Data between 1.7 and 2.4 eV ( $\alpha$ -TiP<sub>4</sub>N<sub>8</sub>) and 1.5 to 1.9 eV ( $\beta$ -TiP<sub>4</sub>N<sub>8</sub>) were used for linear regressions.

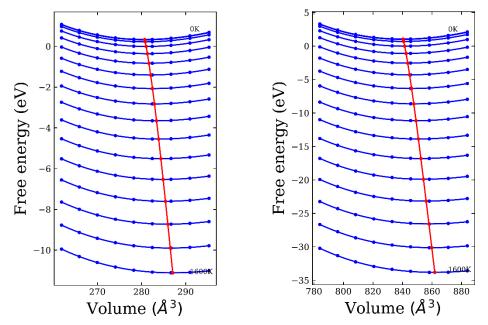


### **DFT-calculations**





**Figure E13:** Phonon band structures and phonon density-of-states of  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (top) and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (bottom).



**Figure E14:** Volume dependent Helmoltz free energies at different temperatures for  $\alpha$ -TiP<sub>4</sub>N<sub>8</sub> (left) and  $\beta$ -TiP<sub>4</sub>N<sub>8</sub> (right). The red lines mark the minima where the Gibbs free energy is obtained, respectively.

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### F Miscellaneous

### F.1 List of Publications within this Thesis

### 1. A Layered Tin Bismuth Selenide with Three Different Building Blocks that Account for an Extremely Large Lattice Parameter of 283 Å

M. Nentwig, L. Eisenburger, F. Heinke, D. Souchay, O. Oeckler

published in: Chem. Eur. J. 2020, 26, 10676–10681.

XRD analyses were performed by Markus Nentwig. Syntheses were performed by Frank Heinke. STEM analyses and cross section sample preparation were performed by Lucien Eisenburger. TEM characterization was performed by Daniel Souchay. Oliver Oeckler supervised the project. The manuscript was written by Markus Nentwig in a leading role with support of all co-authors. All authors revised the manuscript.

# 2.High-PressureHigh-TemperatureSynthesisofMixedNitridosilicatephosphates and Luminescence of $AESiP_3N_7$ :Eu<sup>2+</sup> (AE = Sr, Ba)

L. Eisenburger, O. Oeckler, W. Schnick

published in: *Chem. Eur. J.* **2021**, *27*, 4461–4465.

Syntheses and analyses were performed by Lucien Eisenburger. Oliver Oeckler and Wolfgang Schnick supervised the project. The manuscript was written by Lucien Eisenburger in a leading role with support of all co-authors. All authors revised the manuscript.

Nitridic Analogs of Micas AESi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (AE = Mg, Mg<sub>0.94</sub>Ca<sub>0.06</sub>, Ca, Sr)
 L. Eisenburger, P. Strobel, P. J. Schmidt, T. Bräuniger, J. Wright, E. L. Bright, C. Giacobbe, O. Oeckler, W. Schnick

published in: *Angew. Chem. Int. Ed.* **2022**, 61, e202114902. *Angew. Chemie* **2022**, *134*, e202114902.

Syntheses, XRD and STEM analyses were performed by Lucien Eisenburger. Philipp Strobel and Peter. J. Schmidt performed luminescence measurements and analyses thereof. Thomas Bräuniger performed NMR measurements and aided in interpreteation of data. Jonathan Wright, Eleanonar Lawrence Bright and Carlotta Giacobbe measured X-ray diffraction data at ID11, ESRF, Grenoble. Oliver Oeckler and Wolfgang Schnick supervised the project. The manuscript was written by Lucien Eisenburger in a leading role with support of all co-authors. All authors revised the manuscript.

# 4. High-Pressure Synthesis of $Sc_5P_{12}N_{23}O_3$ and $Ti_5P_{12}N_{24}O_2$ by Activation of the Binary Nitrides ScN and TiN with NH<sub>4</sub>F

L. Eisenburger, V. Weippert, O. Oeckler, W. Schnick

published in: *Chem. Eur. J.* **2021**, *27*, 14184–14188.

Syntheses, XRD and STEM analyses were performed by Lucien Eisenburger. Conducting and analyzing magnetic measurements was done by Valentin Weippert. Oliver Oeckler and Wolfgang Schnick supervised the project. The manuscript was written by Lucien Eisenburger in a leading role with support of all co-authors. All authors revised the manuscript.

#### 5. Discovery of Two Polymorphs of TiP<sub>4</sub>N<sub>8</sub> Synthesized from Binary Nitrides

L. Eisenburger, V. Weippert, C. Paulmann, D. Johrendt, O. Oeckler, W. Schnick published in: *Angew. Chem. Int. Ed.* **2022**, DOI 10.1002/anie.202202014. *Angew. Chemie* **2022**, DOI 10.1002/ange.202202014.

Syntheses, XRD and STEM analyses were performed by Lucien Eisenburger. Conducting and analyzing DFT calculations was done by Valentin Weippert and Dirk Johrendt. Carsten Paulmann aided in SCXRD data collection at P24, DESY, Hamburg. Oliver Oeckler and Wolfgang Schnick supervised the project. The manuscript was written by Lucien Eisenburger in a leading role with support of all co-authors. All authors revised the manuscript.

### F.2 Full List of Publications

# 1. Nucleophilicity of Alkyl Zirconocene and Titanocene Precatalysts, and Kinetics of Activation by Carbenium Ions and by B(C<sub>6</sub>F<sub>5</sub>)<sub>3</sub>

G. Berionni, H. Kurouchi, <u>L. Eisenburger</u>, H. Mayr, *Chem. Eur. J.* **2016**, *22*, 11196–11200.

# 2. The Influence of Nanoscale Heterostructures on the Thermoelectric Properties of Bi-substituted Tl<sub>5</sub>Te<sub>3</sub>

F. Heinke, <u>L. Eisenburger</u>, R. Schlegel, S. Schwarzmüller, O. Oeckler, *Z. Anorg. Allg. Chem.* **2017**, *643*, 447–454.

### 3. The Crystal Structures of Pb<sub>5</sub>Sb<sub>4</sub>S<sub>11</sub> (Boulangerite) – A Phase Transition Explains Seemingly Contradictory Structure Models

P. Schultz, F. Nietschke, G. Wagner, C. Eikemeier, <u>L. Eisenburger</u>, O. Oeckler, *Z. Anorg. Allg. Chem.* **2017**, *643*, 1531–1542.

### 4. $RE_4Ba_2[Si_{12}O_2N_{16}C_3]:Eu^{2+}$ (RE = Lu, Y): Green-Yellow Emitting Oxonitridocarbidosilicates with a Highly Condensed Network Structure

C. Maak, <u>L. Eisenburger</u>, J. P. Wright, M. Nentwig, P. J. Schmidt, O. Oeckler, W. Schnick, *Inorg. Chem.* **2018**, *57*, 13840–13846.

# 5. Ammonothermal Synthesis and Crystal Structure of the Nitridoalumogermanate $Ca_{1-x}Li_xAl_{1-x}Ge_{1+x}N_3$ ( $x \approx 0.2$ )

J. Häusler, <u>L. Eisenburger</u>, O. Oeckler, W. Schnick, *Eur. J. Inorg. Chem.* **2018**, *2018*, 759–764.

# 6. Syntheses and physical properties of the MAX phase boride Nb<sub>2</sub>SB and the solid solutions Nb<sub>2</sub>SB<sub>x</sub>C<sub>1-x</sub> (x = 0-1)

T. Rackl, <u>L. Eisenburger</u>, R. Niklaus, D. Johrendt *Phys. Rev. Mater.* **2019**, *3*, 054001. 7. Ammonothermal Synthesis of the Mixed-Valence Nitrogen-Rich Europium Tantalum Ruddlesden-Popper Phase  $Eu^{III}Eu^{III}_2Ta_2N_4O_3$ 

N. Cordes, M. Nentwig, <u>L. Eisenburger</u>, O. Oeckler, W. Schnick *Eur. J. Inorg. Chem.* **2019**, *2019*, 2304–2311.

8. The Long-Periodic Loop-Branched Chain Structure of the Oxonitridophosphate  $La_{21}P_{40}O_{46}N_{57}$ , Elucidated by a Combination of TEM and Microfocused Synchrotron Radiation

M. Nentwig, S. D. Kloß, L. Neudert, <u>L. Eisenburger</u>, W. Schnick, O. Oeckler *Chem. Eur. J.* **2019**, *25*, 14382–14387.

### 9. Cationic $Pb_2$ Dumbbells Stabilized in the Highly Covalent Lead Nitridosilicate $Pb_2Si_5N_8$

P. Bielec, R. Nelson, R. P. Stoffel, <u>L. Eisenburger</u>, D. Günther, A. K. Henß, J. P. Wright, O. Oeckler, R. Dronskowski, W. Schnick *Angew. Chem. Int. Ed.* **2019**, *58*, 1432–1436. *Angew. Chem.* **2019**, *131*, 1446–1450.

### 10. Targeting Vacancies in Nitridosilicates: Aliovalent Substitution of $M^{2+}(M = \text{Ca, Sr})$ by Sc<sup>3+</sup> and U<sup>3+</sup>

P. Bielec, <u>L. Eisenburger</u>, H. L. Deubner, D. Günther, F. Kraus, O. Oeckler, W. Schnick *Angew. Chem. Int. Ed.* **2019**, *58*, 840–843. *Angew. Chem.* **2019**, *131*, 850-853.

### 11. BaP<sub>6</sub>N<sub>10</sub>NH:Eu<sup>2+</sup> as a Case Study-An Imidonitridophosphate Showing Luminescence

S. Wendl, <u>L. Eisenburger</u>, M. Zipkat, D. Günther, J. P. Wright, P. J. Schmidt, O. Oeckler, W. Schnick

*Chem. Eur. J.* **2020**, *26*, 5010–5016.

### 12. Nitridophosphate-Based Ultra-Narrow-Band Blue-Emitters: Luminescence Properties of *AE*P<sub>8</sub>N<sub>14</sub>:Eu<sup>2+</sup> (*AE* = Ca, Sr, Ba)

S. Wendl, L. Eisenburger, P. Strobel, D. Günther, J. P. Wright, P. J. Schmidt, O. Oeckler,

W. Schnick

Chem. Eur. J. 2020, 26, 7292–7298.

# 13. A Layered Tin Bismuth Selenide with Three Different Building Blocks that Account for an Extremely Large Lattice Parameter of 283 Å

M. Nentwig, <u>L. Eisenburger</u>, F. Heinke, D. Souchay, O. Oeckler *Chem. Eur. J.* **2020**, *26*, 10676–10681.

#### 14. BaGe<sub>8</sub>As<sub>14</sub>: a semiconducting sodalite-type compound

V. Weippert, T. Chau, K. Witthaut, <u>L. Eisenburger</u>, D. Johrendt *Chem. Commun.* **2021**, *57*, 1332–1335.

### 15. High Thermoelectric Properties in the Sodalite Compounds $BaGe_8As_{14}$ and $AGe_7As_{15}(A = Rb, Cs)$

V. Weippert, K. Witthaut, M. Pointner, M. Sachs, <u>L. Eisenburger</u>, F. Kraus, D. Johrendt *Chem. Mater.* **2021**, *33*, 8248–8258.

### 16. High-Pressure High-Temperature Synthesis of Mixed Nitridosilicatephosphates and Luminescence of AESiP<sub>3</sub>N<sub>7</sub>:Eu<sup>2+</sup> (AE = Sr, Ba)

L. Eisenburger, O. Oeckler, W. Schnick *Chem. Eur. J.* **2021**, *27*, 4461–4465.

# 17. High-Pressure Synthesis of $Sc_5P_{12}N_{23}O_3$ and $Ti_5P_{12}N_{24}O_2$ by Activation of the Binary Nitrides ScN and TiN with NH<sub>4</sub>F

<u>L. Eisenburger</u>, V. Weippert, O. Oeckler, W. Schnick *Chem. Eur. J.* **2021**, *27*, 14184–14188.

18. Nitridic Analogs of Micas AESi<sub>3</sub>P<sub>4</sub>N<sub>10</sub>(NH)<sub>2</sub> (AE = Mg, Mg<sub>0.94</sub>Ca<sub>0.06</sub>, Ca, Sr)
L. Eisenburger, P. Strobel, P. J. Schmidt, T. Bräuniger, J. Wright, E. L. Bright, C. Giacobbe, O. Oeckler, W. Schnick
Angew. Chem. Int. Ed. 2022, 61, e202114902.
Angew. Chem. 2022, 134, e202114902.

# 19. Discovery of Two Polymorphs of TiP<sub>4</sub>N<sub>8</sub> Synthesized from Binary Nitrides L. Eisenburger, V. Weippert, C. Paulmann, D. Johrendt, O. Oeckler, W. Schnick Angew. Chem. Int. Ed. 2022, DOI 10.1002/anie.202202014. Angew. Chem. 2022, DOI 10.1002/ange.202202014.

### F.3 Conference Contributions

#### Heterostructured multinary thallium tellurides

Frank Heinke, <u>Lucien Eisenburger</u>, Stefan Schwarzmüller, Frederik Nietschke and Oliver Oeckler

Talk, Hirschegg-Seminar für Festkörperchemie, 2016, Hirschegg, Austria

#### Cation disorder in $Ca_{1-x}Li_xAl_{1-x}Ge_{1+x}N_3$ ( $x \approx 0.2$ ) by STEM-HAADF

Lucien Eisenburger, Jonas Häusler, Wolfgang Schnick, Oliver Oeckler

*Talk*, 26<sup>th</sup> Annual Meeting of the German Crystallographic Society (DGK), **2018**, Essen, Germany

### $\label{eq:kationen-und} Kationen- und Anionen-Ausordnung in der Ruddlesden-Popper-Phase Eu^2+Eu^{3+}_2Ta_2N_4O_3$

Lucien Eisenburger, Niklas Cordes, Markus Nentwig, Juliane Stahl, Theresa Block, Rainer Pöttgen, Dirk Johrendt, Wolfgang Schnick, Oliver Oeckler

Talk, Mitteldeutsches Anorganisches Nachwuchs Symposium, 2018, Jena, Germany

### Influence of dopant concentration on dopant distribution in $Ca_{2-x}Eu_xSi_5N_8$ (x = 0.02, 0.1, 0.15) by STEM-HAADF

Lucien Eisenburger, Monika Pointner, Oliver Oeckler and Wolfgang Schnick

Poster, The 17<sup>th</sup> European Conference on Solid State Chemistry, **2019**, Lille, France

# Crystal structure determination of microcrystalline compounds by the combination of TEM and microfocused synchrotron radiation

Daniel Günther, Johnathan P. Wright, Gavin B. M. Vaughan, Sebastian Wendl, <u>Lucien</u> <u>Eisenburger</u>, Wolfgang Schnick and Oliver Oeckler,

*Talk*, Microscopy Conference, **2019**, Berlin, Germany

#### Structural characterization of barium (imido-)nitridophosphate networks

Daniel Günther, Sebastian Wendl, <u>Lucien Eisenburger</u>, Wolfgang Schnick and Oliver Oeckler *Poster*, 27<sup>th</sup> Annual Meeting of the German Crystallographic Society (DGK), **2019**, Leipzig, Germany

#### Nitridoberyllosilicates – Promising Host Materials for Solid-State Lighting

Tobias Giftthaler, Philipp Strobel, <u>Lucien Eisenburger</u>, Oliver Oeckler and Wolfgang Schnick *Poster*, 10th International Symposium on Nitrides, **2019**, Barcelona, Spain

#### Nitridosilicatephosphates by high-pressure / high-temperature synthesis

<u>Lucien Eisenburger</u>, Oliver Oeckler and Wolfgang Schnick *Talk*, Joint Polish-German Crystallographic Meeting, **2020**, Wrocław, Poland

#### NH<sub>4</sub>F-Assisted Multi-Anvil Syntheses of Novel Nitrides

<u>Lucien Eisenburger</u>, Oliver Oeckler and Wolfgang Schnick *Talk*, International Workshop on "Development of multianvil technology and its applications to Earth and material sciences", **2021**, virtual venue

#### Disorder in layered tellurides $Tt_1Pn_2Te_4$ (Tt = Ge, Sn, Pb; Pn = As, Sb, Bi)

Lennart Staab, <u>Lucien Eisenburger</u>, Sebastian. Geisler and Oliver Oeckler *Talk*, 29<sup>th</sup> Annual Meeting of the German Crystallographic Society (DGK), **2021**, Hamburg, Germany/online

### F.4 Deposited Crystal Structures

Compound	CSD
SrSiP <sub>3</sub> N <sub>7</sub>	2050660
BaSiP <sub>3</sub> N <sub>7</sub>	2050661
$Sc_5P_{12}N_{23}O_3$	2084626
$Ti_5P_{12}N_{24}O_2$	2084627
$CaSi_3P_4N_{10}(NH)_2$	2096670
α-TiP4N8	2121162
β-TiP4N8	2121163
disordered β-TiP <sub>4</sub> N <sub>8</sub>	2121164
$HT \beta - TiP_4N_8$	2121161

**Table F.1:** Compounds discussed within this thesis and their respective depository numbers.

### G Curriculum Vitae

#### **Personal Details**

Name	Lucien André Eisenburger
Date of Birth	11.06.1991
Place of Birth	Düsseldorf, Germany
Nationality	German

01/2018 - 09/2022Research assistant in the groups of Prof. Dr. Wolfgang Schnick<br/>(Department of Chemistry, Ludwig-Maximilians-Universität<br/>München) and of Prof. Dr. Oliver Oeckler (Institute for Material<br/>Science and Crystallography, Leipzig University, Leipzig)<br/>Doctoral thesis: "NH4F-Assisted High-Pressure Synthesis of Novel<br/>Nitrides and Atomic-Resolution Chemical Mapping by STEM"

09/2017-02/2018 **Research assistant** in the group of Prof. Dr. Oliver Oeckler (Institute for Material Science and Crystallography, Leipzig University, Leipzig)

02/2017-08/2017 Master thesis (Inorganic Chemistry) Title: "Transmissionselektronenmikroskopie an nitridischen Festkörpern"

07/2016-02/2017 **Managing Director** Fußorthopädie Eisenburger GmbH, 56457 Westerburg

02/2016-05/2016 Research internship (Inorganic Chemistry) Universität Leipzig Prof. Dr Oliver Oeckler

04/2015-02/2017	Master of Science (Chemistry)
	Ludwig-Maximilians-Universität München
09/2014-01/2015	Bachelor Thesis (Inorganic Chemistry)
09/2014-01/2013	
	Ludwig-Maximilians-Universität München, Prof. Dr. Wolfgang
	Schnick
	Title: "Beiträge zur Strukturaufklärung nitridischer
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09/2011-09/2014	Bachelor of Science (Chemistry and Biochemistry)
	Ludwig-Maximilians-Universität München, Germany
04/2010-06/2011	Military Service
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02/2010	Allgemeine Hochschulreife
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