# Force and Affinity in Cellulosomal Complexes

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# Force and Affinity in Cellulosomal Complexes

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### ZUSAMMENFASSUNG

In dieser Arbeit werden die molekularen Mechanismen der Organisation des Cellulosoms - ein komplexes extrazelluläres Proteinnetzwerk - als Modellsystem für Protein-Protein Interaktionen mittels biophysikalischer Methoden untersucht. Dieses extrazelluläre Organell ermöglicht bestimmten Bakterien die Zersetzung von Cellulose, indem es Enzyme und Cellulose-Bindedomänen auf gerüstartigen Proteinstrukturen in synergistischer Weise kombiniert. Die einzelnen Komponenten werden hierbei von einer Klasse von Rezeptor-Liganden-Paaren namens Cohesin- Dockerin in ihrer Stöchiometrie und Anordnung funktionell kombiniert.

Ein Teil dieser Arbeit besteht in der Entschlüsselung der molekularen Bindemechanismen des Cohesins CohE, welches das Bakterium *Ruminococcus flavefaciens* mit seinem Cellusom verbindet. Durch die Kombination von Einzelmolekül-Kraftspektroskopie mit Molekulardynamik-Simulationen konnte die aussergewöhnliche Belastbarkeit der Interaktionen von CohE mit zwei homologen Dockerinen entschlüsselt werden. Hierbei wurde insbesondere der Einfluss der Kraftpropagation innerhalb eines Proteinkomplexes auf dessen mechanische Widerstandsfähigkeit untersucht. Die physiologische Verankerung über den carboxyl-Terminus von CohE erwies sich als deutlich robuster im Vergleich zu einer nicht nativen N-terminalen Verankerung.

Um den Kontrast zwischen hoher mechanischer Belastbarkeit bei moderaten Affinitäten im nano- bis mikromolaren Bereich besser verstehen zu können, wandte ich mich der Bestimmung der kinetischen Ratenkonstanten  $k_{off}$  und  $k_{on}$  zu, deren Quotient die Gleichgewichtskonstante bildet. Während es eine kleine Dissoziationskonstante dem Bakterium ermöglichen würde die von ihm exprimierte Nanomaschinerie fest an sich zu binden, könnte ein höheres  $k_{off}$  und  $k_{on}$  einen dynamischeren Austausch von Cellulosomen innerhalb des Mikrobioms ermöglichen. Zusätzlich stellte sich die Frage, ob die Verankerungsgeometrie auch in Abwesenheit von Kraft Einfluss auf das Bindeverhalten nehmen würde. Nachdem initiale Messungen mittels Oberflächenplasmonenresonanzspektroskopie inkonsistent waren, wurde eine neuartige, enzymbasierte Kopplungsstrategie für oberflächengebundene Affinitätsbestimmungen entwickelt. Hiermit konnte CohE funktional und spezifisch auf Sensoroberflächen immobilisiert werden. Es zeigte sich, dass in Abwesenheit von externer Kraft die Verankerungsgeometrie von CohE keinen Einfluss auf das Bindeverhalten hat. Dies bestärkt im Umkehrschluss die Hypothese, dass mechanische Stabilitäten stets geometrieabhängig zu untersuchen sind.

Im Rahmen dieser Arbeit wurden auch methodische Verbesserungen in der Einzelmolekülkraftspektroskopie erzielt. Zum einen wurde eine Strategie entwickelt, um Proteindomänen zeitsparend *in vitro* zu exprimieren und ohne weitere Aufreinigung spezifisch auf Objektträgern zu verankern. Die darauffolgende enzymatische Peptidligation eines Dockerins via Sortase A erlaubt es nun, mit hohem Durchsatz Entfaltungsstudien an Proteinen mithilfe der Cohesin-Dockerin Interaktion durchzuführen. Weiterhin ermöglichte es dieselbe Sortase-vermittelte Peptidligation, die gängigen Polyethylenlinker durch Elastin-ähnliche Peptide zu ersetzen. Dies verhindert Artefakte, die sonst durch Polyethylenlinker bei Protein-Kraftspektroskopie über 100 pN entstünden.

Zuletzt wurde der Entfaltungsprozess einer Cohesin-Domäne aus *Acetivibrio cellulolyticus* untersucht, deren Familie in vorangegangenen Studien teils bimodale Entfaltungskraftverteilungen zeigte. Durch die Kombination zweier Messmodi konnte die Kraft-Ladungsrate über fünf Größenordnungen variiert werden. Es konnte gezeigt werden, dass das dabei beobachtete Verhalten mit einer Konformationsänderung während der experimentellen Zeitskala zwischen verschiedenen, gefalteten Konformationen konsistent ist.

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Part I Context

# ⊥ Introduction

Understanding molecular interactions is a key challenge in the ongoing effort to unravel the self-organizing nature of life. Proteins navigate a vast conformational space until they adopt a functional, folded state - a topic that may at first glance not seem best explored by the reductionist approach of a physics. In recent decades however, techniques have emerged that allow us to manipulate and probe single biomolecules, to investigate how they interact with each other, or how they respond to external stimuli. The invention and development of techniques such as singlemolecule force spectroscopy or single-molecule fluorescence methods provide access to the length and timescales relevant to biomolecular interactions. More importantly, they allow us to see more diverse behaviors otherwise hidden by ensemble averaging. This however does not mean that more traditional ensemble methods of investigating biomolecular interactions such as isothermal titration calorimetry (ITC) or surface plasmon resonance (SPR) should be disregarded. It is rather a combination of all these tools that enables us to gain a better understanding of how the building blocks of life interact.

Within this thesis, the cellulosome, a large extracellular network of proteins expressed by certain bacteria, that aids in the degradation of cellulose, is used as a model system for studying protein-protein interactions. It consists of a multitude of scaffold-like proteins, enzymes and cellulose-binding domains, interconnected by a class of receptor-ligand pairs called cohesin-dockerin.

**Publication 3.1** investigates the remarkable mechanical strength of a cohesindockerin pair, which non-covalently anchors cellulose-binding proteins to the bacterial cell wall of the cellulose degrading bacterium *Ruminococcus flavefaciens*. Steered allatom molecular dynamic simulations revealed that - contrary to intuition - the surface contact area between the binding partners increases under load.

**Publication 3.2** provides further mechanistic insight into how this protein complex is able to withstand the remarkable forces measured. To study how force propagates through the complex, the cohesin was anchored in two geometries, via the carboxylterminus (which in nature is coupled to the bacterium's cell wall), as well as via the amine-terminus. A substantial decrease in force resilience in the non-native anchoring geometry was revealed. By analyzing all-atom molecular dynamics simulations with thermodynamic fluctuation theory, force propagation pathways could be determined. This helped to elucidate the mechanics that stabilize the complex, revealing pathways with strong components normal to the pulling direction. Based on the knowledge gained from and with the tools developed for these two publications, a homologous receptor-ligand pair was investigated in **publication 3.3**. It non-covalently anchors a large multi-domain cellulosome scaffold to *R. flavefaciens*. From MD-simulations it was predicted to be even stronger than the system characterized in publications 3.1 and 3.2. This was confirmed experimentally, whereas the affinity of this complex was found to be lower than for the weaker complex. This highlights the importance of distinguishing between the force-induced unbinding along a certain reaction coordinate, and the sum of all unbinding pathways, which give rise to the off-rate of a complex in the absence of force.

Driven by these results, we sought to obtain a more complete understanding of the role these complexes serve in nature. The contrast between very high force resilience and moderate affinities in the nano- to micro-molar range (see publication 3.3) as determined by ITC prompted the question which kinetic rate constants  $k_{on}$ and  $k_{off}$  give rise to them. While low off-rates would allow the bacterium to hold on to the nano-machinery it expressed, higher off-rates might provide for a more flexible composition of the cellulosome. As initial measurements conducted via SPR (using non-specific amine coupling) were inconsistent, a novel coupling strategy for surface-based affinity assays was developed, detailed in **publication 3.4**. Using a biolayer interferometer (BLI) as measurement platform, three different enzymatic coupling methods were used to covalently and site-specifically immobilize a protein of interest. Making use of the same short tags employed in the SMFS experiments of this thesis, this circumvents problems that arise from the conventionally employed non-specific coupling via primary amines. This way, cohesin E(R.f.) was functionally immobilized on BLI sensors in both geometries studied in publication 3.2. It was found that there was no apparent difference between immobilization geometries in BLI, *i.e.* in the absence of force, in contrast to under force. These findings underline the importance and influence of the force loading geometry and the resulting force propagation in biomolecular complexes.

**Publication 3.5** investigates the unfolding behavior of a cohesin domain from *Acetovibrio cellulolyticus*. Verdorfer et al. [1, 2] found some members of this family of cohesins to exhibit bimodal unfolding force distributions, which became unimodal upon ligand binding. To assess whether this was caused by a heterogeneity in folding states or unfolding pathways, a combination of constant loading rate and constant speed SMFS was conducted. This allowed recording of unfolding forces over a very wide range of loading rates. The shape of the unfolding force distributions changed in a force loading rate-dependent manner - consistent with the domain switching between at least two different folded states within the experimental timescale. Monte Carlo simulations modelling a simple system consisting of two folded states, which interchange in a non force-dependent manner, were qualitatively consistent with the recorded SMFS data. While this behavior has been described before in a theoretical study by Pierse and Dudko [3], this is the first time such a behavior was shown for a protein domain by SMFS. Is to publications 3.6 and 3.7.

**Publication 3.6** presents a streamlined approach to expressing proteins of interest via *in vitro* transcription and translation reactions and immobilizing them on a sample surface spatially separated via silicon microwell masks. Subsequently, CttA-XDoc - which is part of the high-force receptor-ligand system described in publications

3.1, 3.2 and 3.3 - is enzymatically ligated to the protein of interest. This allows for high specificity and throughput, while eliminating the need to express the protein of interest as a fusion protein with the pulling handle.

**Publication 3.7** replaces the in SMFS commonly used polyethylene-glycol (PEG) linkers with elastin-like polymers (ELP), improving three aspects of protein based SMFS. As their amino acid sequence is genetically encoded, they can be designed to have terminal sequences recognized by peptide-ligases such as Sortase A. Not only does this provide a site-specific means of anchoring proteins, it also assures that only full-length ELPs are functionalized, ensuring linker monodispersity. Using ELPs as linkers also results in a homogeneous system, where the whole stretched polymer is a peptide. This is especially advantageous for measurements where forces of more than  $\sim 100 \, \text{pN}$  are reached, as PEG exhibits a force-dependent phase transition around that force.

Refinements to AFM instrumentation as well as careful data analysis allowed the investigation of a possible mechano-activation of focal adhesion kinase (FAK) (**publication 3.8**). FAK is a signaling molecule involved in cell adhesion and migration. By combining *in silico* and *in vitro* SMFS, it is shown that mechano-activation can below the unfolding force of FAK. This is consistent with a proposed force - or distance - induced conformational change of FAK that activates it.

Chapter 1. Introduction

# 2 Context

### 2.1 The Cellulosome

First described by R. Lamed and E. A. Bayer in a 1983 study [4], the cellulosome is a bacterial multi-domain, extracellular protein network capable of efficent plant fiber degradation. It is organized by, and constructed around a receptor-ligand interaction termed cohesin-dockerin. A cellulosome typically contains a variety of different enzymes and adhesion proteins fused to dockerin domains, which non-covalently bind to the cohesin-bearing cellulosome backbones. These backbones are multi-domain proteins called scaffoldins, which usually contain several cohesin domains, separated by peptide linkers. Scaffoldins also may contain Dockerin domains themselves, allowing for an intricate stacking of different scaffoldins. This self-assembling nature allows bacteria to fine tune the catalytic activity to their substrate. But it also makes it an excellent toolbox for the thriving field of synthetic biology, as the cellulosome provides many orthogonal cohesin-dockerin pairs, which differ in affinity, specificity, robustness against force, and size. Recently, David et al. have discovered a member of the human gut microbiome, *Ruminococcus champanellensis* [5], to express cellulosomes.

### 2.2 Affinity Determination and Kinetic Rate Measurements

The correct determination of intermolecular affinities is essential for many areas of research. The perfect technique to determine binding behavior should not suffer from unspecific surface binding effects, require little material, be sensitive enough to detect low molecular-weight changes, provide kinetic rate data, and be label-free. Because no actually existing technique satisfies all of these qualities, some systems may require more than one technique to reliably characterize an interaction. Quantitative methods to determine intermolecular affinities can be divided by three criteria - is one of the binding partners immobilized to a surface, can they resolve kinetic rates (or do they just give access to the equilibrium binding constant), and are they label-free. Within this work, isothermal calorimetry (ITC), as well as biolayer interferometry (BLI), were employed to study the binding behavior of receptor-ligand systems in the absence of force.

### 2.2.1 Isothermal Titration Calorimetry

While ITC has the advantage of being a truly label-free technique, requiring no molecular modifications and being largely unaffected by suface adsorpion, it does not provide kinetic rate data (with the exception of the recently developed (but debated) kinetic ITC [6]). It works by tracking the heat differential required to keep temperature constant in two liquid cells, one of which contains reference buffer, and one which contains sample to be analyzed. Consequently, ligand binding in the sample cell needs to produce a sufficiently large temperature change to be observable. Therefore, the change in Gibbs free energy upon binding  $\Delta G = \Delta H - T\Delta S$  must have a sufficient enthalpic component to be observable by ITC. At the same time, the binding sigmoid's sharpness is a function of receptor concentration in the cell. The cell concentration needs to be tuned such that sufficient data points are acquired around the sigmoid's turning point, which provides an upper limit for concentrations that can be chosen. Therefore, very high affinity binders (roughly  $K_d < 1$  nM) cannot be reliably probed via ITC (although this limit will shift with cell volume).

### 2.2.2 Biolayer Interferometry

Observing ligand binding in both, biolayer interferometry as well as surface plasmon resonance (SPR), is based on changes in optical density at a sensor surface upon binding. Surface based techniques, such as BLI, SPR or quartz crystal microbalance often provide high sensitivity. They are suitable to study high-affinity binders, as the desired signal correlates to the occupancy of the surface-bound receptors. This comes at the cost of having to manage unspecific surface binding effects. In order to mitigate these issues, great effort has gone into the development of surface coatings designed to minimize unspecific adsorption [7]. Nevertheless, the experimentalist often has to painstakingly optimize buffer conditions and passivation reagents for every new interaction investigated. Referencing the measurement signal with a second signal acquired from a non-functionalized surface does not provide a means to fully control for unspecific surface adsorption, because the ligand concentration at a sensor surface drastically increases upon specific receptor binding. This often causes unspecific surface binding to increase with the occupancy of surface-coupled receptors. A second challenge for surface based assays is the surface immobilization of the molecules to be studied. While SPR and BLI often is termed label-free, surface immobilization often involves antibody-capture methods or covalent modification to primary amines. A possible solution for this issue is presented in publication 3.4. Here, enzyme-based site-specific ligation reactions are used to covalently functionalize BLI sensors under mild reaction conditions.

### ANALYZING KINETIC RATE MEASUREMENTS

Kinetic Rates for a one-to-one binding scenario between receptor and ligand can be extracted in a straight-forward manner by fitting equations 2.1 and 2.2. Here,  $y_{max}$  represents the maximal binding signal, that can be achieved when all receptors are occupied with a ligand, and  $y_{off}$  is the binding signal at the beginning of ligand dissociation. Ideally, for several concentrations of ligand, association and dissociation cycles are recorded, and the fitted globally. This means that  $k_{on}$ ,  $k_{off}$ , and  $y_{max}$  are

optimized for data from all concentrations simultaneously. However, if the sensor's binding capacity degrades over the course of the measurements, it might become necessary to optimize  $y_{max}$  for each concentration separately. Figure 2.2 gives an example of BLI data showing binding and unbinding of a receptor-ligand interaction under variation of ligand concentations.

$$y = (y_0 - y_{off}) * e^{-k_{off} * (t - t_{0, dissociation})}$$
(2.1)

$$y = y_{max} \frac{1}{1 + \frac{k_{off}}{k_{on}*[c]}} * \left(1 - e^{-([c]k_{on} + k_{off})*(t - t_{0,association})}\right)$$
(2.2)

If full ligand dissociation cannot be achieved after each measurement cycle (for example because of very small off-rates), one option is to force ligand unbinding by altering buffer conditions. However, this so called regeneration procedure needs to be established for each interaction anew, and it should be ensured that binding curves are repeatable after regeneration. Sometimes, no viable regeneration conditions can be found. In these cases, another measurement protocol can be employed, that does not require complete ligand dissociation before moving on to another ligand concentration. Figure 2.3 gives an example for this measurement protocol. In order to correctly analyze such data, one needs to account for the already occupied receptor population at the beginning of each association cycle. To do so, equation 2.2 solved for



**Figure 2.1.** Illustration of affinity measurements via BLI. (**A**) Non-specific immobilization of the receptor on the sensor in different geometries as a result of several accessible amine-groups. (**B**) Specific and site-directed immobilization of a receptor to a sensor. All receptors are homogeneously orientated. The red arrows in (A) represent different binding geometries with possibly different kinetics, whereas specific attachment (B) provides a uniform population of binders. (**C**) The principle of a BLI kinetic experiment. A receptor-functionalized sensor is immersed into a ligand solution. The increasing signal shows binding of the ligand. When the sensor signal has reached a steady state, the rates of ligand association and dissociation are equal—the system has reached equilibrium. The sensor is then moved to a buffer solution, the receptor starts to dissociate and the detected signal decreases again. Figure adapted from publication 3.4.

the time t, yielding equation 2.3. This can be used to calculate a 'time offset' for every association step, by inserting parameters  $k_{on}$ ,  $k_{off}$  and the current concentration [c]. The so calculated  $t_{offset}$  can then be inserted into equation 2.4.

$$t_{offset} = \frac{-ln\left(1 - \left(\frac{y_{start}}{y_{max}*[c]}\right) * \left(\frac{k_{off}}{k_{on}} + [c]\right)\right)}{k_{on}*[c] + k_{off}}$$
(2.3)

$$y = y_{max} \frac{1}{1 + \frac{k_{off}}{k_{on}*[c]}} * \left(1 - e^{-([c]k_{on} + k_{off})*(t - t_{0,association} + t_{offset})}\right)$$
(2.4)

Equations 2.3 and 2.3 can now again be simultaneously fitted to the data under variation of  $k_{on}$ ,  $k_{off}$ , and  $y_{max}$ .



**Figure 2.2.** This figure illustrates a typical series of kinetic rate measurements via BLI. After each dissociation, sufficient time is given for all ligand to dissociate. (Data from publication 3.4).



**Figure 2.3.** An example for a so called kinetic titration series, recorded by BLI. Binding of a receptor-ligand interaction is measured under variation of ligand concentration. While the time given for ligand dissociation between association streps needs to be sufficient for correct determination of the off-rate, complete ligand dissociation does not need to be reached in this measurement protocol. (Data taken from publication 3.4).

### 2.3 Single-Molecule Force Spectroscopy

Single-molecule force spectroscopy today encompasses a variety of techniques, capable of studying molecular force responses over a wide range of forces. While forces below ~ 25 pN are best studied by magnetic or optical tweezer assays, AFM-based SMFS is uniquely suited to study systems in the force range from ~ 50 pN up to nanonewtons [8]. This makes it an ideal technique to study the mechanical properties of cellulosomal components, which can differ strongly in their capability to withstand forces. For example, a cohesin-dockerin pair from *Clostridium perfringens* [9] was found to unbind at approximately ~ 60 pN. These unbinding forces are surpassed more than ten-fold by the cohesin-dockerin systems studied within this thesis (using similar pulling velocities and force probes). AFM based SMFS is most commonly conducted in three different measurement modes: Either the distance between cantilever and surface is increased with constant velocity ("constant speed"), or the force is either increased linearly over time ("force ramp") or kept constant ("force clamp").



**Figure 2.4.** Illustration showing the experimental arrangement of a receptor-ligand based SMFS experiment, as well as example curves for the constant speed and constant loading rate measurement modes. (A) Schematic illustration of the experimental setup. The Ctta XDoc:CohE (*R.f.*) receptor ligand system was used to specifically and reliably probe Coh3 (*A.c.*), which was expressed as a fusion protein with XDoc. After each measurement, the surface is moved laterally to avoid probing the same Coh3 twice and the retraction speeds and loading rates are varied. (**B**, **C**) Example traces showing Coh3 unfolding in constant loading rate and constant speed mode. Unfolded peptide behaves as an entropic spring, as fits of a worm-like chain model (red lines) to both stretches in (C) illustrate. (**D**) Contour length transformation of the constant loading rate trace in (B), illustrating how Coh3 unfolding events are assigned if they match the expected contour length increment of 46 nm. Figure adapted from publication 3.5.

# 2.3.1 Measurement Modes

### CONSTANT SPEED

Because of the ease of implementation, this is by far the most commonly used measurement protocol. Accurate distance control is easily achieved using piezo-actuators with additional position sensors, and by controlling the distance, the duration of a measurement cycle is limited by the time it takes for the retraction distance to approach the contour length of the stretched molecule. SMFS experiments usually rely on linkers to mitigate surface interactions, which have a force response similar to an entropic spring. Therefore, force increases with distance in a non-linear fashion.

## Force Ramp

AFM-based force ramp SMFS requires a much more carefully designed feedback mechanism in comparsion to constant speed SMFS. Abrupt changes in force proceeding e.g. protein unfolding events demand a well tuned control loop in order to restore the set force quickly, without overshooting it. Piezoelectric elements that exhibit small hystereses are favorable, allowing purpose-build instruments to compensate for sudden changes in force within ~ 10 ms [10]. Importantly, force ramp mode allows to record data at comparably low loading rates, as demonstrated in publication 3.5. This effect is especially pronounced for high-force complexes, as in constant speed mode force increases steeply, when the extension approaches the contour length of the system probed.

## Force Clamp

In contrast to constant speed or constant loading rate mode, force clamp SMFS provides direct access to life time statistics. This is essential when investigating phenomena like catch-bonding [11], where the defining criterium is a bond lifetime that increases with force, setting it apart from common slip-bonds. However, force clamp SMFS becomes challenging when an investigated system consists of several components with differing life times.

## 2.3.2 Kinetic Rate Models and Polymer Elasticity

In an effort to quantitatively assess force-biased molecular unfolding or unbinding, kinetic rate models of varying complexity have been developed [3, 12–14]. The most commonly employed model was developed by Bell, Evans, and Schulten. It models a unbinding or unfolding process as the thermal crossing of a one-dimensional free energy barrier, separating two states. Equation 2.5 gives an expression for the off-rate from the bound state. In the Bell picture [15],  $v_0$  is the microscopic attempt frequency, that can be combined with the energy barrier height  $\Delta G$  to form the zero-force off-rate  $k_{off}^0$  along the reaction coordinate  $\Delta x$ . The application of force now reduces the energy barrier height, thus increasing the off-rate.

$$k_{off}(F) = \nu_0 \left( \frac{F\Delta x}{k_B T} - \frac{\Delta G}{k_B T} \right) = k_{off}^0 exp \left( \frac{F\Delta x}{k_B T} \right)$$
(2.5)

To be applicable by the experimentalist, theory must be relatable to quantities that can be reliably extracted from measured data. Under the assumption of a constant force loading rate  $\dot{F}$ , analytical expressions for the barrier crossing forces (eq. 2.6) as well as for the most probable crossing force (eq. 2.7) can be derived.

$$p(F) = \frac{k_{off}^0}{\dot{F}} \exp\left(\frac{\Delta xF}{k_BT} - k_{off}^0 k_BT \frac{\exp\left(\frac{\Delta xF}{k_BT}\right) - 1}{\Delta x\dot{F}}\right)$$
(2.6)

$$F_{most\_probable}(\dot{F}) = \frac{k_B T}{\Delta x} \log \left( \frac{\dot{F} \Delta x}{k_{off}^0 k_B T} \right)$$
(2.7)

LINKER CORRECTIONS TO THE BELL-EVANS MODEL

While these expressions are directly applicable to data recorded in force ramp mode, constant speed experiments will usually not provide constant force loading rates due to the elastic response of linker molecules. Several groups have sought to quantify the resulting deviations [16, 17], and to provide experimentalists with theory consistent with non-constant loading rates that result from entropic springs. Ray et al. have derived an analytical expression for a bond rupture distribution p(F), that is derived for the force-distance response of a worm-like chain [18]. Considering the forces reached in some receptor-ligand systems (e.g. publication 3.3), this approach can be used to derive a similar expression for a polymer model that describes the elastic response of a peptide linker in the high-force limit more accurately. Using the polymer elasticity response for high forces from Livadaru et al. [19],  $x = L_c \left(1 - \frac{k_b T}{cFb}\right)$  and using  $F_b = \frac{k_b T}{b}$  as a lower bound, the bond survival probability S(t) can be calculated analogous to Ray et al. [18]:

$$S(t) = \exp\left[-\int \frac{1}{f} k_{off}^{0} \exp(f/F^{\dagger}) df\right]$$
  
=  $\exp\left[-k_{off}^{0} \left(\int_{0}^{F} \frac{1}{k_{c}v} e^{f/F^{\dagger}} df + \int_{F_{b}}^{F} \frac{F_{b}}{f^{2}} e^{f/f^{\dagger}} df\right)\right]$   
=  $\exp\left[-\frac{k_{off}^{0} F^{\dagger}}{v k_{c}} \left(e^{F/F^{\dagger}} \left(1 - \frac{L_{c}F_{b}k_{c}}{cF^{\dagger}F}\right) + \frac{L_{c}F_{b}k_{c}}{cF^{\dagger^{2}}} \left(\frac{F^{\dagger}}{F_{b}} e^{F_{b}/F^{\dagger}} + \operatorname{Ei}\left(\frac{F}{F^{\dagger}}\right) - \operatorname{Ei}\left(\frac{F_{b}}{F^{\dagger}}\right)\right) - 1\right)\right]$   
(2.8)

Here,  $k_c$  denotes the cantilever spring constant, Ei the exponential integral, and v the retraction velocity. By performing the derivative  $-\frac{dS(F)}{dF} = p(F)$ , we get Eq.2.9.

$$p(F) = \frac{k_{off}^0}{vk_c} \left( 1 + \frac{L_c F_b k_c}{cF^2} \right) \exp\left[ \frac{F}{F^{\dagger}} - \frac{k_{off}^0}{vk_c} \left( F^{\dagger} e^{F/F^{\dagger}} \left( 1 - \frac{L_c F_b k_c}{cF^{\dagger}F} \right) + \frac{k_c L_c}{c} e^{F_b/F^{\dagger}} + \frac{L_c F_b k_c}{cF^{\dagger}} \left( \operatorname{Ei}\left(\frac{F}{F^{\dagger}}\right) - \operatorname{Ei}\left(\frac{F_b}{F^{\dagger}}\right) \right) - F^{\dagger} \right) \right]$$
(2.9)

Part II Results

# 3 Research Articles

# 3.1 Load-Tightening of a Protein-Protein Binding Interface

This publication explores the mechanical resilience of a receptor-ligand pair, which anchors the cellulose-adhesion machinery of a cellulose degrading bacterium to its cell wall. Its ability to withstand exceptionally high forces, which may have evolved to withstand the turbulent environment of the bovine rumen, is explored with a combined approach of AFM based SMFS and all-atom MD simulations. Comparing how the contact surface area between the recetor-ligand pair changes when under force, we surprisingly find that it increases with force. My contribution to this work includes the preparation of Cohesin E in its native anchoring geometry as well as performing SMFS experiments and analyzing the recorded data.

# Ultrastable Cellulosome-adhesion Complex Tightens Under Load

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### ARTICLE

ellulosomes are protein networks designed by nature to degrade lignocellulosic biomass<sup>1</sup>. These networks comprise intricate assemblies of conserved subunits including catalytic domains, scaffold proteins, carbohydrate binding modules (CBMs), cohesins (Cohs), dockerins (Docs) and X-modules (XMods) of unknown function. Coh:Doc pairs form complexes with high affinity and specificity<sup>2</sup>, and provide connectivity to a purvised of collucordin provedre with varying form complexes with light annuly and specificity, and provide connectivity to a myriad of cellulosomal networks with varying Coh:Doc network topology<sup>3–5</sup>. The most intricate cellulosome known to date is produced by *Ruminococcus flavefaciens* ( $R_{f}$ )<sup>67</sup> and contains several primary and secondary scaffolds along with over 220 Doc-bearing protein subunits<sup>8</sup>. The importance of cellulolytic enzymes for the production of

renewable fuels and chemicals from biomass has highlighted an urgent need for improved fundamental understanding of how cellulosomal networks achieve their impressive catalytic activity<sup>9</sup>. Two of the mechanisms known to increase the catalytic activity<sup>9</sup> cellulosomes are proximity and targeting effects<sup>10</sup>. Proximity refers to the high local concentration of enzymes afforded by incorporation into nanoscale networks, while targeting refers to specific binding of cellulosomes to substrates. Protein scaffolds and CBM domains are both critical in this context as they mediate interactions between comparatively large bacterial cells and cellulose particles. As many cellulosomal habitats (for example, cow rumen) exhibit strong flow gradients, shear forces will accordingly stress bridging scaffold components mechanically in vivo. Protein modules located at stressed positions within these networks should therefore be preselected for high mechanostability. However, thus far very few studies on the mechanics of carbohydrate-active proteins or cellulosomal network components have been reported<sup>11</sup>. In the present study we sought to identify cellulosomal network

junctions with maximal mechanical stability. We chose an XMod-Doc:Coh complex responsible for maintaining bacterial adhesion to cellulose in the rumen. The complex links the *R. flavefaciens* cell wall to the cellulose substrate via two CBM domains located at the N-terminus of the CttA scaffold, as shown in Fig. 1a. The

crystal structure of the complex solved by X-ray crystallography  $^{12}$  is shown in Fig. 1b. XMod-Doc tandem dyads such as this one are a common feature in cellulosomal networks. Bulk biochemical assays on XMod-Docs have demonstrated that XMods improve assays on AMode Docs have demonstrated that AMode improve Doc solubility and increase biochemical affinity of Doc:Coh complex formation<sup>13</sup>. Crystallographic studies conducted on XMod-Doc:Coh complexes have revealed direct contacts between XMods and their adjacent Docs<sup>12,14</sup>. In addition, many XMods (for example, PDB 2B59, 1EHX, 3PDD) have high  $\beta$ -strand content to feld with  $N_{\rm end}$  of the mining the mode for the content and fold with N- and C-termini at opposite ends of the content and toid with N- and C-termini at opposite ends of the molecule, suggestive of robust mechanical clamp motifs at work<sup>15,16</sup>. These observations all suggest a mechanical role for XMods. Here we perform AFM single-molecule force spectroscopy experiments and steered molecular dynamics simulations to understand the mechanostability of the XMod Doc:Coh cellulosomal ligand-receptor complex. We conclude that the high mechanostability we observe originates from molecular mechanisms, including stabilization of Doc by the adjacent XMod domain and catch bond behaviour that causes the complex to increase in contact area on application of force.

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### **Results and Discussion**

Single-molecule experiments. We performed single-molecule force spectroscopy (SMFS) experiments with an atomic force miscroscope (AFM) to probe the mechanical dissociation of XMod-Doc:Coh. Xvlanase (Xvn) and CBM fusion domains on the XMod-Doc and Coh modules, respectively, provided identifiable unfolding patterns permitting screening of large data sets of force-distance curves<sup>17–19</sup>. Engineered cysteines and/or peptide tags on the CBM and Xyn marker domains were used to covalently immobilize the binding partners in a site-specific manner to an AFM cantilever or cover glass via poly(ethylene glycol) (PEG) linkers. The pulling configuration with Coh-CBM immobilized on the cantilever is referred to as configuration I, as shown in Fig. 1c. The reverse configuration with Coh-CBM on the cover glass is referred to as configuration II. In a typical



Figure 1 | System overview. (a) Schematic of selected components of the R. flavefaciens cellulosome. The investigated XMod-Doc:Coh complex responsible for maintaining bacterial adhesion to cellulose is highlighted in orange. (b) Crystal structure of the XMod-Doc:Coh complex. Ca<sup>2+</sup> ions are shown as orange spheres. (c) Depiction of experimental pulling configuration I, with Coh-CBM attached to the cantilever tip and Xyn-XMod-Doc attached to the glass surface

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experimental run we collected about 50,000 force extension traces from a single cantilever. We note that the molecules immobilized on the cantilever and glass surfaces were stable over thousands of pulling cycles.

We sorted the data by first searching for contour length increments that matched our specific xylanase and CBM fingerprint domains. After identifying these specific traces (Fig. 2a), we measured the loading rate dependency of the final Doc:Coh ruptures based on bond history. To assign protein subdomains to the observed unfolding patterns, we transformed the data into contour length space using a freely rotating chain model with quantum mechanical corrections for peptide backbone stretching (QM-FRC, Supplementary Note 1, Supplementary Fig. 1)<sup>20,21</sup>. The fit parameter-free QM-FRC model describes protein stretching at forces >200 pN more accurately than the commonly used worm-like chain (WLC) model<sup>20,22</sup>. The resulting contour length histogram is shown in Fig. 2b. Peak-to-peak distances in the histogram represent contour length per stretched amino acid of 0.365 nm and accounting for the folded length of each subdomain, we compared the observed increments to the polypeptide lengths of individual subdomains of the Xyn-XMod-Doc and Coh-CBM fusion proteins. Details on contour length estimates and domain assignments are shown in Supplementary Table 1. Unfolding patterns in configuration I showed PEG stretching

Unfolding patterns in configuration 1 showed PEG stretching followed by a three-peaked Xyn fingerprint (Fig. 1a, top trace, green), which added 90 nm of contour length to the system. Xyn unfolding was followed by CBM unfolding at ~150 pN with 55 nm of contour length added. Finally, the XMod-Doc:Coh complex dissociated at an ultra-high rupture force of ~600 pN. The loading rate dependence of the final rupture event for curves of subtype 1 is plotted in Fig. 2c (blue). The measured complex rupture force distributions are shown in Supplementary Fig. 2.

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Less frequently (35–40% of traces) we observed a two-step dissociation process wherein the XMod unfolded before Doc:Coh rupture as shown in Fig. 2a (middle trace, orange). In these cases, the final dissociation exhibited a much lower rupture force (~300 pN) than the preceding XMod unfolding peak, indicating the strengthening effect of XMod was lost, and XMod was no longer able to protect the complex from dissociation at high force. The loading rate dependency of Doc:Coh rupture occurring immediately following XMod unfolding is shown in Fig. 2c (grey).

In configuration II (Fig. 2a, bottom trace), with the Xyn-XMod-Doc attached to the cantilever, the xylanase fingerprint was lost after the first few force extension traces acquired in the data set. This indicated the Xyn domain did not refold within the timescale of the experiment once unfolded, consistent with prior work<sup>17,18</sup>. CBM and XMod unfolding events were observed repeatedly throughout the series of acquired force traces in both configurations I and II, indicating these domains were able to refold while attached to the cantilever over the course of the experiment.

We employed the Bell-Evans model<sup>23</sup> (Supplementary Note 2) to analyse the final rupture of the complex through the effective distance to the transition state ( $\Delta x$ ) and the natural off-rate ( $k_{off}$ ). The fits to the model yielded values of  $\Delta x = 0.13$  nm and  $k_{off} = 7.3 \times 10^{-7} \text{s}^{-1}$  for an intact XMod, and  $\Delta x = 0.19$  nm and  $k_{off} = 4.7 \times 10^{-4} \text{s}^{-1}$  for the 'shielded' rupture following XMod unfolding (Fig. 2c). These values indicate that the distance to the transition state is increased following XMod unfolding, reflecting an overall softening of the binding interface. Distances to the transition state observed for other ligand–receptor pairs are typically on the order of ~0.7 nm (ref. 17). The extremely short  $\Delta x$  of 0.13 nm observed here suggests that mechanical unbinding for this complex is highly coordinated. We further analysed the unfolding of  $\Delta x = 0.15$  and  $k_{off} = 2.6 \times 10^{-6} \text{s}^{-1}$ . The loading



Figure 2 | Experimental SMFS unfolding traces. (a) Unfolding fingerprints from pulling configuration I (curves 1 & 2) and configuration II (curve 3). The QM-FRC model (dashed lines) was used to estimate the contour lengths of the unfolded modules. (b) Contour length histogram obtained from 127 force extension traces (Config. I). The peak-to-peak increments correspond to Xyn, CBM and XMod amino-acid sequence lengths. (c) Dynamic force spectra for the final Doc:Coh complex rupture peaks obtained from 2,122 force-extension traces. The blue points show Doc:Coh ruptures that occurred with an intact XMod, while grey points show ruptures immediately following XMod unfolding. Black circles and diamonds represent the most probable rupture force/loading rate obtained by Gaussian fitting at each pulling speed. Error bars are ±1 s.d. Dashed lines are least square fits to the Bell-Evans model.

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unfold in several substeps starting at  $\sim$  400 pN. Snapshots at different time steps are labelled I-V and are shown in (c). Steps IV and V are shown at smaller scale. (d) Average number of hydrogen bonds between Doc:Coh (black) and XMod-Doc (red). XMod-Doc contact is dominated by the insert sequences 1-3. (e) Root mean squared deviation (RMSD) of Doc (black) and Coh (red). key residues in contact throughout the entire rupture process and Among the 223 Doc sequences from *R. flavefaciens*, six

key residues in contact throughout the entire rupture process and particularly immediately before rupture. These residues are presented in Fig. 3a,c,d and Supplementary Figs 5,6. The simulation results clearly reproduced key hydrogen bonding contacts previously identified<sup>12</sup> as important for Doc:Coh recognition (Supplementary Fig. 5).

The main interacting residues are shown in Fig. 3a,b. Both Coh and Doc exhibit a binding interface consisting of a hydrophobic centre (grey) surrounded by a ring of polar (green) and charged residues (blue, positive; red, negative). This residue pattern suggests the hydrophilic side chains protect the interior hydrophobic core from attack by water molecules, compensating for the flat binding interface that lacks a deep pocket. The geometry suggests a penalty to unbinding that stabilizes the bound state. Further, we analysed the contact surface areas of interacting residues (Fig. 3b–e). The total contact area was found to increase due to rearrangement of the interacting residues when the complex is mechanically stressed, as shown in Fig. 3e and Supplementary Movie 2. Doc residues upon mechanical loading, resulting in increased stability and decreased accessibility of water into the hydrophobic core of the bound complex (Fig. 3b). These results suggest that a catch bond mechanism is responsible for the remarkable stability<sup>33</sup> under force and provide a molecular mechanism which the XMod–Doc:Coh complex uses to summon mechanical strength when needed, while still allowing relatively fast assembly and disassembly of the complex at equilibrium. The residues that increase most in contact area (Fig. 3c,d) present promising candidates for future mutagenesis studies.

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Among the 223 Doc sequences from *R. flavefaciens*, six subfamilies have been explicitly identified using bioinformatics approaches<sup>8</sup>. The XMod–Doc investigated here belongs to the 40-member Doc family 4a. A conserved feature of these Doc modules is the presence of three sequence inserts that interrupt the conserved duplicated F-hand motif Doc structure. In our system, these Doc sequence inserts make direct contacts with XMod in the crystallized complex (Fig. 1) and suggest an interaction between XMod and Doc that could potentially propagate to the DocCoh binding interface. To test this, an independent simulation was performed to unfold XMod (Fig. 4). The harmonic restraint was moved to the C-terminus of XMod only, while leaving Doc and Coh unrestrained. The results (Fig. 4b) showed XMod unfolded at forces slightly higher than but similar to the XMod–Doc:Coh complex rupture force determined from the standard simulation at the same pulling speed. This suggested XMod unfolding was observed in  $\sim$ 35–40% of traces. Furthermore, analysis of the H-bonding between Doc and XMod unfolded, dominated by contact loss between the three Doc insert sequences and XMod. Interestingly, XMod unfolding contact loss between Doc and Coh at a later stage ( $\sim$ 200 ns) well after XMod had lost most of its contact with Doc, even though no force was being applied across

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As shown in Fig. 4e, the root mean squared deviation (RMSD) of Doc increased throughout the simulation as XMod unfolded. Coh RMSD remained stable until it started to lose H-bonds with Doc. Taken together this suggests that, as XMod unfolded, Coh and Doc became more mobile and lost interaction strength, potentially explaining the increase in  $\Delta x$  from 0.13 to 0.19 nm on unfolding of XMod in the experimental data sets. Apparently the XMod is able to directly stabilize the Doc:Coh interface, presumably through contact with Doc insert sequences that then propagate this stabilizing effect to the Doc:Coh binding interface.

In summary, we investigated an ultrastable XMod-Doc:Coh complex involved in bacterial adhesion to cellulose. While previously the role of XMod functioning in tandem XMod-Doc dyads was unclear<sup>12,14</sup>, we show that XMod serves as a mechanical stabilizer and force-shielding effector subdomain in the ultrastable ligand-receptor complex. The Doc:Coh complex presented here exhibits one of the most mechanically robust protein-protein interactions reported thus far, and points towards new mechanically stable artificial multi-component biocatalysts for industrial applications, including production of second-generation biofuels.

### Methods

Site-directed mutagenesis. Site-directed mutagenesis of R. flavefaciens strain Site-directed mutagenesis. Site-directed mutagenesis of *R. flavefaciens* strain FD1 chimeric cellulosomal proteins. A pET28a vector containing the previously cloned *R. flavefaciens* CohE from ScaE fused to cellulose-binding module 3a (CBM3a) from *C. thermocellum*, and a pET28a vector containing the previously cloned *R. flavefaciens* XMod-Doc from the CrtAt scaFioldin fused to the XymT6 xylanase from *Geobacillus stearothermophilus*<sup>12</sup> were subjected to QuikChange mutagenesis<sup>34</sup> to install the following mutations: A2C in the CBM and T129C in the xylanase, respectively. For the construction of the native configuration of the CohE-CBM A2C fusion protein Gibson assembly<sup>35</sup> was used. For further analysis CohE-CBM A2C was modified with a QuikChange PCR<sup>36</sup> to replace the two cysteins (C2 and C63) in the protein will aalnine and serien (C2A and C635). All mutagenesis products were confirmed by DNA sequencing analysis.

The XynT6-XDoc T129C was constructed using the following primers: 

The CBM-CohE A2C was constructed using the following primers  $5'-ttaactttaagaaggagatataccatgtgcaatacaccggtatcaggcaatttgaag-3\\5'-cttcaaattgcctgataccggtgattgcacatggtatatctccttcttaaagttaa-3'$ 

The CohE-CBM C2A C63S was constructed using the following phosphorylated primers:

5'-ccgaatgccatggccaatacaccgg-3' 5'-cagaccttctggagtgaccatgctgc-3'

**Expression and purification of Xyn-XMod-Doc.** The T129C Xyn-XMod-Doc protein was expressed in *E. coli* B1.21 cells in kanamycin-containing media that also contained 2 mM calcium chloride, overnight at 16 °C. After harvesting, cells were lysed using sonication. The lysate was then pelleted, and the supernatant fluids were applied to a Ni-NTA column and washed with tris-buffered saline (TBS) buffer containing 20 mM imidazole and 2 mM calcium chloride. The bound protein was eluted using TBS buffer containing 250 mM imidazole and 2 mM calcium chloride. The solution was dialysed with TBS to remove the imidazole, and then concentrated using an Amicon centrifugal filter device and stored in 50% (v/v) glycerol at -20 °C. The concentrations of the protein stock solutions were determined to be ~5 mg ml<sup>-1</sup> by absorption spectrophotometry.

Expression and purification of Coh-CBM. The Coh-CBM C2A, C63S fusion Expression and purification of Coh-CBM. The Coh-CBM C2A, C635 fusion protein was expressed in  $E_{\rm cold}$  B21(DE3) RPL in knamycin and chlor-amphenical containing ZYM-5052 media<sup>37</sup> overnight at 22 °C. After harvesting, cells were lysed using sonication. The lysate was then pelleted, and the supernatant fluids were applied to a Ni-NTA column and washed with TBS buffer. The bound protein was eluted using TBS buffer containing 200 mM imidazole. Imidazole was removed with a polyacrylamide gravity flow column. The protein solution was concentrated with an Amicon centrifugal filter device and stored in 50% (v/v) glycerol at  $-80^\circ$ C. The concentrations of the protein solutions were determined to be ~5 mg ml<sup>-1</sup> by absorption spectrophotometry.

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Sample preparation. In sample preparation and single-molecule measurements calcium supplemented TBS buffer (Ca-TBS) was used (25 mM TRIS, 72 mM NaCl, 1 mM CaCl<sub>2</sub>, pH 7.2). Cantilevers and cover glasses were functionalized according to previously published protocols<sup>18,38</sup>. In brief, cantilevers and cover glasses were cleaned by UV-ozone treatment and piranha solution, respectively. Levers and glasses were silanized using (3-aminopropy)1-dimethyl-ethoxysilane (APDMES) to introduce surface amine groups. Amine groups on the cantilevers and cover glasses were subsequently conjugated to a \$LDa NH3-PEG-Mal linker in sodium borate buffer. Disulfide-linked dimers of the Xyn-XMod-Doc proteins were reduced for 2 h at room temperature using a TCEP disulfide reducing bead slurry. The protein/ bead mixture was rinsed with Ca-TBS measurement buffer, centrifuged at 850 r.c.f. for 3 min, and the supernatart was collected with a micropiette. Reduced proteins for 3 min, and the supernatant was collected with a micropipette. Reduced proteins were diluted with measurement buffer (1:3 (v/v) for cantilevers, and 1:1 (v/v) for cover glasses), and applied to PEGylated cantilevers and cover glasses for 1 h. Both cantilevers and cover glasses were then rinsed with Ca-TBS to remove unbound proteins and stored under Ca-TBS before force spectroscopy measurements. Site-specific immobilization of the Coh-CBM-ybbR fusion proteins to previously PEGylated cantilevers or coverglasses was carried out according to previously published protocols<sup>39</sup>. In brief, PEGylated cantilevers or coverglasses were incubated with Coenzyme A (CoA) (20 mM) stored in coupling buffer (50 mM sodium phosphate, 50 mM NaCl, 10 mM EDTA, pH 7.2) for 1 h at room sodium phosphate, 50 mM NaCl, 10 mM EDTA, pH 7.2) for 1 h at room constrained with Ca-TBS to remove unbound CoA. Coh-CBM-ybbR fusion proteins were then covalently linked to the CoA surfaces or levers by incubating with Sfp phosphopantetheim) transferase for 2 h at room 37°. Finally, surfaces or levers were subjected to a final rinse with Ca-TBS and stored under Ca-TBS before measurement. cover glasses), and applied to PEGylated cantilevers and cover glasses for 1 h. Both

**Single-molecule force spectroscopy measurements.** SMFS measurements were performed on a custom built AFM<sup>40</sup> controlled by an MFP-3D controller from Asylum Research running custom written Igor Pro (Wavemetrics) software. Cantilever spring constants were calibrated using the thermal noise/equiparition method<sup>41</sup>. The cantilever was brought into contact with the surface and withdrawn at constant speed ranging from 0.2 to 6.4 µm s<sup>-1</sup>. An x-y stage was actuated after each force-extension trace to expose the molecules on the cantilever to a new molecule at a different surface location with each trace. Typically 20,000–50,000 force-extension curves were obtained with a single cantilever in an experimental run of 18–24h. A low molecular density on the surface was used to avoid formation of multiple bonds. While the raw data sets contained a majority of unusable curves due to lack of interactions or nonspecific adhesion of molecules to the cantilever to a combination of automated data processing and manual classification by searching for contour length increments that matched the lengths of our specific protein fingerprint domains: Ym (~ 89 nm) and CBM (~ 56 nm). After identifying these specific traces, we measured the loading rate dependency of the final Doc:Coh ruptures based on bond history.

**Data analysis.** Data were analysed using previously published protocols<sup>17,18,22</sup>. Force extension traces were transformed into contour length space using the QM-FRC model with bonds of length b = 0.11 nm connected by a fixed angle  $\gamma = 41^{\circ}$  and assembled into barrier position histograms using cross-correlation. Detailed description of the contour length transformation can be found in Supplementary Note 1 and Supplementary Fig. 1. For the loading rate analysis, the loading rate at the point of rupture was extracted by applying a line fit to the force vs time trace in the immediate vicinity before the rupture peak. The loading rate was determined from the slope of the fit. The most probable rupture forces and loading rates are determined by applying Gaussian fits to histograms of rupture forces and loading rate at each pulling speed.

speed.

Molecular dynamics simulations. The structure of the XMod-Doc:Coh complex had been solved by means of X-ray crystallography at 1.97 Å resolution and is available at the protein data bank (PDB:41U3). A protonation analysis performed in VMD<sup>42</sup> did not suggest any extra protonation and all the amino-acid residues were simulated with standard protonation states. The system was then solvated, keeping also the water molecules present in the crystal structure, and the net charge of the protein and the calcium income are neutralized using redum atoms are compter of the protein and the calcium ions was neutralized using sodium atoms as counter ions, which were randomly arranged in the solvent. Two other systems, based on the aforementioned one, were created using a similar salt concentration to the one used in the experiments (75 mM of NaCl). This additional salt caused little or no change in SMD results. The overall number of atoms included in MD simulations varied from 300,000 in the majority of the simulations to 580,000 for the unfolding of the X-Mod.

of the X-Mod. The MD simulations in the present study were performed employing the NAMD molecular dynamics package<sup>45,44</sup>. The CHARMM36 force field<sup>45,46</sup> along with the TIP3 water model<sup>47</sup> was used to describe all systems. The simulations were done assuming periodic boundary conditions in the NpT ensemble with temperature maintained at 300 K using Langevin dynamics for pressure, kept at 1bar, and temperature coupling. A distance cut-off of 11.0Å was applied to short-range, non-bonded interactions, whereas long-range electrostatic interactions were

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treated using the particle-mesh Ewald (PME)<sup>48</sup> method. The equations of motion were integrated using the r-RESPA multiple<sup>48</sup> method. The equations of motion steps. The time step of integration was chosen to be 2 is for all simulations performed. Before the MD simulations all the systems were submitted to an energy minimization protocol for 1,000 steps. The first two nanoseconds of the simulations served to equilibrate systems before the production runs that varied from 40 to 450 ns in the 10 different simulations that were carried out. The equilibration step consisted of 500 ps of simulation where the protein backhone was restrained and 1.5 ns where the system was completely free and no restriction or force was applied. During the equilibration the initial temperature was set to zero and was constantly increased by 1 K every 100 MD steps until the desired temperature (300 K) was reached. To characterize the coupling between Doc and Coh, we performed SMD simulations<sup>49</sup> of constant velocity stretching (SMD-CV protocol) employing three different pulling speeds: 1.25, 0.625 and 0.25 A ns<sup>-1</sup>. In all simulations, SMD was employed by restraining the position of one end of the XMod-Doc domain harmonically (center of mass of ASN5), and moving a second restraint point, at the end of the Coh domain (center of mass of GLY210), with constant velocity in the end the knew next course in the mode where the time of the spring. The force ender the here next constant we increase the time of the time of the spring. The force

desired direction. The procedure is equivalent to attaching one end of a harmonic spring to the end of a domain and pulling on the other end of the spring. The force applied to the harmonic spring is then monitored during the time of the molecular dynamics simulation. The pulling point was moved with constant velocity along the z-axis and due to the single anchoring point and the single pulling point the system is quickly aligned along the z-axis. Owing to the flexibility of the linkers, this approach reproduces the experimental set-up. All analyses of MD trajectories were carried out employing VMD<sup>42</sup> and its plug-ins. Secondary structures were assigned using the Timeline plug-in, which employs STRIDE criteria<sup>50</sup>. Hydrogen bonds were assigned based on two geometric criteria for every trajectory frame saved: first, distances between acceptor and hydrogen should be <3.5 Å; second, the angle between hydrogen-donor-acceptor should be <30°. Surface contact areas of interacting residues were calculated employing Volarea<sup>31</sup> implemented in VMD. The area is calculated using a probe radius defined as an in *silico* rolling spherical probe that is screened around the area of Doc exposed to Coh and also Coh area exposed to Doc.

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### ARTICLE

### ARTICLE

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Additional information Accession codes: Plasmids used in this study are available through Addgene (https:// www.addgenco.rg) under the following accession codes: Xylanase-Xmodule-Dockerin: 60865; Cohesin-CBM: 60866.

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$\label{eq:module_states} \begin{array}{ c c c c c c } \hline \begin{tabular}{ c c c c c c c } \hline \begin{tabular}{ c c c c c c c c c c c c c c c c c c c$						
No. amino acids, $N_A$ 260 (378)159117205119Folded length, $L_F$ [nm]62722Expected increment, $\Delta L_E$ [nm]8956367242Observed increment, [nm]90 ± 4 $55 \pm 3$ $34 \pm 2$ upplementary Table 1: Domain assignment of observed contour length increments. The expectedntour length increment ( $\Delta L_E$ ) for each protein domain was calculated according to $\Delta L_E = N_A \cdot$ 365 nm - $L_F$ , where $L_F$ is the folded length, $N_A$ is the number of amino acids, and 0.365 nm <sup>2</sup> is thength per stretched amino acid. $L_F$ was measured for Xyn, CBM, and XDoc:Coh from PDB structures85, 1NBC, and 4IU3, respectively. For the Xyn domain, only amino acids located C-terminal of the129 mutation which served as attachment point are considered. Errors for the observed incrementsere determined from Gaussian fits to the combined contour length histogram shown in Fig. 2b.	Module	Xylanase	CBM	X-module	Cohesin	Dockerin
Folded length, $L_F$ [nm] 6 2 7 2 2 Expected increment, $\Delta L_E$ [nm] 89 56 36 72 42 Observed increment, [nm] 90 ± 4 55 ± 3 34 ± 2 upplementary Table 1: Domain assignment of observed contour length increments. The expected ntour length increment ( $\Delta L_E$ ) for each protein domain was calculated according to $\Delta L_E = N_A \cdot$ 365 nm - $L_F$ , where $L_F$ is the folded length, $N_A$ is the number of amino acids, and 0.365 nm <sup>2</sup> is the ngth per stretched amino acid. $L_F$ was measured for Xyn, CBM, and XDoc:Coh from PDB structures R85, 1NBC, and 4IU3, respectively. For the Xyn domain, only amino acids located C-terminal of the 129 mutation which served as attachment point are considered. Errors for the observed increments are determined from Gaussian fits to the combined contour length histogram shown in Fig. 2b.	No. amino acids, $N_A$	260 (378)	159	117	205	119
Expected increment, $\Delta L_E$ [nm] 89 56 36 72 42 Observed increment, [nm] 90 ± 4 55 ± 3 34 ± 2 – – – upplementary Table 1: Domain assignment of observed contour length increments. The expected income length increment ( $\Delta L_E$ ) for each protein domain was calculated according to $\Delta L_E = N_A \cdot 365 \text{ nm} - L_F$ , where $L_F$ is the folded length, $N_A$ is the number of amino acids, and 0.365 nm <sup>2</sup> is the ngth per stretched amino acid. $L_F$ was measured for Xyn, CBM, and XDoc:Coh from PDB structures 855, INBC, and 4IU3, respectively. For the Xyn domain, only amino acids located C-terminal of the 129 mutation which served as attachment point are considered. Errors for the observed increments are determined from Gaussian fits to the combined contour length histogram shown in Fig. 2b.	Folded length, $L_F$ [nm]	6	2	7	2	2
Observed increment, [nm] $90 \pm 4$ $55 \pm 3$ $34 \pm 2$ – – – upplementary Table 1: Domain assignment of observed contour length increments. The expected intour length increment ( $\Delta L_E$ ) for each protein domain was calculated according to $\Delta L_E = N_A \cdot$ $365 \text{ nm} - L_F$ , where $L_F$ is the folded length, $N_A$ is the number of amino acids, and 0.365 nm <sup>-2</sup> is the ngth per stretched amino acid. $L_F$ was measured for Xyn, CBM, and XDoc:Coh from PDB structures 385, INBC, and 4IU3, respectively. For the Xyn domain, only amino acids located C-terminal of the 129 mutation which served as attachment point are considered. Errors for the observed increments are determined from Gaussian fits to the combined contour length histogram shown in Fig. 2b.	Expected increment, $\Delta L_E$ [nm]	89	56	36	72	42
	contour length increment $(\Delta L_E)$ for ea ).365 nm – $L_F$ , where $L_F$ is the folded l ength per stretched amino acid. $L_F$ was IR85, INBC, and 4IU3, respectively. Fo C129 mutation which served as attachm were determined from Gaussian fits to th	ich protein d length, $N_A$ is measured foi r the Xyn do nent point ar he combined	omain wa the numb · Xyn, CB main, only e consider contour le	s calculated a er of amino a M, and XDoc y amino acids red. Errors fo ength histogra	according to ccids, and 0 ::Coh from 1 located C- r the observ m shown in	$\Delta L_E = N_A \cdot$ .365 nm <sup>2</sup> is the PDB structures terminal of the ved increments a Fig. 2b.

# Supplementary Notes

Supplementary Note 1: QM-FRC Model for Polymer Elasticity

The freely rotating chain model  $^3$  considers bonds of length b, connected by a fixed angle  $\gamma.$  The torsional angles are not restricted. The stretching behavior in the FRC picture is given by

$$\frac{x}{L} = \begin{cases} \frac{Fa}{3k_BT} & \text{for} \quad \frac{Fb}{k_BT} < \frac{b}{p} \\ 1 - \left(\frac{4Fp}{k_BT}\right)^{-\frac{1}{2}} & \text{for} \quad \frac{b}{p} < \frac{Fb}{k_BT} < \frac{p}{b} \\ 1 - \left(\frac{cFb}{k_BT}\right)^{-1} & \text{for} \quad \frac{p}{b} < \frac{Fb}{k_BT} \end{cases}$$
(1)

where  $a=b\frac{1+\cos\gamma}{(1-\cos\gamma)\cos\frac{\gamma}{2}}$  is the Kuhn length, and  $p=b\frac{\cos\frac{\gamma}{2}}{|\ln(\cos\gamma)|}$  is the effective persistence length in the FRC picture.

To account for backbone elasticity of the polypeptide chain at high force, quantum mechanical *ab-initio* calculations can be used to obtain the unloaded contour length at zero force. A polynomial approximation to these calculations can be used to obtain the unloaded contour length at zero force  $L_0$ :

$$F = \gamma_1 \left(\frac{L}{L_0} - 1\right) + \gamma_2 \left(\frac{L}{L_0} - 1\right)^2 \tag{2}$$

where the  $\gamma_1=27.4\,\mathrm{nN}\text{,}$  and  $\gamma_2=109.8\,\mathrm{nN}$  are the elastic coefficients reported for polypeptides^4.

Supplementary Note 2: Bell-Evans Model for Mechanically Induced Receptor Ligand Dissociation

The Bell-Evans model was used to estimate the distance to the transition state ( $\Delta x$ ) and the natural off-rate ( $k_{off}$ ) of individual rupture events:

$$\langle F \rangle = \frac{k_B T}{\Delta x} \ln \frac{\Delta x \cdot \dot{F}}{k_{off} k_B T}$$
(3)

where  $k_B$  is Boltzmann's constant, T is the temperature and  $\dot{F}$  is the loading rate at the point of rupture.

	nentary Methods
Materials	3
Silicon nits spring cons 2.4 cm in a dimethyl e PEG-Male TCEP Dist standard c tris(hydrox p.a), NaCl mM, pH 8. TRIS, 75 n were filtere USA) prior	ride cantilevers (Biolever mini, BL-AC40TS-C2, Olympus Corporation) with a nominal stant of 100 pN/nm (25 kHz resonance frequency in water) were used. Circular coverglasses, diameter, were obtained from Menzel Gläser (Braunschweig, Germany). 3-Aminopropyl ethoxysilane (APDMES) was purchased from ABCR GmbH (Karlsruhe, Germany). NHS- etimide (5 kDa) was purchased from Rapp Polymer (Tübingen, Germany). Immobilized ulfide Reducing Gel was obtained from Thermo Scientific (Pittsburgh, PA). The following chemicals were obtained from Carl Roth (Karlsruhe, Germany) and used as received: symethyl)aminomethane (TRIS, >99% p.a.), CaCl <sub>2</sub> (>99% p.a.), sodium borate (>99.8% (>99.5% p.a.), ethanol (>99% p.a.), and toluene (>99.5% p.a.). Borate buffer was 150 .5. The measurement buffer for force spectroscopy was Tris-buffered saline (TBS, 25 mM nM NaCl, pH 7.2) supplemented with CaCl <sub>2</sub> to a final concentration of 1 mM. All buffers ed through a sterile 0.2 $\mu$ m polyethersulfone membrane filter (Nalgene, Rochester, NY, r to use.
Protein S	Sequences
Sequences tags and re	of protein constructs used in this work are listed here. Domains as well as engineered esidues are color-coded.
Xyn-XMo	bdDoc
Xylanase T Linker or e X-module Dockerin t	F129C extra residues type III
M       S       H       H       I         A       V       E       P       Y         F       N       F       C       G         G       K       P       N       I         H       Q       S       H       I         H       Q       S       H       I         H       Q       S       H       I         V       V       F       K       G         V       V       F       H       I         G       V       V       I       I         G       V       V       I       I         G       V       V       I       I         G       V       V       I       I         G       V       V       I       I         G       V       P       V       I       I         G       V       V       I       I       I         G       V       V       I       I       I         J       D       A       T       A       I         F       S       A       F       I	H       H       H       H       H       H       H       H       H       H       K       N       A       D       S       Y       A       K       K       P       H       I       S       A       L       N       A       P       Q       L       D       Q       R       Y       K       N       E       F       T       I       G       A       L       N       K       N       F       T       I       G       R       Y       K       N       E       F       N       K       N       R       N       N       N       N       N       N       N       N       N       N       N       N       N       N       N       N       N



# 3.2 Force Propagation and the Influence of Anchoring Geometry

In an effort to better understand the molecular mechanisms, which allow the CohE:CttA-XDoc receptor-ligand pair to withstand remarkable forces (as we had found in a previous study (3.1)), in this article we investigate how force propagates through a protein complex under load. Using all-atom MD simulations, we use network-based correlation analysis to identify stiff paths through which force is transmitted. We find that for the natively occuring anchoring geometry, force is directed with strong normal components through the complex' binding interface. When anchoring the Cohesin non-natively at its N-terminus, the complexes ability to withstand force is compromised. I contributed to this manuscript by performing SMFS experiments together with Constantin Schöler. Additionaly, Wolfgang Ott and I cloned and expressed CohE with tags in both anchoring geometries, native and non-native. I also contributed to the analysis and interpretation of SMFS data.

# Mapping Mechanical Force Propagation through Biomolecular Complexes

Constantin Schoeler<sup>†</sup>, Rafael C. Bernardi<sup>†</sup>, Klara H. Malinowska, Ellis Durner, Wolfgang Ott, Edward A. Bayer, Klaus Schulten, Michael A. Nash, and Hermann E. Gaub

<sup>+</sup>these authors contributed equally to this publication

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configuration, we found that the complex dissociated along two competing pathways with very different mechanical characteristics.

Our new dynamic network analysis protocol reveals how different mechanical behaviors are attributable to differences in the direction of force transmission across the binding interface. Together, the experiments and simulations depict a simple physical mechanism for achieving high complex rupture forces: the complex directs force along pathways orthogonal to the pulling axis.

Single-Molecule Pulling Experiments and SMD. For SMFS experiments, XMod-Doc was produced as a fusion protein with an N-terminal Xylanase (Xyn) domain. Coh was produced as either an N- or C-terminal fusion domain with a carbohydrate binding module (CBM). These fusion domains were used for site specific immobilization to a glass surface and AFM cantilever to achieve the two loading configurations shown in Figure 1A and further served as marker domains with known unfolding length increments to validate single-molecule interactions and sort SMFS data traces.<sup>36</sup>

For the native pulling configuration found intext For the native pulling configuration found in vivo, Coh-CBM and XMod-Doc are loaded from their C- and N-termini, respectively (Figure 1A). A representative unbinding trace for the native pulling configuration is shown in Figure 1B. We measured the loading rate dependence of complex rupture using both experimental and SMD data sets (unbinding trace from SMD shown in Figure 3A) and plotted them on a combined dynamic force spectrum (Figure 1E). The linear Bell model produced fit parameters for the effective distance to the transition state  $\Delta x = 0.13$  nm, and the zero-force off rate  $k_{off} =$  $4.7 \times 10^{-4}$  s<sup>-1</sup>. Both experimental and simulation data are well described by a single Bell expression, despite the differences in loading rates between experiments and simulation. The observation suggests that the application of force does not siemificantly chance  $\Delta x$  for this particular configuration.

Significantly access that the appreciation of hore does not significantly change  $\Delta x$  for this particular configuration. To test the influence of pulling geometry on mechanical stability, we performed SMFS and SMD on the system where Coh was pulled from the opposite terminus (i.e., non-native Nterminus, cf. Figure 1A). Unlike the native pulling geometry, this geometry exhibited two clearly distinct unbinding pathways that are characterized by different force ranges (high or low) at which the complex dissociated. We refer to these pathways as non-native high force (HF) (Figure 1C) and non-native low force (LF) (Figure 1D).

AFM data traces classified as non-native HF showed similar characteristics as those in the native pulling configuration (cf. Figure 1B,C,F). The non-native LF traces, however, exhibited a markedly different unfolding behavior (Figure 1D). Xyn unfolding (highlighted in orange) was regularly observed, but CBM unfolding was only very rarely observed. The complex usually did not withstand forces high enough to unfold CBM when rupturing along the non-native LF path. Among nonnative LF curves, we regularly found an additional contour length increment of 17–19 nm consistent with unfolding of ~60 amino acids located at the N-terminus of Coh. This unfolding occurred immediately following Xyn unfolding (Figure 1D, red), or alternatively prior to Xyn unfolding, or with a substep (Supplementary Figure S1). Taken together, it appears that partial Coh unfolding from the N-terminus destabilizes the complex, causing lower rupture forces (Figure 1G).

The experimental rupture forces from the non-native HF class were indistinguishable from those arising in the native

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configuration. To confirm this, we performed additional measurements where both Coh configurations were alternately probed with the same Xyn-XMod-Doc functionalized cantilever (Supplementary Figure S2), eliminating inaccuracies introduced through multiple cantilever calibration. Most probable rupture forces at a pulling speed of 800 nm s<sup>-1</sup> of 606 and 597 pN for the native configuration and non-native HF class, respectively, were determined in the Bell Evans model (Figure IF, Supplementary eq S2), demonstrating that the native and non-native HF classes are experimentally indistinguishable. For the LF class, we analyzed the final complex rupture event

For the LF class, we analyzed the Innal complex rupture event and plotted the combined dynamic force spectrum (Figure 1G). Here, simulated and experimentally observed data were not well described by a single Bell expression. In such cases nonlinear models have been developed to obtain kinetic and energetic information from dynamic force spectra.<sup>57,38</sup> To fit the combined data, we used the nonlinear Dudko–Hummer– Szabo (DHS) model (Supplementary eq S3) and obtained values of  $\Delta x = 0.42$  nm and  $k_{off} = 0.005$  s<sup>-1</sup>. The DHS model further provides the free energy difference  $\Delta G$  between the bound state and the transition state as a fit parameter, which was found to be  $\Delta G = 129 k_B T$ . The model fit produced a distance to transition that was much longer than observed for the native configuration. Independent SMD simulations for the non-native pulling configuration were found to also lead to HF and LF unbinding scenarios (see below, Figure 4A,D, respectively).

The differential solvent contact area was calculated from SMD simulations to estimate the intermolecular contact area in the Doc:Coh complex. In the native configuration, the simulated Doc:Coh contact area increased by 14% and 9% for Coh and Doc, respectively (Figure 1H). For the non-native HF class, the contact area increased by 11% and 12% for Coh and Doc, respectively (Figure 1I). In the non-native LF class, the contact area increased by only 7% for Coh and decreased by 3% for Doc (Figure 1J). Evidently, an increased surface contact area for Doc in the native and non-native HF pathways correlated with high mechanostability of the system.

**Force Propagation Theory: A Simple Model.** To further understand the observed unbinding pathways, we sought to identify paths through the molecule along which the externally applied load propagates. From thermodynamic fluctuation theory,<sup>39,40</sup> it is known that the correlation of fluctuations of atoms *i* and *j* and the force **F**<sub>i</sub> on atom *i* are related through

$$\langle \Delta \mathbf{r}_i \Delta \mathbf{r}_j^T \rangle = k_{\rm B} T \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i}$$
(1)

where  $\Delta \mathbf{r}_i = \mathbf{r}_i(t) - \langle \mathbf{r}_i(t) \rangle$  and  $\mathbf{r}_i$  is the position of atom *i*. The derivative on the right-hand side of eq 1 states that neighboring atoms *i* and *j* will move with high correlation due to an external force  $\mathbf{F}_i$  acting on atom *i* if the coupling between them is strong. Hence, a given element of a correlation matrix  $M_{ij} = \langle \Delta \mathbf{r}_i \Delta \mathbf{r}_j \rangle$  will be large in the case of a strong interaction potential between *i* and *j*. When force is propagated through a molecule, soft degrees of freedom will be stretched out along the path of force propagation, while stiff degrees become more important for the dynamics of the system.

Consequently, paths with high correlation of motion describe the paths along which force propagates through the system. To illustrate this behavior for a toy system, we employed the NAMD<sup>41</sup> SMD<sup>42</sup> constant velocity protocol to a test pattern of identical spheres connected with harmonic springs of different

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# stiffness (Figure 2A). The position of one sphere was fixed during the simulation, while another sphere on the opposite



Figure 2. Network analysis test simulation. (A) Simulated pattern of atoms depicted by spheres. Connecting lines between atoms represent harmonic springs with different stiffnesses (red, k; blue, 5k; yellow, 7.5k; black, 10k). The green atom was fixed (anchor), while a second green atom was withdrawn at constant speed (arrow). Black and yellow atoms and their adjacent springs were introduced to maintain the general shape of the pattern. (B) Deformed sphere pattern at the end of the simulation. (C) Edges between nodes are weighted by the corresponding correlation matrix elements. (D) The path with highest correlation of motion is shown in red.

side of the structure was withdrawn at constant velocity. The strained structure at the end of the simulation is shown in Figure 2B. We assigned weights to the lines between spheres according to the Pearson correlation coefficient  $C_{ij}$  (Supplementary eq S4) between those network nodes (Figure 2C). The Pearson correlation factor  $(\langle \Delta \mathbf{r}_i^2(t) \rangle \langle \Delta \mathbf{r}_j^2(t) \rangle)^{-1/2}$  and was chosen to make our analysis mathematically more tractable. For a detailed discussion on this choice of correlation measure, see Supporting Information. In a harmonic potential approximation, the equipartition theorem can be applied to this normalization factor resulting in the following expression for  $C_{ij}$ :

$$C_{ij} = \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i} \sqrt{k_{i,\text{eff}} k_{j,\text{eff}}} \tag{2}$$

where  $k_{i,\text{eff}} = \left(\frac{1}{k_{x_i}} + \frac{1}{k_{y_i}} + \frac{1}{k_{z_i}}\right)^{-1}$  and  $k_{x_i}$  is the curvature of the

potential on atom *i* in the *x* direction. For a full derivation, see Supporting Information. Equation 2 illustrates how Pearson correlation is a suitable measure to identify the stiff paths in our simple model. We then used dynamical network analysis implemented in VMD<sup>49</sup> to find the path of highest correlation (Figure 2D). As expected from eq 1, we found this path to be the one connected by the stiff springs. Force Propagation through XMod-Doc:Coh Complex.

Force Propagation through XMod-Doc:Coh Complex. The simple pattern of spheres validated our general approach of using local correlations to identify load-bearing pathways through networks. We next employed dynamical network

#### Letter

analysis to understand force propagation through the XMod-Doc:Coh complex.

The dynamic networks for the native configuration (unloaded and loaded) are shown in Figure 3B,C, respectively. While the network shows multiple suboptimal paths in the unloaded scenario, the loaded case exhibits a well-defined main path along which force propagates through the system. Interestingly, in the loaded configuration, force propagates through both binding helices of Doc, which results in a force path with large normal components to the unbinding axis close to the binding interface as illustrated in Figure 3D. It had been shown for another ultrastable protein, namely, silk crystalline units, that curving force paths distribute tension through the entire system.<sup>31</sup> A strategy that assumes an indirect path would therefore allow the system to have more time to absorb the tension from the applied force. The result here supports the view that directing the force along a path with significant perpendicular components to the pulling axis leads to high mechanical stability. In a simple mechanical picture, a certain amount of mechanical work, namely  $dW = \mathbf{F} \cdot \mathbf{d}_S$  is required to separate the two binding interfaces by a distance  $\Delta z$  and break the interaction. In this simplified picture, ds points along the unbinding axis, whereas the force  $\mathbf{F}$  is locally largely perpendicular to this direction. Consequently, a larger force is required to break the interaction than in a scenario where the force path would point along the unbinding axis.

To validate this picture, we repeated the same analysis for the non-native HF and non-native LF pathways. The HF simulation (Figure 4A) exhibited only a small stretching of the flexible N-terminal region of Coh and complex dissociation at approximately 800 pN and a pulling distance around 10 nm. However, the LF case shown in Figure 4D exhibited a stepwise N-terminal Coh unfolding, dissociating at a force of about 480 pN at a pulling distance of about 25 nm. This behavior confirmed our assignment of the experimentally observed 17– 19 nm contour length increment to Coh unfolding up to residue 62 in PDB 41U3.

While the experimental data did not show a detectable difference between the native configuration and the non-native HF class, the propagation of force takes place along a different pathway (Figure 4B). For N-terminal Coh pulling, helix 3 of Doc is not involved in the propagation of force as it is for the native geometry. In the native configuration, force propagates through the center of Coh, while for non-native HF the path is shifted toward the side of the molecule. Despite these differences, there is a common feature between the native and non-native HF pathways. At the binding interface, the pathway again shows pronounced components perpendicular to the unbinding axis (cf. Figure 4C), suggesting that this feature is indeed responsible for the exceptional mechanical strength observed for these two unbinding pathways.

sobserved for these two unbinding pathways. Figure 4E shows the force propagation pathway for the nonnative LF class prior to rupture. Due to the unfolding of the Nterminal Coh segment, the propagation of force is shifted even further away from the central portion of Coh than for the nonnative HF class. Interestingly, force is propagated through the small helical segment of Coh (ALA167-GLN179), a portion of the molecule that is not involved in force propagation for any of the other analyzed trajectories. Unlike in the aforementioned scenarios, there is no pronounced tendency for perpendicular force components at the binding interface for the non-native LF class. In fact, the force is propagated along a path largely parallel to the pulling axis (cf. Figure 4F). In cases where force

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 $\mathbf{3}$ 

which were randomly arranged in the solvent. Total system size was approximately 580k atoms. The MD simulations in the present study were performed employing the molecular dynamics package NAMD<sup>10;11</sup>. The CHARMM36 force field<sup>12;13</sup> along with the TIP3 water model<sup>14</sup> was used to describe all systems. The simulations were carried out assuming periodic boundary conditions in the NpT ensemble with temperature maintained at 300 K using Langevin dynamics for pressure, kept at 1 bar, and temperature coupling. A distance cut-off of 11.0 Å was applied to short-range, non-bonded interactions, whereas long-range electrostatic interactions were treated using the particle-mesh Ewald (PME)<sup>15</sup> method. The equations of motion were integrated using the r-RESPA multiple time step scheme<sup>11</sup> to update the van der Waals interactions every two steps and electrostatic interactions every four steps. The time step of integration was chosen to be 2 fs for all simulations performed. The first two nanoseconds of the simulations served to equilibrate systems before the production runs, which varied from 200 ns to 1.3  $\mu$ s in the different simulations. To characterize the coupling between dockerin and cohesin, we performed SMD simulations<sup>16</sup> of constant velocity stretching (SMD-CV protocol) with pulling speed of 0.25 Å/ns. In all simulations, SMD was employed by restraining the position of one end of the XMod-Doc domain harmonically, and moving a second restraint point, at the end of the Coh domain, with constant velocity in the desired direction. The procedure is equivalent to attaching one end of a harmonic spring to the end of a domain and pulling on the other end of the spring. The force applied to the harmonic pulling spring is then monitored during the time of the molecular dynamics simulation. All analyses of MD trajectories were carried out employing  $\rm VMD^{17}$  and its plugins. Surface contact areas of interacting residues were calculated employing Volarea<sup>18</sup> implemented in VMD. The area is calculated using a probe radius defined as an *in* silico rolling sphere that is scanned around the area of the dockerin exposed to the cohesin and also the cohesin area exposed to the docker in. The Network View plugin  $^{19}$  on VMD  $^{17}$  was employed to perform dynamical network analysis. A network was defined as a set of nodes, all  $\alpha$ -carbons, with connecting edges. Edges connect pairs of nodes if corresponding monomers are in contact, and 2 nonconsecutive monomers are said to be in contact if they fulfill a proximity criterion, namely any heavy atoms (nonhydrogen) from the 2 monomers are within 4.5 Å of each other for at least 75% of the frames analyzed. As suggested by Sethi et al.<sup>20</sup>, nearest neighbors in sequence are not considered to be in contact as they lead to a number of trivial suboptimal paths. The dynamical networks were constructed from 20 ns windows of the total trajectories sampled every 400 ps. The probability of information transfer across an edge is set as  $w_{ij} = -log(|C_{ij}|)$ , where  $C_{ij}$  is the correlation matrix calculated with Carma<sup>21</sup>. Using the Floyd-Warshall algorithm, the suboptimal paths were then calculated. The tolerance value used for any path to be included in the suboptimal path was -log(0.5) = 0.69. To calculate the relevance of off-diagonal terms in the correlation matrix we employed Carma to calculate a correlation matrix where x, y, z components of each atom were considered independently.

2 Protein Sequences

Sequences of protein constructs used in this work are listed here. Domains as well as engineered tags and residues are color-coded.

2.1 HIS-Xyn(T128C)-XDoc

X-module Dockerin type III



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### 3 Supplementary Discussion

The Pearson correlation matrices of the Xmod-Doc:Coh complex before and after applying force in the native pulling configuration are presented in Supplementary Figure S3 and S4, respectively. For the unloaded complex, movements within Doc domain are seen to be highly correlated, while XMod is seen to be divided into two anti-correlated sub-domains, one comprising the  $\beta$ -sheet fragment close to the N-terminus (residues 5-15 and 45-66) and the other constituting the rest of the domain. Intra-domain correlations of Coh exhibit more a complex pattern to which both secondary (anti-parallel  $\beta$ -strands and  $\beta$ -sheet at the binding interface) and tertiary structure (vicinity of C- and N-termini) contribute. Some of the inter-domain correlations in the complex originate from spatial vicinity and direct interactions, specifically at the Doc:Coh binding interface and at XMod contacts with Doc inserts. However, coupling between distant parts of the complex is also present. For example, fluctuations of the non-binding part of Coh are correlated with the N-terminal part of XMod and strongly anti-correlated with Doc domain.

# 4 Supplementary Notes

# 4.1 Constant Barrier Distance Model

The constant barrier distance model<sup>16</sup>, also referred to as the Bell-Evans model<sup>22</sup>, is commonly used to estimate the distance to the transition state  $\Delta x$  and the natural off-rate  $k_0$  of mechanically induced receptor ligand dissociation from single-molecule force spectroscopy experiments. It predicts that the most probable rupture force  $\langle F \rangle$  is linearly dependent on the logarithm of the force loading rate<sup>16</sup>:

$$\langle F(r) \rangle = \frac{k_B T}{\Delta x} \ln \frac{\Delta x \cdot r}{k_0 k_B T} \tag{S1}$$

where  $k_B$  is Boltzmann's constant, T is the temperature and r is the loading rate at the point of rupture.

The probability density distribution of rupture forces at given loading rate r in this model is given as  $^{16}$ :

$$p(F) = \frac{k_0}{r} \exp\left[\frac{\Delta x}{k_B T} F - \frac{k_0 \cdot k_B T}{\Delta x \cdot r} \left(e^{\frac{\Delta x}{k_B T}F} - 1\right)\right]$$
(S2)

# 4.2 Dudko-Hummer-Szabo Model

The Dudko-Hummer-Szabo (DHS)  $^{23;24}$  model describes a non-linear dependence for the most probable rupture force on loading rate:

$$\langle F(r) \rangle = \frac{\Delta G}{\nu \Delta x} \left\{ 1 - \left[ \frac{k_B T}{\Delta G} \ln \left( \frac{k_B T k_0}{\Delta x r} e^{\frac{\Delta G}{k_B T} + \gamma} \right) \right]^{\nu} \right\}$$
(S3)

where  $\Delta G$  is the free energy of activation and  $\gamma = 0.577$  is the Euler-Mascheroni constant. The model parameter  $\nu$  defines the single-well free-energy surface model used ( $\nu = \frac{2}{3}$  for linear-cubic and  $\frac{1}{2}$  for cusp free-energy. For  $\nu = 1$  and  $\Delta G \rightarrow \infty$  independent of  $\nu$  the Eqs. (S1) and (S2) are recovered.

4.3 Pearson Correlation and covariance matrix

# 4.3.1 Validation

An  $N \times N$  matrix of Pearson correlation coefficients  $C_{ij}$  (Supporting Eq. S4) was calculated from each atom's x, y, z position throughout the simulation trajectory, which inherently ignores off-diagonal elements of the atomic  $3 \times 3$  submatrices  $D_{ij}^{mn}$  from the full normalized  $3N \times 3N$  covariance matrix (*i.e.*, correlations along orthogonal axes are neglected, see Supporting Eqs. (S5) and (S6)) and Supporting Fig S8.

Although this quasi-harmonic approximation is commonly employed in correlation analy $sis^{19;25-29}$ , it is not a priori justified for complicated biomolecular interactions<sup>30</sup>. To validate the use of Pearson correlations, we therefore first analyzed independently the contributions from diagonal and off-diagonal elements of each 3 x 3 covariance submatrix for each pair of  $\alpha$ -carbons within the structure (Fig. S9A and B). Both with and without applied force, the off-diagonal elements roughly follow Gaussian distributions centered around a correlation value of 0. Interestingly, as force was applied, the standard deviation of the distribution of off-diagonal correlation values decreased from  $\sigma_{unloaded} = 0.45$  to  $\sigma_{loaded} = 0.29$ . This indicated a lesser influence of off-diagonal elements on the highly (anti-)correlated motion within the system under force (see Supporting Discussion 3). The diagonal elements of the sub-matrices that are used for calculating the Pearson correlation values showed a dramatically different behavior. Both in the unloaded and loaded state, the resulting distributions were strongly shifted towards highly correlated motion, and the shape of the distribution remained mostly unchanged after application of force. Since our analysis relies on the identification of paths of highest correlation through proximate residues, the quasi-harmonic approximation implied by the use of Pearson correlation is justified, especially for suboptimal pathway analysis. The resulting distributions of on- and off-diagonal matrix elements of each covariance submatrix for the loaded configuration HF class (Fig. S10A) and LF class (Fig. S10B) exhibited the same characteristics as previously described for the native configuration, with off-diagonal elements showing symmetric correlations around zero and diagonal elements showing highly correlated motions.

4.3.2 Supplementary Equations

The Pearson correlation coefficient  $C_{ij}$  used in our dynamical network analysis protocol is given by:

$$C_{ij} = \frac{\langle \Delta \mathbf{r}_i(t) \cdot \Delta \mathbf{r}_j(t) \rangle}{\left( \left\langle \Delta \mathbf{r}_i(t)^2 \right\rangle \left\langle \Delta \mathbf{r}_j(t)^2 \right\rangle \right)^{\frac{1}{2}}}$$
(S4)

where  $\Delta \mathbf{r}_{i}(t) = \mathbf{r}_{i}(t) - \langle \mathbf{r}_{i}(t) \rangle$ . The full  $3N \times 3N$  covariance matrix  $M_{ij}$  for atoms *i* and *j* consists of  $3 \times 3$  submatrices of the form:

$$\left\langle \Delta \mathbf{r}_{i}\left(t\right) \Delta \mathbf{r}_{j}\left(t\right)^{T} \right\rangle = M_{ij} = \begin{pmatrix} M_{ij}^{xx} & M_{ij}^{xz} & M_{ij}^{xz} \\ M_{ij}^{yx} & M_{ij}^{yy} & M_{ij}^{yz} \\ M_{ij}^{zx} & M_{ij}^{zy} & M_{ij}^{zz} \end{pmatrix}$$
(S5)

The full normalized correlation matrix is calculated from  $M_{ij}$ :

$$D_{ij}^{mn} = \frac{M_{ij}^{mn}}{\sqrt{M_{ij}^{mm}M_{ij}^{nn}}}$$
(S6)

Consequently, the Pearson correlation coefficient is calculated as the trace of the normalized  $3 \times 3$  submatrices  $(C_{ij} = \operatorname{Tr} D_{ij}).$ 

 $4.3.3\,$  Derivation of Main Text Equation 2

Eq. 1 from the main text reads:

$$\left\langle \Delta \mathbf{r}_i \Delta \mathbf{r}_j^T \right\rangle = k_B T \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i} \tag{S7}$$

Combining Eqs. (S7) and (S4) yields:

$$C_{ij} = k_B T \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i} \cdot \left( \left\langle \Delta \mathbf{r}_i^2\left(t\right) \right\rangle \left\langle \Delta \mathbf{r}_j^2\left(t\right) \right\rangle \right)^{-\frac{1}{2}}$$
(S8)

For an arbitrary potential  $U_{i}\left(\mathbf{r}\right)$  of atom i, a Taylor expansion around the potential minimum (set to be at 0) yields:

$$U_i(\mathbf{r}) = 0 + \underbrace{\mathbf{r}_i^T \nabla U(0)}_{=0} + \frac{1}{2} \mathbf{r}_i^T \underline{H}(0) \mathbf{r}_i + \dots$$
(S9)

where  $\underline{H}(0)$  is the Hessian matrix evaluated at the potential minimum. Assuming Schwarz' theorem holds for  $U_i(\mathbf{r}), \underline{H}(0)$  is a symmetric matrix and therefore has real eigenvalues and orthonormal eigenvectors. Hence, a change to the eigenbasis of  $\underline{H}\left(0\right)$  is a rotation of the coordinate system. In this new basis the Hessian is diagonal:

$$\underline{H}(0) \to \underline{H}'(0) = \begin{pmatrix} k_{x'} & 0 & 0\\ 0 & k_{y'} & 0\\ 0 & 0 & k_{z'} \end{pmatrix}$$
(S10)

This yields a simple expression for the second order term in Eq. (S9):

$$U_{i}(\mathbf{r}') = \frac{1}{2} \mathbf{r}'^{T} \underline{H}'(0) \mathbf{r}' = \frac{1}{2} \left( k_{x'} x'^{2} + k_{y'} y'^{2} + k_{z'} z'^{2} \right)$$
(S11)

Now we inspect the normalization of  $C_{ij}$ :

$$\langle \Delta \mathbf{r}_{i}^{2}(t) \rangle = \langle \mathbf{r}_{i}^{2}(t) - 2\mathbf{r}_{i}(t) \langle \mathbf{r}_{i}(t) \rangle + \langle \mathbf{r}_{i}(t) \rangle^{2} \rangle \tag{S12}$$

In the harmonic approximation of the potential of atom i,  $\langle \mathbf{r}_i(t) \rangle = 0$ , and therefore  $\langle \Delta \mathbf{r}_i(t)^2 \rangle = \langle \mathbf{r}_i^2(t) \rangle$ . In the basis of  $\underline{H}'(0)$  this becomes:

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$$\langle \mathbf{r}'_{i}^{2}(t) \rangle = \langle x'_{i}(t)^{2} + y'_{i}(t)^{2} + z'_{i}(t)^{2} \rangle = \langle x'_{i}(t)^{2} \rangle + \langle y'_{i}(t)^{2} \rangle + \langle z'_{i}(t)^{2} \rangle$$
(S13)

Applying the equipartition theorem to this result yields:

$$\langle x_i'(t)^2 \rangle = \frac{k_B T}{k_{x_i}'} \tag{S14}$$

And therefore:

$$\langle \Delta \mathbf{r}'_{i}(t)^{2} \rangle = k_{B}T \left( \frac{1}{k'_{x_{i}}} + \frac{1}{k'_{y_{i}}} + \frac{1}{k'_{z_{i}}} \right) = \frac{k_{B}T}{k'_{i,eff}}$$
(S15)

Plugging this result into Eq. (S8), one finds:

$$C_{ij} = k_B T \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i} \cdot \left(\frac{k_B T}{k'_{i,eff}}\right)^{-\frac{1}{2}} \left(\langle \Delta \mathbf{r}_j \left(t\right)^2 \rangle\right)^{-\frac{1}{2}}$$
(S16)

Repeating the above steps for atom j yields the final result:

$$C_{ij} = k_B T \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i} \cdot \left(\frac{k_B T}{k'_{i,eff}}\right)^{-\frac{1}{2}} \left(\frac{k_B T}{k'_{j,eff}}\right)^{-\frac{1}{2}}$$
(S17)  
$$= \frac{\partial \mathbf{r}_j}{\partial \mathbf{F}_i} \cdot \sqrt{k'_{i,eff} \cdot k'_{j,eff}}$$
(S18)




















# CHAPTER 3. RESEARCH ARTICLES



[28] Rebecca W Alexander, John Eargle, and Zaida Luthey-Schulten. Experimental and computational determination of tRNA dynamics. FEBS Lett., 584(2):376–386, January 2010.[29] Rafael C Bernardi, Isaac Cann, and Klaus Schulten. Molecular dynamics study of enhanced Man5B enzymatic activity. Biotechnol. Biofuels, 7(83):1–8, 2014. [30] Oliver F Lange and Helmut J Grubmüller. Generalized correlation for biomolecular dynamics. Proteins, 62(4):1053–1061, March 2006. 21

# 3.3 Comparing two Cellulosomal Anchoring Interactions

This publication investigates the receptor-ligand pair ScaB-XDoc:CohE (*R.f.*), one of the strongest receptor:ligand interactions found to date. It is a homologue of the interaction between Ctta-XDoc and CohE, which are described in publications 3.1 and 3.2. Despite the very similar binding interfaces of CttA-XDoc and ScaB-XDoc and comparable binding constants, we find the latter to be substantially stronger. We attribute this increase in mechanical strength to the larger X-module domain of ScaB, which unfolds at higher forces and more seldomly than its CttA counterpart. Using all-atom MD simulations, we can corroborate the previously assumed role of the Xmodule as a mechanical stabilizer, and propose that it is essential for distributing force to the Dockerin's binding interface. I contributed to the design of the experimental part of the study, and performed SMFS experiments and analysis. Interpretation of the results as well as manuscript preparation was shared between Michael A. Nash, Rafael C. Bernardi, and myself.

# Mechanisms of Nanonewton Mechanostability in a Protein Complex Revealed by Molecular Dynamics Simulations and Single-Molecule Force Spectroscopy

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 $K_{\rm D}$  of CttA-XDoc:CohE obtained by isothermal titration calorimetry (see Supporting Information Figure S6) was 18 nM while the  $K_{\rm D}$  of ScaB-XDoc:CohE was 149 nM, both of which are comparable to other Coh:Doc interactions. The discrepancy between the equilibrium analysis, where Ctta-XDoc:CohE exhibits higher affinity, and the mechanical analysis, where ScaB-XDoc:CohE exhibits higher rupture forces, highlights how mechanical dissociation can direct unbinding reactions along pathways that are distinct from those dominant in the absence of force.

High-Force Mechanism. An increase in contact area upon application of force was observed in previous SMD simulations of the CttA-XDoc:CohE complex.<sup>12</sup> Taking advantage of the recently developed PyContact,<sup>43</sup> we further analyzed our simulation trajectories by quantifying the contact surface area between CohE and ScaB-XDoc in both loaded and unloaded scenarios. For the loaded scenario, we analyzed all trajectories immediately prior to complex rupture over a simulation window of 4.0 ns. For the unloaded scenario, we analyzed equilibrium MD simulations, also over a 4.0 ns window. The total contact area was found to increase upon loading due to a rearrangement of interacting residues, as shown in Figure 5A. We further analyzed the contact surface area per residue (Figure 5B, C), and plotted the contact surface area for main interacting residues as shown in Figure 5D and E.

ScaB-XDoc (Figure 5D) and CohE (Figure 5E) both exhibit a binding interface comprising a hydrophobic core (white) surrounded by polar (green) and charged residues (blue, positive; red, negative). This residue pattern suggests that the hydrophilic side chains prevent access of solvent to the interior core. Such hot spots of binding are a common motif in protein-protein binding interfaces, referred to as O-rings.<sup>44</sup> Contact area between the two proteins increased upon mechanical loading based on rearrangement of side chains and closing off of the hydrophobic core to solvent (Figure SF, G).

**Exemplary Force Propagation Pathways.** To investigate how strain propagates through the protein complex, we employed cross-correlation based network analysis to study one-step rupture events and rupture events following Xmod unfolding. From thermodynamic fluctuation theory, one can infer that paths with high correlation of motion describe the paths along which force propagates through a molecular system.<sup>15,45</sup> The relationship between fluctuations of atoms *i* and *j* and the force  $\vec{F}_i$  on atom *i* can be described by

$$\langle \Delta \vec{r}_i \Delta \vec{r}_j^T \rangle = k_{\rm B} T \frac{\partial r_j}{\partial \vec{l}_i^2} \tag{1}$$

where  $\Delta \vec{r}_i = \vec{r}_i(t) - \langle \vec{r}_i(t) \rangle$  and  $\vec{r}_i$  is the position of atom *i*. If an external force  $\vec{F}_i$  acts on atom *i* and the potential between atoms *i* and *j* is steep, the two atoms will move with high positional correlation. By calculating a given element of the correlation matrix  $M_{ij} = \langle \Delta \vec{r}_i \Delta \vec{r}_j^T \rangle$ , we can quantify the strength of an interaction potential between *i* and *j*. The deformation response of the molecular complex under applied force will be dominated by the stiff interaction potentials, while the soft potentials become insignificant.<sup>15</sup> The propagation of force through the molecules and across the binding interface can then be visualized along the pathways with high correlations of motion. The high-correlation path with the smallest length, namely, the smallest number of amino acid residues, is considered to be the optimal path. When multiple paths of

similar length (number of nodes) are found to have high correlation, then these are considered suboptimal paths. It is important to notice that not all residues along these paths need be considered to have the same importance for force propagation. Instead, only residues or interactions that occur in the highest number of suboptimal pathways need to be conserved to guarantee an effective force propagation pathway in the complex.<sup>15,29</sup>

Although the quasi-harmonic approximation of Pearson correlation is commonly employed in correlation analysis of molecular systems,<sup>46–51</sup> it is not a priori justified for complicated biomolecular interactions.<sup>52</sup> However, since our analysis relies on the identification of paths of highest correlation through proximate residues, the quasi-harmonic approximation implied using Pearson correlation is justified, particularly for suboptimal pathway analysis.<sup>15</sup>

In Figure 6, the force propagation pathway analysis. In Figure 6, the force propagation pathways through both ScaB-XDoc:CohE (Figure 6A, C) and CttA-XDoc:CohE (Figure 6B, D) complexes are depicted. Figure 6A and B shows correlation networks obtained from one-step rupture event trajectories for ScaB-XDoc:CohE and CttA-XDoc:CohE, respectively. Figure 6C and D shows correlation networks obtained from simulations exhibiting Xmod unfolding prior to rupture. The depicted paths through the system are those along which stress primarily propagates under load, obtained just prior to rupture over a simulation window of 4.0 ns.

Just prior to rupture over a simulation window of 4.0 ns. For the one-step rupture rajectories, it was previously shown for CttA-XDoc:CohE that force propagated through both binding helices of the Doc and along pathways with large normal components with respect to the pulling axis.<sup>12</sup> The thickness of the lines in Figure 6 indicates how important these pathways are. That is calculated by considering the multiple possible paths of similar length with high-correlation. Important nodes (amino acid residues) are those that are present in a large number of paths. Likewise, important edges (a line connecting two nodes) are those edges that are part of a large number of paths. Therefore, the most important pathways are those that are shown with thicker lines. Although clear differences between the four force pathways shown in Figure 6 are evident, one can easily observe that these important pathways all have "sharp turns" near the interface. These force-propagation routes nonparallel to the pulling axis for ScaB-XDoc:CohE are advantageous for achieving high dissociation forces.

It was previously shown for another ultrastable protein, namely, silk crystalline units, that curving and branching force paths distributed tension through the entire system.<sup>53</sup> A strategy that assumes an indirect path would allow the system to more evenly utilize the interface between binding partners. Such behavior can be simplified in a mechanical picture, where a certain amount of mechanical work  $dW = \vec{F} \cdot \vec{d}$  is required to separate two binding interfaces by a distance  $\Delta z$  and break the interaction. If force  $\vec{F}$  is locally perpendicular to the direction of the unbinding axis  $d\vec{s}$ , a larger force is required to break the interaction than in a scenario where the force path points along the unbinding axis.

Analysis of the correlation maps indicated that the calciumbinding loops were key in distributing forces through the Coh:Doc binding interface. In all scenarios, force propagated through the calcium loops. Even after unfolding of Xmod, force-propagation through the folded Doc remained largely unchanged. As shown in Figure 4C, complex rupture forces were greatly reduced when the Xmod was unfolded, but

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#### Article

remained on the order of 200 pN at the given loading rates. Overall, our correlation analysis supports the view that directing force along a path with significant perpendicular component to the pulling axis leads to high mechanical stability.

## CONCLUSION

The detailed molecular mechanisms of cellulosomes are of broad interest, particularly given their applications in biotechnology  $^{54}$  and their recent discovery as part of the human gut microbiome.  $^{55,56}$  The ScaB-XDoc:CohE receptor ligand complex is essential for bacterial degradation of cellulose, serving as an adhesion domain and anchoring its organism to a primary carbon source.

The multistep computational pipeline we described to understand the ScaB-XDoc:CohE complex involved homology modeling, wide-sampling of steered molecular dynamics simulations, contact area analysis, and finally dynamic network analysis. We developed this pipeline in order to structurally characterize the ScaB-XDoc:CohE complex, computationally predict its high resistance to applied forces, and understand the molecular mechanisms involved in force resistance. We furthermore were able to confirm the predicted high forces experimentally using parallel high volume acquisition of single-molecule AFM force traces.

The additional strength observed for the ScaB-XDoc:CohE complex is in agreement with its function. The actual load on this complex, with its attached Doc-bearing enzymes and other subunits, would seen to justify the improved mechanical stability of the ScaB-XDoc:CohE complex over the CttA-XDoc:CohE. The primary differences in behavior between ScaB-XDoc and the previously characterized CttA-XDoc arise from the Xmod region, which is significantly longer in ScaB. Unfolding of the Xmod domain prior to complex rupture caused a strong destabilization of the complex, and eliminated the improved mechanical stability of ScaB-XDoc:CohE over CttA-XDoc:CohE complexes. As the ScaB's Xmod unfolded far less frequently and at higher forces than its CttA analogue, its larger structure rendered it more robust, improving the stabilizing effect it had on the Doc:CohE complex. Taken together, these results depict the Xmod's function as a structural support of the Doc's binding helices under high loads, acting in part as a mechanical shield to protect the adjacent Doc from unfolding. Our wide-sampling simulation approach with many simulation replicas was critical for interpreting the mechanical behavior of these receptor-ligand complexes, particularly in describing the mechanisms mechanical stability. In the future, successful merger of SMD simulations with single-molecule mechanical experiments will benefit from the analysis pipeline presented here in order to gain insight into the fascinating interplay between equilibrium protein binding affinity and adhesion mechanics.

#### METHODS

**Structural Model.** The structure of the CttA-XDoc:CohE complex has been solved by means of X-ray crystallography at 1.97 Å resolution and is available in the Protein Data Bank (PDB: 4IU3).<sup>20</sup> A resolution and is available in the Protein Data Bank (PDB: 4103). The second system, namely, ScaB-XDoc, had no structure available and was solved by a homology modeling strategy.<sup>30</sup> The construction of Coh models was performed using MODELLER 9.17<sup>31</sup> software that employs spatial restriction techniques based on the 3D-template structure. The best model was selected by analyzing the stereo-chemical quality check using PROCHECK<sup>55</sup> and overall quality by

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ERRAT server.<sup>58</sup> The structures of two sequence gaps in the loop regions marked by arrows in Figure 1C were resolved by a loop optimization protocol employing ROSETTA.<sup>33</sup> Using NAMD<sup>34,55</sup> through its QwikMD interface,<sup>56</sup> structures were subjected to 100 ns through its QwikwiD interface, surctures were subjected to for its of equilibrium MD to ensure conformational stability. All structures shown are from postequilibration MD simulations. Molecular Dynamics Simulations. MD simulations were

performed employing the NAMD molecular dynamics package.<sup>34</sup> The CHARMM36 force field<sup>59</sup> along with the TIP3 water model<sup>60</sup> was used to describe all systems. Simulations were carried out assuming periodic boundary conditions in the NpT ensemble with temperature maintained at 300 K using Langevin dynamics for pressure, kept at 1 bar, and temperature coupling. A distance cutoff of

pressure, kept at 1 bar, and temperature coupling. A distance cutoff of 14.0 Å was applied to short-range, nonbonded interactions, whereas long-range electrostatic interactions were treated using the particlemesh Ewald (PME) method. The equations of motion were integrated using the r-RESPA multiple time step scheme<sup>34</sup> to update the van der Waals interactions every two steps and electrostatic interactions every two steps and electrostatic interactions every four steps. The time step of integration was chosen to be 2 fs for all simulations performed. In a wide-sampling strategy,<sup>61</sup> 75 steered molecular dynamics (SMD) runs were carried out for a total of 1.5  $\mu$ s for each system. To characterize the coupling between dockerin and cohesin, SMD simulations<sup>38</sup> with constant velocity stretching (SMD-CV protocol) at 5.0 Å/ns pulling speed were carried out. In all simulations, SMD was employed by restraining the position the N-terminal of ScaB-XDoc or CttA-XDoc, while pulling on the C-terminus of CohE. The procedure is equivalent to attaching one end of a harmonic spring to the end of the domain with a spring, analogous one domain and pulling on the other domain with a spring, analogous to the experiment. The force applied to the harmonic spring was then monitored during the time of the molecular dynamics simulation.

All analyses of MD trajectories were carried out employing VMD<sup>32</sup> and its plugins. Surface contact areas of interacting residues were calculated employing PyContact.<sup>43</sup> The Network View plugin<sup>29</sup> on VMD was employed to perform dynamical network analysis. A network was defined as a set of nodes comprising all *a*-carbons connected together by edges. Edges connect pairs of nodes if corresponding monomers are in contact, and two nonconsecutive monomers are said to be in contact if they fulfill a proximity criterion, monomers are said to be in contact if they fulfill a proximity criteron, namely, any heavy atoms (non-hydrogen) from the two monomers are within 4.5 Å of each other for at least 75% of the frames analyzed. As suggested by Sethi et al.,<sup>29</sup> nearest neighbors in sequence are not considered to be in contact as they lead to a number of trivial suboptimal paths. Conceptually, suboptimal paths can be understood as allosteric signaling pathways or force propagation pathways. Suboptimal paths are defined as paths that are slightly longer than the cotimal path with a eiven suboptimal path sciling a node not buopmina paths are centred as paths that are signed integrit time the optimal path, with a given suboptimal path visiting a node not more than once. These multiple communication paths are nearly equal in length, and not all residues along these paths need be

considered important for allostery. Instead, only residues or interactions that occur in the highest number of suboptimal pathways need to be conserved to guarantee an effective pathway for allosteric communication. The thickness of the effective pathway for allosteric communication. The thickness of the edges connecting the nodes reveals the least and most used paths. Allostery can be understood in terms of pathways of residues that efficiently transmit energy, here in the form of mechanical stress,  $^{15,26,63}$  between different binding sites. The dynamical networks were constructed from 4 ns windows of the total trajectories. Using the Floyd–Warshall algorithm, the suboptimal paths were then calculated. The totarcac value used for any path to be included in the suboptimal path was  $-\ln 0.5 = 0.69$ . To calculate the relevance of off-diagonal terms in the correlation matrix, we employed Carma to calculate a correlation matrix where x, y, and z components of each atom were considered independently.

calculate a correlation matrix where x, y, and z components of each atom were considered independently. **Protein Preparation for Experiments.** All proteins were expressed from pET28a vectors using standard induction and expression protocols in NiCo21(DE3)RIPL cells, which were cultivated in ZYM-5052 autoinduction media<sup>64</sup> supplemented with kanamycin and chloramphenicol. After pelleting, cells were lysed by sonication and then centrifuged at 4 °C, 39 000 rcf for 60 min. The

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supernatant was filtered through a 0.22 PES membrane (Carl Roth + Co. KG, Karlsruhe, Germany) and applied to Ni-NTA columns (HisTrap FF, GE Healthcare Europe GmbH, Freiburg, Germany). After washing with 6 column volumes of a buffer containing 25 mM TRIS, pH 8.4, 300 mM NaCl, 20 mM imidazole, and 0.5% (v/v) Triton X-100, the bound fraction was eluted with a buffer containing 25 mM TRIS, pH 8.4, 300 mM NaCl, and 300 mM imidazole. All protein solutions were concentrated using Amicon centrifugal filter units (10K MWCO, Merck KGaA, Darmstadt, Germany), followed by buffer exchange against Ca-TBS (25 mM TRIS, pH 7.2, 75 mM NaCl, and 1 mM CaCl.) buffer using ZebaSpin columns (Zeba spin desalting columns 7K, Thermo Fisher Scientific Inc.). Proteins were stored at -80 °C, with glycerol added to 10% (v/v). **Surface Functionalization**. Glass surfaces and silicon nitride cantilevers (BioLever mini BL-AC40TS-C2, Olympus, Tokio, Japan) were silanized with (3-aminopropyl)-dimethyl-ethoxysilane (APDMES, ABCR GmbH, Karlsruhe, Germany). Utilizing silicon masks (CultureWell Reusable Gaskets, Grace Bio-Laboratories, Bend, OB), two snitally senarated sonts on the silanized elas surfaces were

Surface Functionalization. Glass surfaces and silicon nitride cantilevers (BioLever mini BL-AC40TS-C2, Olympus, Tokio, Japan) were silanized with (3-aminopropyl)-dimethyl-ethoxysilane (APDMES, ABCR GmbH, Karlsruhe, Germany). Utilizing silicon masks (CultureWell Reusable Gaskets, Grace Bio-Laboratories, Bend, OR), two spatially separated spots on the silanized glass surfaces were PEGylated with *a*-maleimindo-hexanoic-*ao*-NHS polyethylene glycol (NHS-PEG5000-Mal, Rapp Polymere, Tübingen, Germany) dissolved into 25 mM in 4-(2-hydroxyethyl)-1-piperazineethanesulfonic acid buffer (HEPES), 50 mM, pH 7.5. Cantilevers were PEGylated using the same conditions. Next, the PEGylated surfaces and cantilevers were coupled to Coenzyme A (CoA, 1 mM) in sodium phosphate buffer, pH 7.2. Finally ybbR-CBM(CG3)-SCB-XDoc or bbR-CBM(CG3)-CtLA-XDoc was covalently immobilized onto the two spatially separated spots on the glass slide via their ybbR-tags in an Sip-catalyzed ligation at room temperature for 30 min. Each Dockerin was diluted to 0.5  $\mu$ M in Ca-TBS supplemented with 20 mM MgC1, while the Sip enzyme was added to 1  $\mu$ M. CohE-CBM(CG3)-ybbR was coupled to cantilevers at a concentration of 20  $\mu$ M under the same conditions.

Single-Molecule Force Spectroscopy. Measurements were performed in Ca-TBS buffer using custom built AFM instruments (driven vertically bp PI-731 piezo actuators and laterally bp a 25 x 25 mm piezomotor (U-751) in combination with a 100 × 100 nm (P-734) stage, Physik Instrumente, Germany) in conjunction with MFP-3D AFM controllers (Asylum Research, Santa Barbara, CA). Upon approaching the sample surface with the cantilever tip, the complex between CohE and either CttA-XDoc or ScaB-XDoc was formed and the cantilever was retracted from the surface at constant velocities of 100, 200, 400, 800, 1600, 3200, and 6400 nm/s. After each forceextension curve was acquired, the sample was moved laterally by 100 nm in order to probe a different molecule. Every several hundred measurements, the glass slide was moved laterally between protein spots, such that alternatingly CohE-ScaB-Doc and CohE-CttA-Doc complexes were probed throughout the measurement. In this manner, thousands of force--extension curves was using contour length analysis, and identifying only those traces in which two CBM unfolding length were identified by filtering the data sets using contour length analysis, such extension extension actures were at ropture event scatter plots describing the rupture of the XDoc:CohE complexes.

#### ASSOCIATED CONTENT

#### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/jacs.9b06776.

Protein sequences, ITC data, rupture force histograms, dynamic force spectrum combining in vitro and in silico experiments (PDF)

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#### Article

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S1



S2















# 3.4 Surface Chemistry for Affinity Assays

In this publication, we utilize enzymatic, covalent surface immobilization strategies to perform kinetic rate measurements of receptor-ligand pairs on a biolayer interferometer. This allows us to provide the same anchoring geometry as in SMFS experiments. It also allows us to exclude possible artifacts introduced by the usually employed sensor functionalization via primary amines. My contribution to this publication was to the design of the study, the execution of experiments as well as as the writing of the manuscript. I also wrote analysis software capable of fitting kinetic models to the acquired data.

# Enzyme-Mediated, Site-Specific Protein Coupling Strategies for Surface-Based Binding Assays

Wolfgang Ott<sup>+</sup>, Ellis Durner<sup>+</sup>, and Hermann E. Gaub

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#### GDCh Angewandte Communications International Edition: DOI: 10.1002/anie.201805034 Protein Immobilization Hot Paper German Edition: DOI: 10.1002/ange.201805034 Enzyme-Mediated, Site-Specific Protein Coupling Strategies for Surface-Based Binding Assays Wolfgang Ott<sup>+</sup>, Ellis Durner<sup>+</sup>, and Hermann E. Gaub\* Abstract: Covalent surface immobilization of proteins for charges, which in turn may cause unfolding and aggregation issues.<sup>[6]</sup> Additionally, proteins with a low isoelectric point binding assays is typically performed non-specifically via lysine residues. However, receptors that either have lysines near their might not be sufficiently protonated, and thus remain negatively charged. binding pockets, or whose presence at the sensor surface is electrostatically disfavoured, can be hard to probe. To over-Another challenge with non-specific surface chemistry is come these limitations and to improve the homogeneity of that proteins often contain more than one reactive residue, surface functionalization, we adapted and optimized three which leads to inhomogeneous surface anchoring. Consedifferent enzymatic coupling strategies (4'-phosphopantequently, sensorgrams of a binding experiment represent the theinyl transferase, sortase A, and asparaginyl endopeptidase) superimposed response of multiple populations of differently for biolayer interferometry surface modification. All of these attached receptors. Varying attachment sites may strongly enzymes can be used to site-specifically and covalently ligate influence binding kinetics solely due to the molecules' proteins of interest via short recognition sequences. The orientation. The binding behaviour can be altered or binding enzymes function under mild conditions and thus immobilizamay even be prevented, especially with receptors immobition does not affect the receptors' functionality. We successfully lized via reactive residues close to their binding interface (Figure 1 A).<sup>[7]</sup> employed this enzymatic surface functionalization approach to study the binding kinetics of two different receptor-ligand pairs. In this study, we expand the toolbox for surface functionalization by adapting advances in enzyme-based protein The binding properties of receptor-ligand complexes have modification strategies to overcome the limitations of nonbeen studied in vitro with numerous assays developed during specific pull-down strategies in binding assays. The employed the last decades.<sup>[1-3]</sup> Mainly, covalent approaches for receptor enzymes are a 4'-phosphopantetheinyl transferase from immobilization have been established to precisely determine Bacillus subtilis (Sfp),<sup>[8]</sup> an evolved sortase A (SrtA) from Staphylococcus aureus (d59SrtA, P94R/D160N/D165A/ K190E/K196T),<sup>[9]</sup> and an engineered asparaginyl endopeption-rate $(k_{on})$ , off-rate $(k_{off})$ , and equilibrium constant $(K_d)$ .<sup>[4]</sup> For these methods, the receptor is immobilized onto a surface dase from the plant Oldenlandia affinis (OaAEP1) and a change in signal upon ligand application is evaluated. (C247A).<sup>[10]</sup> All of these enzymes recognize specific amino While sometimes the terminology "ligand-analyte" is used, throughout this article the molecule immobilized to the acid sequences (tags) and covalently attach these tags to other sensor surface is called the receptor and its binding partner the ligand. In general, accessible side chains of corresponding amino acids (amine-, carboxyl-, or thiol-groups)<sup>[5]</sup> can be employed to covalently link the receptor to a surface. However, non-specific attachment requires an electrostatically driven surface pre-concentration step, where the pH and salt conditions of the buffer must be chosen such that the sensor surface and the receptors are predominantly oppositely charged. This pre-concentration step requires a buffer of low ionic strength in order to prevent screening of surface [\*] Dr. W. Ott,<sup>[+]</sup> E. Durner,<sup>[+]</sup> Prof. Dr. H. E. Gaub Luchrstuhl für Angewandte Physik and Center for NanoScience Luchrstuhl für Angewandte Physik and Center for NanoScience Luchig-Maximilians-Universität München Amalienstrasse 54, 80799 Munich (Germany) Figure 1. Schematic of BLL Kinetics, A) Non-specific immobilization of the receptor on the sensor in different geometries as a result of several accessible amine-groups. B) Specific and site-directed immobilization and of a receptor to a sensor. All receptors are homogeneously orientated The red arrows in (A) represent different binding geometries with Center for Integrated Protein Science Munich (CIPSM), Ludwig-Maximilians-Universität München Butenandtstrasse 5–13, 81377 Munich (Germany) possibly different kinetics, whereas specific attachment (B) provides a uniform population of binders. C) The principle of a BLI kinetic E-mail: gaub@lmu.de experiment. A receptor-functionalized sensor is immersed into a ligand [+] These authors contributed equally to this work solution. The increasing signal shows binding of the ligand. When the sensor signal has reached a steady state, the rates of ligand associa-Supporting information (including experimental details) and the tion and dissociation are equal—() the system has reached equilibrium. The sensor is then moved to a buffer solution, the receptor starts to ORCID identification number(s) for the author(s) of this article can Ó be found under https://doi.org/10.1002/anie.201805034. dissociate and the detected signal decreases again. 12666 Wiley Online Library © 2018 Wiley-VCH Verlag GmbH & Co. KGaA, Weinheim Angew. Chem. Int. Ed. 2018, 57, 12666-12669

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amino acid sequences (SrtA and OaAEP1) or to Coenzyme A (CoA; Sfp). In case of SrtA and OaAEP1, the tags have to be at the termini of the protein, whereas the ybbR-tag (11 amino acids) recognized by Sfp can also be internal (if accessible) since its ligation mechanism does not rely on peptidase activity. These tags can be fused to proteins and employed in surface pull-down strategies, hence allowing homogeneous loading of a surface (Figure 1B). In single-molecule approaches, such as single-molecule force spectroscopy, sitespecific reactions<sup>[11-13]</sup> are already well established and ensure reliable mechano-probing of receptor-ligand systems without removing the proteins from the surface or cantilever. We adapted these enzyme-based techniques, which enabled us to link a receptor of interest to a sensor surface in very mild reaction conditions while using only low micromolar quantities of receptor.

We chose a biolayer interferometer (BLI) as a development platform because of its fast and flexible assay format. However, it should be noted that the approach presented here is applicable to other surface sensitive techniques, such as surface plasmon resonance (SPR) or quartz crystal microbalance (QCM), since the receptor immobilization relies on the same chemistry. The underlying principle of a BLI makes use of light reflection at interfaces between media of different optical densities to analyse the spectral shift of interference signals upon binding—which effectively modifies the optical path length—to the sensor.<sup>[14,15]</sup> The interference signal changes whenever binding/unbinding to the sensor fibre occurs (Figure 1 C).

In order to establish our enzyme-based BLI binding assays, we selected two different systems (Table S1, Figure S1 in the Supporting Information). Firstly, we chose GFPbinding nanobodies (LaG9).<sup>[16]</sup> Nanobodies are small functional single-chain antibodies<sup>[17]</sup> and are popular tools in diagnostic as well as in therapeutic applications. As a second system, we chose the mechanically highly robust cohesindockerin type III complex (CohE–XDocIII) from *Ruminococcus flavefaciens*. As previous single-molecule force spectroscopy studies have shown, its unbinding behaviour under force depends on the anchoring geometry of the cohesin. When immobilized via its C-terminus, a most probable rupture force of around 700 pN (at 100 nN s<sup>-1</sup>)<sup>[18]</sup> is observed, in contrast to only 100 pN (at 0.7 nN s<sup>-1</sup>)<sup>[19]</sup> when anchored via its N-terminus. With the site-specific immobilization strategies presented here (Figure 2), we were able to probe the geometry dependence in the absence of force.

Experimental details, traces for Sfp-, SrtA- and OaAEP1based sensor modifications (Figure S4–S19), and an overview of all possible immobilization geometries (Figure S20) can be found in the Supporting Information. Once the sensors were site-specifically loaded with the protein of interest, they were equilibrated in the same measuring buffer throughout all experiments.

In order to compare the different immobilization strategies, a kinetic binding series with each coupling approach was recorded. Figure 3 A shows an example sensorgram of an SrtA-based experiment. Despite using another GFP variant which differs in the binding epitope (Figure S21), we obtained similar binding kinetics to the reported ones ( $K_d$  = 3.5 nm,

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Figure 2. Overview of the different covalent, site-specific immobilization techniques. Left: Sfp catalyses the reaction between ybbR-tag of the TagCFP2 and coenzyme a (CoA). First, the amine-group of PDEA reacts with the EDC/NHS-activated carboxyl groups of the sensor. PDEA can then undergo a thiol exchange reaction with CoA, which presents a free thiol. Middle: SrAI links C-terminal LPETGG with N-terminal GGG. In the case shown here, a C-GGGGG peptide was reacted with the EDC/NHS-PDEA-activated sensor. Right: OaAEP1 recognizes the C-terminal amino acids GLP. EDC/NHS-activated sensors were reacted with the amine-groups of a KK-GSGS-NGL peptide. All three immobilization methods yield a homogeneous TagGFP2-modified sensor ready for binding kinetic measurements.

 $k_{on} = 2.3 \times 10^6 \,\mathrm{m^{-1} s^{-1}}, \ k_{off} = 8.0 \times 10^{-3} \,\mathrm{s^{-1}})$  determined with SPR<sup>[16]</sup> (compare Figure 3B).

The obtained kinetic rates were independent of the functionalization method (specific and non-specific). The site-specific approach anchors proteins at their termini and decreases the chance of binding site obstruction (spatial separation of surface coupling and ligand binding), which thus allows us to determine the unaltered (un)binding rates. This increased reliance is an intrinsic advantage of our site-specific surface functionalization. Based on this, we can compare the data with the non-specifically anchored proteins and conclude that the multiple lysine anchoring possibilities do not obstruct the binding behaviour in the case of TagGFP2-LaG9 interaction. TagGFP2 contains 17 lysines that may take part in the non-specific immobilization procedure. Hence, it is not surprising that enough primary amines non-adjacent to the binding epitope that do not disturb binding are available as anchoring sites. Other receptor-ligand systems might be more strongly affected by the non-specific anchoring (see cohesindockerin interaction below). Especially if the surface area at the ligand binding site is charged such that it is electrostatically favoured to make surface contact during the pre concentration step, the binding site could be obstructed.



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Figure 3. Binding kinetics of the TagGFP2-LPETGG receptor with the GLP-LaG9-HIS ligand. A) An example sensorgram of LaG9 ligand binding to the TagGFP2-LPETGG receptor at different concentrations (25, 50, 100 and 200 nм). B) The kinetic rates obtained from performing global fits to sensorgrams for each immobilization method. Values and the respective standard errors are obtained from triplicates.

TagGFP2-NGL and TagGFP2-ybbr could not be fused to a BLI sensor. However, both tags were functional to fuse protein domains in an in vitro bulk reaction. Thus, it appears that both tags are sterically hindered by the GFP domain when used for surface functionalization. A longer linker between GFP domain and the recognition sequence could possibly provide both enzymes (Sfp, OaAEP1) more space to ligate the protein to the sensor.

This enzyme-based and site-specific surface reaction also allowed us to probe the inverse geometry with the nanobody now immobilized to the sensor. Two kinetic titration series, one using the SrtA-based and the other using the non-specific immobilization approach, were recorded (Figure S10). Fits deviated notably from a 1:1 binding model, which might be explained by either ligand-surface interactions or by potential avidity effects should TagGFP2 present more than one binding interface. The ability to site-specifically anchor both binding partners allows us to exclude that the deviations from a simple 1:1 binding model stem from heterogeneous surface preparation due to non-specific protein anchoring. Based on the conducted experiments, we were able to show that all three enzymes can be used for sensor functionalization.

The advantages of defined surface immobilization emerge more clearly when investigating the CohE–XDocIII interaction of *R. flavefaciens.* This cohesin–dockerin pair has already been characterized in bulk studies<sup>[20]</sup> as well as singlemolecule studies.<sup>[18,19]</sup> However, we were not previously able to non-specifically immobilize the cohesin in a functional state (Figure S3), possible due to a lysine in its binding pocket.<sup>[18]</sup> However, not only were we able to attach the cohesin site-specifically and in biologically active form with the enzymatic approach, we were also able to do so from

unbinding behaviour of this complex under external force was shown to strongly depend on the anchoring geometry of the cohesin. The complete sensorgram of the sensor modification can be found in Figure S11. Because of the evidently very low off-rate of the cohesin-dockerin complex, and because no regeneration conditions could be found to force ligand dissociation, we chose to perform a kinetic titration experiment. Full dissociation of the complex would take too long and by far exceed the four hours of total experiment time suggested for BLI. Longer measurements would suffer from evaporation of liquids in the microwell plate, thus falsifying concentrations. Initial experiments showed that sensor-drift effects seemed to exceed the actual ligand dissociation due to the low off-rate (Figure S12). Thus, as recommended,<sup>[21]</sup> we modified our protocol such that both sensors are loaded with a receptor; in our case they were functionalized sitespecifically with cohesin. For referencing, one sensor was only dipped into measurement buffer while the other was presented with ligand. Despite having minimized drift, it became clear that the off-rate of the complex is too low to be assessed through BLI. Too little dissociation occurred so that adequate fitting of the data was not possible. For an exact determination of the off-rate, alternative techniques such as SPR or QCM might be more promising, since they are not limited by evaporation effects and can thus measure for extended periods of time. The apparent low off-rate is common in cohesin–dockerin systems.<sup>[22,23]</sup> However, a qualitative statement about the (un)binding behaviour of the cohesin-dockerin complex is possible, namely that on- and off-rates appear to be independent of the anchoring geometry (Figure 4).

either terminus. This was of particular interest since the

This stands in contrast to force spectroscopy experiments.<sup>[18,10]</sup> where the anchoring geometry strongly changes the force necessary to dissociate cohesin and dockerin. By comparing the findings obtained by force spectroscopy with those from site-directed BLI, we can conclude that the different rupture forces are indeed a consequence of force propagation within the receptor–ligand complex, rather than an artefact caused by surface effects or the employed anchoring chemistry.

In summary, the presented strategy provides an efficient means to covalently and site-specifically couple receptors under mild reaction conditions. The employed tags are all small and should not influence the overall functionality of a protein. This makes it viable to use the same constructs for characterization in a surface-based assay as well as for other bulk and single-molecule studies. Moreover, these small tags can be further used for post-translational protein modifica-tions, that is, attachment of a fluorescent dye<sup>[24]</sup> or an additional protein domain,<sup>[13]</sup> or as a pull-down technique.<sup>[13]</sup> Hence, label-free and label-dependent techniques can be used with the same batch of proteins. While the enzymatic approaches presented here are shown for sensor modification in BLI, they can be easily adopted for other assays, such as SPR or QCM. Overall, the site-directed and covalent immobilization techniques present a viable, easily implementable alternative to the non-specific approach. Additionally, no buffer conditions suitable for pre-concentration need to be

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Figure 4. Comparison of N- and C-terminal immobilization of CohE. A) Site-specific surface chemistry enables N- and C-terminal anchoring to be distinguished. B) Kinetic titration series of an OaAEP1-immobiized GLP-CohE (Top) and a SrtA-immobilized CohE-LPETGG (Bottom) sensor with XDocIII ligand binding. SrtA was employed to covalently couple CohE-LPETGG to a polyglycine-modified sensor. OaAEP1 immo-bilized the GLP-CohE to an NGL-modified sensor. A kinetic titration series was performed by subsequently dipping the sensor into different concentrations of XDocIII (80, 160, 320 and 640 nm) with dissociation steps in measurement buffer.

found, which provides a faster way of establishing assays on new receptor-ligand interactions. Also, due to the specific nature of the surface coupling, signal arising from ligands that bind to non-specifically adsorbed receptors can be subtracted since the reference sensor can be prepared by simply omitting the coupling enzyme. Most importantly, receptor-ligand interactions that were previously inaccessible due to reactive residues in their binding interface or due to electrostatic repulsion can now be site-specifically immobilized and characterized with the enzymatic approaches.

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## Conflict of interest

The authors declare no conflict of interest.

Keywords: biolayer interferometry  $\cdot$  protein immobilization  $\cdot$ protein ligation  $\cdot$  site-specific protein modification  $\cdot$ sortase A, Sfp, OaAEP1

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## SUPPORTING INFORMATION

### Methods

All materials employed in this study were at least of analytical purity grade and purchased from Sigma Aldrich Chemie GmbH (Taufkirchen, Germany) and Carl Roth + Co. KG (Karlsruhe, Germany). All buffers were filtered (0.22 µm) and degassed via sonication. The pH of the buffers were adjusted at room temperature.

### Cloning

The employed nanobody LaG9 was published and characterized by Firdy *et al.*<sup>[11]</sup>. It was necessary to modify the plain nanobody sequence with the corresponding tags for enzyme-catalyzed pulldown. All constructs were subcloned into a modified pET28a vector with Gibson assembly<sup>[21]</sup>. In order to perform Gibson assembly, the inserts were ordered as a gene string with overlapping sequences to the plasmid, up- and downstream of the insert or amplified with primers containing the overlaps. The Gibson assembly reaction was performed at 50°C for 1 hour (10 µl 2x HiFi MasterMix, New England Biolabs, Ipswinch, MA, USA, mixed with 0.1 nmol of vector and 0.2 nmol of insert). Later modifications were done with a plasmid linearization reaction via PCR and recircularization reaction (4.5 µl of PCR product, 1 µl ATP (10 mM, Thermo Fisher Scientific Inc., Waltham, MA, USA), 0.5 µl PEG-6000 (Thermo Fisher Scientific Inc.), 1 µl Durl (Thermo Fisher Scientific Inc.), 1 µl Attre (10 WL), Thermo Fisher Scientific Inc.). This reaction was performed to 15 minutes at 37°C and 45 minutes at 22°C.

DHSo cells (Life Technologies GmbH, Frankfurt, Germany) were thawed on ice, 1 µl of the reaction mix was added to the cells to transform them. After 30 minutes on ice a heat shock at 42°C for 1 minute was done. Finally, the cells grew shaking (850 rpm) for 1 hour at 37°C in 1 ml of SOC Medium. Usually 100 µl of the transformed culture was used to streak on a LB-Kanamycin containing plate. Cells were incubated overnight at 37°C. Plasmids of a small amount of clones (less than 5) were amplified and sent to sequencing to verify their sequence.

The gene for the asparaginyl endopeptidase OaAEP1 was ordered via gene string (Thermo Fisher Scientific Inc.) and subcloned via Gibson assembly. Sortase A (StA) was cloned and modified as described by Durner *et al.*<sup>[3]</sup>. YbbR-HIS-XDockerinIII (XdocIII), CohesinE-HIS-LPETGG and GLP-CohesinE-HIS (CohE) were created with the above described linearization and recircularization reaction. The underlying constructs were ybbR-HIS-CBM-XDocIII and CohE-CBM-HIS-ybbR<sup>[4]</sup>. 4'-Phosphopantetheinyl transferase (Sfp) was already available in the laboratory.

#### Protein Expression and Purification

*NiCo21(DE3)* (New England Biolabs) cells were transformed with the appropriate plasmid (50 ng), and incubated overnight at 37°C on LB-Agar plates with kanamycin (50 ng/µl). One clone was inoculated in 5 ml LB-Kanamycin liquid culture (37°C, 200 rpm, 12-16 hours). This starter culture was used to inoculate a larger 400 ml autoinduction media (270/M-5052) culture (100 ng/µl kanamycin)<sup>[5]</sup>. Usually the cultures were incubated for 24 hours (4 hours 37°C, 20 nors 18°C; 120 rpm). Cells were harvested *via* centrifugation (6500 g, 20 minutes, 4°C), the supernatant discarded and the pellets frozen at -80°C until further use.

This starter culture was used to inoculate a larger 400 mi autoinduction media (27M-5052) culture (100 ng/µl kanamycin)<sup>50</sup>. Usualiy the cultures were incubated for 24 hours (4 hours 37°C, 20 hours 18°C; 120 rpn). Cells were harvested via centrifugation (6500 g, 20 minutes, 4°C), the supernatant discarded and the pellets frozen at -80°C until further use. The frozen pellets were resolubilized in 40 ml lysis buffer (50 mM TRIS (tris(hydroxymethyl)aminomethane)-HCI, pH 8.0, 50 mM NaCI, 5 mM MgCl<sub>2</sub>, 10 % (v/v) glycerol, 0.1 % (v/v) Triton X-100 and supplemented with 10 µg/ml DNase I and 100 µg/ml Iyaszyme). To enhance the chemical lysis, the cells were sonicated with a sonication lance twice for 7 minutes on ice (50 % Power, 50 % Cycle; Bandelin Sonoplus GM 70, tip: Bandelin Sonoplus MS 73, Berlin, Germany). The cell ysate was centrifuged (30000 g, 4°C) for 1 hour. The supernatant was filtrated (0.45 µm) and applied to a Ni-NTA column (5 ml HisTrap FF, GE Healthcare Europe GmbH, Freiburg, Germany), since all proteins contained a HIS-tag. The protein was eluted after washing (25 mM TRIS-HCI pH 7.8, 500 mM NaCI, 20 mM imidazole, 0.25 % (v/v) Tween 20, 10 % (v/v) glycerol) with 6 column volumes with elution buffer (25 mM TRIS-HCI, pH 7.8, 500 mM NaCI, 300 mM imidazole, 0.25 % (v/v) Tween 20 (v/v), 10% (v/v) glycerol). The eluted proteins were checked on a SDS-gel for digestion and purity.

Nanobody expressing cultures were inoculated in the autoinduction media minus lactose to prevent autoinduction (ZYM-505). They were grown at 37°C until an OD<sub>800</sub> of 4, induced with 0.1 mM IPTG and the temperature lowered to 18°C. Cells were harvested (6500 g, 10 minutes, 4°C) after 20 hours and directly lysed and purified. Nanobodies were purified with periplasmic extraction<sup>[6]</sup>. The cell pellet was resuspended in lysis buffer (100 mM TRIS-HCI, 500 mM Saccharose, 1 mM EDTA, pH 8.0) and incubated on a reaction tube roller (60 rpm, 4°C, 1 hour). A centrifugation step separated the insoluble part from the nanobody containing supernatant (3220 g, 4°C, 10 minutes). The supernatant was again centrifuged (30000 g, 4°C, 1 hour) and filtered (0.22 µm). The filtrated supernatant was applied on a Ni-NTA column as described above.

minutes). The supernatant was again centrifuged (30000 g, 4°C, 1 hour) and filtered (0.22 µm). The filtrated supernatant was applied on a Ni-NTA column as described above. The buffer of the protein solution was exchanged with ZebaSpin columns (Zeba spin desalting columns 7K, Thermo Fisher Scientific Inc.) to 50 mM TRIS-HCI, pH 7.2, 72 mM NaCl, 1 mM CaCl<sub>2</sub> and 20 % (v/v) of glycerol was added. Nanobodies were stored in 20 mM HEPES pH 8.0, 150 mM NaCl and 10 % (v/v) glycerol. Small aliquots of all proteins were flash frozen in liquid nitrogen and stored at -80°C.







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No binding signal could be detected while dipping into a well with dockerin (1000 nM concentration). The concentration of dockerin was many fold above the reported K₂ of 20.83 nM<sup>(9)</sup>. As a consequence, it can be said that with the non-specific approaches no biofunctional cohesin could be immobilized, requiring alternative immobilization methods.

### Specific Sensor Immobilization

If the specific protocol relied on a disulfide exchange reaction, the now amine reactive sensors were dipped for 10 minutes into a 40 mM solution of PDEA (dissolved in 50 mM borate buffer, pH 8.5) to covalently immobilize the thiol reactive compound to the biosensor. In order to quench remaining amine-reactive EDC/NHS-groups, the PDEA modified sensors were incubated for 10 minutes in 100 mM TRIS+HCI, pH 8.5.

The sensors for TagGFP2-experiments were functionalized and the interaction measured in one experiment. Sensors for Cohesinexperiments were functionalized and parked in measurement buffer before starting the actual binding experiment. This two step approach was necessary to undercut the recommended 4 hours timeframe for an experiment to avoid evaporation of the solutions in the microwell plate.

Sfp has the benefit that the ybbR-tag can be internally introduced in loops and unstructured regions. OaAEP1 works under acidic conditions and might be attractive to immobilize protein complexes with a low isoelectric point. Regarding TagGFP2 immobilization SrtA proved to be most robust.

#### Sfp-based Sensor Modification

Sensors were modified with PDEA as described under basic sensor preparation. Since Coenzyme A (CoA, Merck KGaA Darmstadt, Germany) contains an accessible thiol-group, the PDEA modified sensors can react with them. 1 mM CoA in coupling buffer (50 mM sodium phosphate, 50 mM NaCl, 10 mM EDTA, pH 7.2) was fused to the sensors for 10 minutes. A final quenching step (25 mM TRIS-HCI, pH 7.2, 72 mM NaCl, 1 mM CaCl<sub>a</sub>, 0.1 % (v/v) casein, 0.1 % (v/v) Tween-20) passivated the remaining sensor surface against non-specific attachment (10 minutes). The receptor of choice with a ybbR-tag was covalently attached to the CoA by enzymatic ligation via Sfp (5 µM receptor, 1 µM Sfp in 25 mM TRIS-HCI, pH 7.2, 72 mM NaCl, 1 mM CaCl<sub>a</sub>, 5 mM MgCL, 0.1 % (v/v) casein, 0.1 % (v/v) Tween-20) for 6 minutes until a surface density of ca. 0.6 - 0.8 nm was achieved. The sensors are now ready for kinetic binding measurements.

Figure S4 shows an example signal trace for the Sfp functionalization. A carboxyl-sensor, in MES-buffer equilibrated, dips into the EDC/NHS solution (Figure S4, 1). An increase in signal can be detected. The following PDEA step also shows binding (Figure S4, 2). Next an amine quenching is necessary, because CoA also contains amine-groups. When linked via its amine-groups, CoA cannot be ligated to the ybbR-tag (Figure S4, 3). The thiol exchange reaction between CoA and PDEA can now take place. It should be noted that CoA does not generate any signal upon binding at concentrations below 20 mM (Figure S4, 4). Next, the sensor needs to be passivated with a passivation agent suited for the receptor:ligand pair. In this case, casein worked best (Figure S4, 5). The site-specific Sfp reaction can now take place. In yellow a trace is shown with added Sfp, in blue only the receptor was added, Sfp was omitted (Figure S4, 6). We therefore can correct for the content of non-specific binding during the enzyme-catalyzed reaction (Figure S4, 6, Yellow Trace). The positive binding signal can be later corrected with the non-specific disorption value. The almost vertical jumps in the signal traces are due to changing buffer conditions which cause an abrupt change in the signal.



Supplemental Figure S4. Exemplary sensorgram for the immobilization of ybbR-TagGFP2. First the carboxyl-group is activated with EDC/NHS (1) which reacts with the amine-group of PDEA (2), (3) TRIS quenches all the unreacted, but active EDC/NHS-Groups. Coerazyme A (4) shows no increase in the sensorgram but nevertheless links to the PDEA via its thior-group, (5) is a passivation step with case in to prevent non-specific adsorption of proteins to the sensor. (6) shows the specific Stp reaction (yellow) and the non-specific adsorption of the protein mix to the sensor (bue). (7) is the equilibration to baseline in measurement buffer.

The functionalized sensors are then dipped into wells with GLP-LaG9-HIS with different concentrations (25, 50, 100 and 200 nM) and buffer to measure dissociation of the nanobody (Figure S5).



























```
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Data Analysis Script
 import numpy as np
 import matplotlib as mpl
 import matplotlib.pyplot as plt
 from scipy.optimize import curve fit
 import sys, os
 import re
 import csv
 import argparse
 import StringIO
 import string
import string
import matplotlib.pylab as pylab
 import xml.etree.ElementTree as ET
 import struct
 import glob
 fontsize = 16
rontsize = 16
params = { 'font.size': fontsize,
            'legend.fontsize': fontsize,
            'axes.labelsize': fontsize,
            'axes.titlesize': fontsize,
            'xtick.labelsize': fontsize,
                 'ytick.labelsize': fontsize,
                 'axes.labelpad':
                                              3.0,
                 'text.usetex': True,
           'text.latex.preamble': [
                 r"\usepackage{helvet}",
r"\usepackage{sansmath}",
r"\renewcommand\familydefault{\sfdefault}"
                  r"\usepackage{sfmath}
                  r"\usepackage[detect-all]{siunitx}",
 ]
                }
 pylab.rcParams.update(params)
mpl.rc("figure", figsize=(12, 6))
 lw=3
 papercolors = [\
(0.988, 0.752, 0.015), #Yellow
(0, 0.403, 0.603), #Blue
 (0.192, 0.6, 0.4), #Green
(0.8, 0.207, 0.090), #Bright Red
  (0.529, 0.403, 0.666), #Bright Purple
(0.529, 0.403, 0.666), #Bright Purple
(0.196, 0.141, 0.368), #Medium Purple
(0.180, 0.074, 0.196), #Dark Purple
(0.396, 0.078, 0.054), #Dark Red
(0.909, 0.192, 0.541), #Pink
(0.796, 0.6, 0.015), #Brown
(0.450, 0.760, 0.686), #Dark Teal
(0.603, 0.803, 0.803)] #Bright Teal
 *****
                                                               20
```

```
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   def string2float(value):
              value = value.replace(',','.')
return float(value)
 def plotadd(axis, dd, cycleparm, fit_global, i):
    sensor = cycleparm['sensor']
    coefoff = cycleparm['offrate_parms']
               if fit_global == True:
                         koff = cycleparm['offrate_parms_global'][0]
                          coefon = cycleparm['onrate_parms_global']
              else:
                         koff = coefoff[0]
                          coefon = cycleparm['onrate_parms']
               conc = cycleparm['conc'
              offratefit_offset = cycleparm['offratefit_offset']
             baseline_duration = 10# in seconds
baseline_length = (baseline_duration/\
               (dd[sensor[0]][1]-dd[sensor[0]][0])).astype(int)
             (u(sensor(b))[1]-u(sensor(b)[1]()).asype(Int)
start = cycleparm['assoc_start']-baseline_length
concstring = '{0:.2e}'.format(conc)
axis.plot(dd[sensor[0]][start:\
cycleparm['dissoc_end']] - cycleparm['t0_assoc'],
               dd[sensor[0]+"_resp"][start:cycleparm['dissoc_end']],
             '-', color=papercolors[i], markevery=1,
label=r'$\SI{{{}}}{(nano\Molar}}$'.format(int(conc*1e9)),
              zorder=0)
   ####plot offrate####
              if offratefit_offset == True:
                         plt.plot(dd[sensor[0]][cycleparm['dissoc_start']:\
             cycleparm['dissoc_end']] - cycleparm['dissoc_start'].\
cycleparm['dissoc_end']] - cycleparm['dissoc_'],
offrate(dd[sensor[0]][cycleparm['dissoc_start']:\
cycleparm['dissoc_end']],koff,
cycleparm['y0_dissoc'], cycleparm['t0_dissoc'], coefoff[1]),linestyle="--",
color="k", zorder=1)#, label='koff= {0:.2e}'.format(coefoff[0]))
life offerth file offerth
               elif offratefit_offset == False:
             plt.plot(dd[sensor[0]][cycleparm['dissoc_start']:\
cycleparm['dissoc_end']] - cycleparm['t0_assoc'],
offrate(dd[sensor[0]][cycleparm['dissoc_start']:\
cycleparm['dissoc_end']], koff, cycleparm['y0_dissoc'],
cycleparm['t0_dissoc'], offset=0),linestyle="--",color="k",
              zorder=1)
#, label='koff= {0:.2e}'.format(coefoff[0]))
   ####plot onrate####
              if cycleparm['fit_onrate'] == True:
            if coefon[1:].size > 1:#ymax unlinked
                          else:
                                     plt.plot(dd[sensor[0]][cycleparm['assoc_start']:\
             cycleparm['assoc_end']] - cycleparm['t0_assoc'],
onrate(dd[sensor[0]][cycleparm['assoc_start']:\
cycleparm['assoc_end']], coefon[0], koff,
               coefon[1], conc, cycleparm['t0_assoc']),linestyle="--",
                                                                                                                                    21
```

```
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       color="k", zorder=1)
       #,Label='kon= {0:.2e}'.format(coefon[0]))
 def plotadd_kinetic(axis, dd, cycleparm, fit_global, i):
      sensor = cycleparm['sensor']
coefoff = cycleparm['offrate_parms']
if fit_global == True:
            koff = cycleparm['offrate_parms_global'][0]
            coefon = cycleparm['onrate_parms_global']
       else:
            koff = coefoff[0]
            coefon = cycleparm['onrate parms']
       conc = cycleparm['conc']
       offratefit_offset = cycleparm['offratefit_offset']
       plt.plot(dd[sensor[0]][cycleparm['assoc_start']:\
       cycleparm['assoc_end']], onrate_kinetic(dd[sensor[0]]\
[cycleparm['assoc_start']:cycleparm['assoc_end']], coefon[0],
       koff, coefon[1], cycleparm['ystart'], conc, cycleparm['t0_assoc']),
linestyle="--",color="k")
#,label='kon= {0:.2e}'.format(coefon[0]))
 ##offset in kinetic titration series doesn't make sense
        if offratefit_offset == True:
              plt.plot(dd[sensor[0]][cycleparm['dissoc_start']:\
       pl:/plot(du[sensor[0]][cycleparm['dissoc_start ]:\
    [cycleparm['dissoc_start']:cycleparm['dissoc_end']],
    koff, cycleparm['y0_dissoc'], cycleparm['t0_dissoc'], coefoff[1]),
    linestyle="--",color="k")
#, label='koff= {0:.2e}'.format(coefoff[0]))
 #
 #
 #
 #
 #
      #, table koff= {0:.26} .format(cofoff[0]))
if offratefit_offset == False:
    #for i in range(1, Len(offrates)+1):
    plt.plot(dd[sensor[0]][cycleparm['dissoc_start']:\
    cycleparm['dissoc_end']], offrate(dd[sensor[0]]\
    [cycleparm['dissoc_start']:cycleparm['dissoc_end']],
       koff, cycleparm['y0_dissoc'], cycleparm['t0_dissoc'], offset=0),
       inestyle="--", color="k")
#, label='koff= {0:.2e}'.format(coefoff[0]))
def plotadd_data(axis, dd, cycleparm, i):
    sensor = cycleparm['sensor']
       conc = cycleparm['conc']
       baseline_duration = 60# in seconds
       baseline_length = (baseline_duration/\
       (dd[sensor[0]][1]-dd[sensor[0]][0])).astype(int)
       if i == 0:
            start = cycleparm['baseline_start']
       else:
            start = cycleparm['assoc_start']
       axis.plot(dd[sensor[0]][start:\
      cycleparm['dissoc_end']], dd[sensor[0]+"_resp"]\
[start:cycleparm['dissoc_end']],
'-', color=papercolors[i], markevery=1, zorder=0,
      label=r$\SI{{\}}{{\nan\Molar}}$'.format(int(conc*1e9)))
#Label='{0:.2e}M'.format(conc)
        axis.plot(time, response, color='black')
 def plotrates(axis, onrate, onrate_error, offrate, offrate_error):
       kd = offrate/onrate
       onstring = "k_{(0)}: SI_{(0.2e)} \{(per\mole\per\second)\} \pm
                                                                 22
```

```
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 \SI{{{1:.2e}}}{{}}.format(onrate,onrate_error)
      offstring =
                       "$k_{{off}}
 STString = $k_{(0<sup>+</sup>)}
\SI{{0:.2e}}{{per\second}}pm\SI{{1:.2e}}{{}}*.format(offrate, offrate_error)
kdstring = "$k_d: \SI{{}}{{per\nano\Molar}}".format(int(kd*1e9))
axis.annotate(offstring+'\n'+onstring+'\n'+kdstring, xy=(0.5, 0.1),
xycoords='axes fraction')
 def plotformat(axis):
      axis.legend(loc='upper left', shadow=False, frameon=False)
      despine(axis, top=1, bottom=0, left=0,right=1,xaxis=0,yaxis=0)
      plt.axvline(
      axis.set_xlabel('t $[\si{\second}]$')
      axis.set_ylabel('RU $[\si{\\nano\metre}]$')
 def despine(axis,top=1,bottom=1,left=1,right=1,xaxis=1,yaxis=1):
      if xaxis:
            axis.xaxis.set_visible(False)
      if yaxis:
           axis.yaxis.set_visible(False)
      if top:
           axis.spines["top"].set_visible(False)
      if bottom:
           axis.spines["bottom"].set visible(False)
      if left:
           axis.spines["left"].set_visible(False)
      if right:
           axis.spines["right"].set_visible(False)
 def extract_steps(folder):
      assaydict={}
      datadict={}
      files = glob.glob(folder +'/*_*.frd')
files = [os.path.basename(x) for x in glob.glob(folder +'/*_*.frd')]
      for f in files:
           tree = ET.parse(f)
           root = tree.getroot()
           sensorcoord = root.find('./ExperimentInfo/SensorName').text
           assaydict[sensorcoord] = []
           xdatas = []
ydatas = []
           for entry in root.iter('AssayXData'):
                length = int(entry.attrib['Points'])
                xdatas.append(unpack_data(entry.text))
                if length != len(xdatas[-1]):
    print "x data length not as expected"
                      .
return 1
                assaydict[sensorcoord].append(length)
           for entry in root.iter('AssayYData'):
    length = int(entry.attrib['Points'])
    ydatas.append(unpack_data(entry.text))
                if length != len(ydatas[-1]):
    print "y data length not as expected"
                      .
return 1
      datadict[sensorcoord] = np.hstack(xdatas)
    datadict[sensorcoord + '_resp'] = np.hstack(ydatas)
return assaydict, datadict
 def unpack_data(data):
                                                          23
```

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      size = 4
      decoded = data.decode('base64')
      length = len(decoded)/4
      return struct.unpack('<{}f'.format(length), decoded)</pre>
def onrate_equ(t,kon, koff ,yequ, conc, t0_assoc):
    y = yequ*(1-np.exp(-((conc*kon)+koff)*(t-t0_assoc)))
      return y
def yequ_to_ymax(yequ, koff, kon, conc):
    return yequ/(conc/(conc+koff/kon))
 def kinetic_timeoffset(ystart, kon, koff, maxresp, conc):
      timeoffset = (1-(ystart/(maxresp*conc))*((koff/kon)+conc))
timeoffset = -np.log(timeoffset)/((kon*conc)+koff)
      return timeoffset
def onrate_back(t,kon, koff ,ymax, conc, t0_assoc):
    y = ymax*(conc/(conc+koff/kon))*(1-np.exp(-((conc*kon)+koff)*(t-t0_assoc)))
      return y
 def onrate(t, kon, koff, ymax, conc, t0_assoc):
      y = ymax*(1/(1+koff/(kon*conc)))*(1-np.exp(-((conc*kon)+koff)*(t-t0_assoc)))
      return y
def onrate_kineticequ(t,kon, koff ,yequ, ystart, conc, t0_assoc):
    y = yequ*(1-np.exp(-((conc*kon)+koff)*(t-t0_assoc)))+ystart
      return y
 def onrate_kinetic(t, kon, koff ,ymax, ystart, conc, t0_assoc):
      t_offset = kinetic_timeoffset(ystart, kon, koff, ymax, conc)
y = ymax*(1/(1+koff/(kon*conc)))*(1-np.exp(-((conc*kon)+koff))
*(t-t0_assoc+t_offset)))
      return y
def offrate(t, koff, y0, t0, offset):
    y = (y0-offset)*np.exp(-koff*(t-t0)) + offset
      return y
 def combined_global(t, *args):
      concs = args[0]
      t0s_off = args[2]
      t_on_idx = args[3]
      t_off_idx = args[4]
      ystarts = args[5]
      offsets = args[6]
      fitparms = args[7]
kon_global = fitparms[0]
      ymaxs = fitparms[1:-1]
      koff_global = fitparms[-1]
      calc = []
      for i in range(0, len(concs)):
    if len(ystarts) > 0:#kinetic fit
        if i == 0:
                      ystart = 0
                                                             24
```

```
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              else:
                  ystart = calc[-1][-1]
              calc.append(onrate_kinetic(t[start:(start+t_on_idx[i])],
     kon_global, koff_global, ymaxs[0], ystart, concs[i], t0s_on[i]))
          else:
              if len(ymaxs) > 1:
                   calc.append(onrate(t[start:(start+t_on_idx[i])],
     kon_global, koff_global, ymaxs[i], concs[i], t0s_on[i]))
              else:
                   calc.append(onrate(t[start:(start+t_on_idx[i])],
     kon_global, koff_global, ymaxs[0], concs[i], t0s_on[i]))
start += t_on_idx[i]
          calc.append(offrate(t[start:(start+t_off_idx[i])],
     koff_global, calc[-1][-1], t0s_off[i], offsets[i]))
start += t_off_idx[i]
return np.hstack(calc)
def offrate_global(t, koff, *args):
     y0s = args[0]
     t0s = args[1]
     offsets = args[2]
t_idx = args[3]
     offrate_calc = []
     start = 0
     for i, idx in enumerate(t_idx):
         offrate_calc.append(offrate(t[start:(start+idx)],
     koff, y0s[i], t0s[i], offsets[i]))
start += idx
     return np.hstack(offrate_calc)
 def onrate_global(t, *args):
     koff = args[0]
concs = args[1]
     t0s = args[2]
     t_idx = args[3]
     ystarts = args[4]
     fitparms = args[5]
kon_global = fitparms[0]
     ymaxs = fitparms[1:]
     onrate_calc = []
     start = 0
     for i, idx in enumerate(t_idx):
          if len(ystarts) > 0:#kinetic fit
     onrate_calc.append(onrate_kinetic(t[start:(start+idx)],
kon_global, koff, ymaxs[0], ystarts[i], concs[i], t0s[i]))
          else:
              if len(ymaxs) > 1:
                   onrate_calc.append(onrate(t[start:(start+idx)],
     kon_global, koff, ymaxs[i], concs[i], t0s[i]))
              else:
                  onrate_calc.append(onrate(t[start:(start+idx)],
     kon_global, koff, ymaxs[0], concs[i], t0s[i]))
          start += idx
     return np.hstack(onrate_calc)
def displaysensors(experiment, sensor, dd, assaydict, displaysteps):
    fig, axis = plt.subplots()
    steptimes = []
                                                   25
```

```
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     stepidx = []
      steps = -1
      minresp = np.min(np.hstack((dd[sensor[0]+"_resp"][:],
     dd[sensor[1]+"_resp"][:])))
     dd[sensor[0]+"_resp"][:]))
     for step in assaydict[sensor[0]]:
          steps += step
          stepidx.append(step)
     steptimes.append(dd[sensor[0]][steps])
endidx = sum(stepidx[:displaysteps])
     axis.plot(dd[sensor[0]][:endidx], dd[sensor[0]+"_resp"][:endidx],
'-', color=papercolors[0], markevery=1, zorder=1,
     label='with OaAEP1')
      axis.plot(dd[sensor[1]][:endidx], dd[sensor[1]+"_resp"][:endidx],
       -', color=papercolors[1], markevery=1, zorder=1,
     label='without OaAEP1')
     plt.vlines(x=steptimes[:displaysteps], ymin=minresp-0.2, ymax=maxresp+0.2,
     color = 'black', linestyles='dashed')
for i in range(0, len(steptimes) -1)
     for i in range(0, displaysteps -1):
    textcoord = steptimes[i] + (steptimes[i+1] - steptimes[i])/2
    axis.text(textcoord, maxresp+0.2, "%d" % (i+1), rotation=0,
      horizontalalignment='center')
     despine(axis, top=1, bottom=0, left=1,right=1,xaxis=0,yaxis=1)
axis.set_xlabel('t $[\SI{}{\second}]$')
axis.set_ylabel('RU $[\SI{}{\nano\metre}]$')
     plt.legend(bbox_to_anchor=(0.98,1), loc="upper left", frameon=False)
     plt.savefig(experiment + ".pdf", bbox_inches="tight")
     plt.show()
      Label='{0:.2e}M'.format(conc))
#
      axis.plot(time, response, color='black')
def analyze_experiment(experiment, sensors, fit_kinetic=False, fit_onrate=True,
      reference=True, offratefit_offset = False, fit_global=False, uslocale=True,
      singlefit=False, skipsteps=0, fit_y0=False, ymax_linked=True,
     invert_sensors=False, koff_initial=1e=3, kon_initial=1e3,
offset_initial=0.1, displayonly=False, displaysteps = 7,
fit_stepwise = True):
     print(experiment.center(80, '*'))
     assaydict, dd = extract_steps(experiment)
sensorresults = []
     for sensor in sensors:
          if displayonly == True:
               displaysensors(experiment, sensor, dd, assaydict, displaysteps)
          else:
               fig, axis = plt.subplots()
               sensorresults.append(fit_sensor(sensor, dd, assaydict,
     fit_kinetic, fit_onrate, reference, offratefit_offset, uslocale,
singlefit, skipsteps, fit_y0, koff_initial, kon_initial, offset_initial,
      invert_sensors))
      if len(sensors) == 1 and fit_global == True:
          if fit_stepwise == True:
               koff_global_sol, koff_global_stderr = fit_global_off_onesensor(dd,
      sensors[0], sensorresults)
               kon_global_sol, kon_global_stderr = fit_global_on(dd, sensors[0],
```

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      sensorresults, koff_global_sol[0], fit_kinetic, ymax_linked)
            else:
                 kon_global_sol, kon_global_stderr, koff_global_sol, \
      koff_global_stderr = fit_global_combined(dd, sensors[0], sensorresults,
fit_kinetic, ymax_linked)
            for j, resultdict in enumerate(sensorresults[0]):
                 resultdict['offrate_parms_global'] = koff_global_sol
resultdict['offrate_stderr_global'] = koff_global_stderr
           resultdict['onrate_parms_global'] = kon_global_sol
resultdict['onrate_stderr_global'] = kon_global_sol
resultdict['onrate_stderr_global'] = kon_global_stderr
print(' global fitting '.center(80, '*'))
print "ymaxtest", kon_global_sol[1:]
print "global k_on: {0:.2e}, stdev: {1:.2e}".format(kon_global_sol[0],
kon_global_stder[0])
            kon_global_stderr[0])
           print koff_global_sol, "test"
print "global k_off: {0:.2e}, stdev: {1:.2e}".format(koff_global_sol[0],
            koff_global_stderr[0])
            print "global K_d: {0:.2e}".format(koff_global_sol[0]/kon_global_sol[0])
      if len(sensors) == 1 and displayonly == False:
            for j, resultdict in enumerate(sensorresults[0]):
    if fit_kinetic == True:
                      plotadd kinetic(axis, dd, resultdict, fit global, j)
                      plotadd_data(axis, dd, resultdict, j)
                 else:
                      plotadd(axis, dd, resultdict, fit_global, j)
      if len(sensors) > 1 and fit_kinetic == False and fit_global == False:
            for j, results in enumerate(sensorresults):
                 for resultdict in results:
                      plotadd(axis, dd, resultdict, fit_global, j)
      if singlefit == 1:
           print "handle me"
      if displayonly == False:
            plotformat(axis)
            ticks = [tick for tick in fig.gca().get_yticks() if tick >=0]
fig.gca().set_yticks(ticks)
      plt.savefig(experiment + ".pdf", bbox_inches="tight")
#plotrates(axis, kon_global_sol[0], kon_global_stderr[0],
      #koff_global_sol[0], koff_global_stderr[0])
      plt.show()
      return
 def fit_global_off_onesensor(dd, sensor, sensorresults):
      globaloff_t = []
globaloff_resp = []
      koffs_local = []
      y0s = []
      t0s = []
      offsets = []
       t_idx = []
      for i, resultdict in enumerate(sensorresults[0]):
      globaloff_t.append(dd[sensor[0]]\
[resultdict['dissoc_start']:resultdict['dissoc_end']])
globaloff_resp.append(dd[sensor[0]+"_resp"]\
      [resultdict['dissoc_start']:resultdict['dissoc_end']])
                                                            27
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             koffs_local.append(resultdict['offrate_parms'][0])
             y0s.append(resultdict['y0_dissoc'])
t0s.append(resultdict['t0_dissoc'])
       offsets.append(resultatic['offset_dissoc'])
    t_idx.append(globaloff_t[i].size)
globaloff_resp_arr = np.hstack(globaloff_t)
globaloff_t_arr = np.hstack(globaloff_t)
       koff_med = np.median(np.asarray(koffs_local))
       coefoff,covaroff = curve_fit(lambda t, koff: \
offrate_global(globaloff_t_arr, koff, y0s, t0s, offsets, t_idx),
globaloff_t_arr, globaloff_resp_arr,
p0=(1e-3), bounds=((0), (np.inf)), method='trf')

       return coefoff, np.sqrt(np.diag(covaroff))
 def fit_global_combined(dd, sensor, sensorresults, fit_kinetic, ymax_linked):
       global_t = []
global_resp = []
       kons_local = []
       koffs_local = []
       t0s_on = []
t0s_off = []
maxresponses = []
       concs = []
       t_on_idx = []
       t_off_idx = []
       ystarts = []
offsets = []
       for i, resultdict in enumerate(sensorresults[0]):
       global_t.append(dd[sensor[0]]\
[resultdict['assoc_end']])
       global_resp.append(dd[sensor[0]+"_resp"]\
[resultdict['assoc_start']:resultdict['assoc_end']])
             kons_local.append(resultdict['onrate_parms'][0])
             maxresponses.append(resultdict['onrate_parms'][1])
             concs.append(resultdict['conc'])
t0s_on.append(resultdict['t0_assoc'])
             t_on_idx.append(global_t[-1].size)
             global_t.append(dd[sensor[0]]\
       [resultdict['dissoc_start']:resultdict['dissoc_end']])
       global_resp.append(d[sensor[0]+"_resp"]\
[resultdict['dissoc_start']:resultdict['dissoc_end']])
koffs_local.append(resultdict['offrate_parms'][0])
t0s_off.append(resultdict['offset_dissoc'])
offsets.append(resultdict['offset_dissoc'])
t off idv_append(resultdict] t [ dissoc'])
              t_off_idx.append(global_t[-1].size)
             if fit_kinetic == True:
    ystarts.append(resultdict['ystart'])
       global_resp_arr = np.hstack(global_resp)
global_t_arr = np.hstack(global_t)
       koff_med = np.median(np.asarray(koffs_local))
       kon_med = np.median(np.asarray(kons_local))
       maxresponse_med = np.median(np.asarray(maxresponses))
if ymax_linked == True:
             fitparams_start = (kon_med, maxresponse_med, koff_med)
             no_of_onparms = 2
                                                                    28
```



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     else:
          fitparams_start = (kon_med,) + tuple(maxresponses)
          #fitparams_start = np.copy(maxresponses)
#fitparams_start = np.insert(fitparams_start,0,kon_med)
     fitparams = np.empty_like(fitparams_start)
coefon,covaron = curve_fit(lambda t, *fitparams: \
onrate_global(globalon_t_arr, koff, concs,
     points__biological starts, fitparams), globalon_t_arr, globalon_resp_arr,
p0=fitparams_start,# bounds=((0, 0), (np.inf, np.inf)),
     method='trf')
     return coefon, np.sqrt(np.diag(covaron))
def fit_sensor(sensor, dd, assaydict, fit_kinetic, fit_onrate,
     reference, offratefit_offset, uslocale, singlefit, skipsteps, fit_y0,
koff_initial, kon_initial, offset_initial, invert_sensors):
print(' local fitting '.center(80, '*'))
     if reference == True:
          if invert_sensors == False:
               dd[sensor[0]+"_resp"] = dd[sensor[0]+"_resp"] - \
     dd[sensor[1]+"_resp"]
          else:
              dd[sensor[0]+"_resp"] = dd[sensor[1]+"_resp"] - \
     dd[sensor[0]+"_resp"]
      no_of_concentrations = len(sensor[2:])
     if singlefit == 1:
          no_of_measurements = 1
     elif fit kinetic==True:
          no_of_measurements=(len(assaydict[sensor[0]])-1-skipsteps)/2
          #-1 for baseline
      else:
          no_of_measurements=(len(assaydict[sensor[0]])-skipsteps)/3
      #print
               'no_of_meas vs concs", no_of_measurements, no_of_concentrations
     if no_of_concentrations > no_of_measurements:
          print "more concentrations given than measurement cycles found in log"
     cycleparm = []
testcounter = 0
     for i in range(0, no_of_measurements):
          conc = sensor[i+2]
          if singlefit == 0 and fit_kinetic == False:
               assoc_start =\
     int(np.sum(assaydict[sensor[0]][:(skipsteps+1+(i*3))]))
              baseline_start =\
     assoc_start - int(assaydict[sensor[0]][skipsteps+(i*3)])
               t0_baseline = dd[sensor[0]][baseline_start]
               dissoc_start =\
     int(np.sum(assaydict[sensor[0]][:(skipsteps+2+(i*3))]))
     dissoc_end = int(np.sum(assaydict[sensor[0]]\
[:(skipsteps+3+(i*3))])) - 1
          elif fit kinetic == True:
              assoc_start = 
      int(np.sum(assaydict[sensor[0]][:(skipsteps+1+(i*2))]))
               baseline_start =\
      assoc_start - int(assaydict[sensor[0]][skipsteps+(i*2)])
               t0_baseline = dd[sensor[0]][baseline_start]
dissoc_start =\
      int(np.sum(assaydict[sensor[0]][:(skipsteps+2+(i*2))]))
                                                    30
```



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      # if offratefit_offset == True:
# print "offset = {0:.2e} +/- {1:2e}".format(coefoff[1],
#offrates_err[2])
            dissoc_end_y =\
      offrate(dd[sensor[0]][dissoc_start:dissoc_end], coefoff[0],
y0_dissoc, t0_dissoc, offset = 0)[-1]
    if fit_onrate == True:
                 if fit_kinetic == True:
      else:
                       coefon, covaron = curve_fit(lambda t, kon, maxresponse: \
      onrate(t, kon, coefoff[0], maxresponse, conc, t0_assoc),
      kon_stderrs = np.sqrt(np.diag(covaron))
                  koff_stderrs = np.sqrt(np.diag(covaroff))
print "concentration: ", conc
print "local k_on = {0:.2e}, stdev: {1:.2e}".format(coefon[0],
      kon_stderrs[0])
print "local ymax = {0:.2f}, stdev: {1:.2e}".format(coefon[1],
      kon_stderrs[1])
    print "local k_off = {0:.2e}, stdev: {1:.2e}".format(coefoff[0],
      koff_stderns[0])
    print "local K_d: {0:.2e}".format(coefoff[0]/coefon[0])
    print(''.center(80, '-'))
            cycleparm.append(dict([
            ('t0_baseline', t0_baseline),
      ('t0_dissoc', t0_dissoc),
('t0_assoc', t0_assoc),
('tend_dissoc', tend_dissoc),
('tend_assoc', tend_assoc),
       ('baseline_start', baseline_start),
      ( baseline_start , baseline_star
('assoc_start', assoc_start),
('baseline_end', baseline_end),
('baseline_y', baseline_y),
('y0_dissoc', y0_dissoc),
('dissoc_end_y', dissoc_end_y),
('assoc_end', assoc_end),
      ('offrate_covar', covaroff),
('onrate_parms', coefon),
('onrate_covar', covaron),
         'dissoc_end', dissoc_end),
      ('conc', conc),
('conc', conc),
('offratefit_offset', offratefit_offset),
('offset_dissoc', offset),
       ('fit_onrate', fit_onrate)]))
```

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<pre>if fit_kinetic == True:</pre>	errors)
<pre>sensors_180417_GGG_LaG9_vs_ybbR_TagGFP_4 = \     [('G1', 'H1', 25e-9, 50e-9, 100e-9, 200e-9)] analyze_experiment("180417_GGG-LaG9_vs_ybbR-TagGFP_4",     sensors_180417_GGG_LaG9_vs_ybbR_TagGFP_4, fit_kinetic=True,     fit_global=True, fit_stepwise=False, skipsteps=6, koff_initial=1     kon_initial=1e5)</pre>	.e-4,
<pre>SortaseA_TagGFP_LPETGG_3 = [('C1', 'D1', 25e-9, 50e-9, 100e-9, 200e- analyze_experiment("SortaseA_TagGFP-LPETGG_3", SortaseA_TagGFP_LPETG fit_kinetic=False, ymax_linked = False, fit_global=True, fit_y0= offratefit_offset=False, skipsteps=10, reference=True, fit_stepw</pre>	9)] GG3, F <b>alse</b> , rise= <b>False</b> )



# 3.5 Conformational Switching of a Protein Fold

In this article, we are able to show conformational transitions between different folded states of a cellulosomal protein by SMFS. We combine different measurement modes to acquire a very wide dynamic loading rate spectrum of the unfolding event of a Cohesin domain. This allows us to capture the features neccesary to distinguish the unfolding from at least two different states from the unfolding of one state via different unfolding pathways. Finally, we perform Monte Carlo simulations to illustrate that the observed behaviour is qualitatively consistent with a simple model of two folded states. Design and execution of this study was performed primarily by myself, as well as the drafting of the manuscript.

# Conformational Transitions of a Protein Fold Observed by Dynamic Single-Molecule Force Spectroscopy

Ellis Durner, Magnus S. Bauer, Markus A. Jobst, Wolfgang Ott, Lukas F. Milles, and Hermann E. Gaub

at the time of print, manuscript under review at PRL


### Abstract

When studying the mechanical properties of protein folds by Single-Molecule Force Spectroscopy, in some cases non-homogeneous behavior is observed. In these cases, the underlying process of forced unfolding from two distinct states must be discerned from two unfolding pathways of a single state. For these processes, theoretical predictions have been made based on transition state models of unfolding trajectories along a free energy landscape induced by a given force loading rate. Here, we experimentally test these predictions with atomic force microscopy-based singlemolecule force spectroscopy. Using a wide force-loading rate dynamic force spectrum and Monte Carlo simulations we experimentally validate the theoretical distinctions on a mechanostabilitystate-switching cohesin protein fold from Clostridia.

### INTRODUCTION

Single-Molecule Force Spectroscopy has evolved into a key tool for the investigation of proteins under mechanical stress[1, 2]. The most commonly employed approach to explore the energy landscape of a protein fold or the force-induced unbinding of a receptor-ligand system, is to mechanically load the system under well defined boundary conditions. A force probe is used to stretch single molecules in a known anchoring geometry, which can be controlled by site-specific surface immobilization strategies [3-5]. Most commonly, this is done either at constant retraction velocities  $\dot{z}$  or at constant force loading rates  $\dot{F}$  and until the protein unfolds or the complex dissociates[6]. The resulting dynamic force spectrum can often be interpreted with the Bell-Evans model [7], where the most probable rupture force of a single energy barrier increases linearly with the logarithm of the force loading rate. While this will in most cases be a gross oversimplification of the roughly  $3^N$ -dimensional energy landscape of a fold consisting of N amino acids, it describes the unbinding process from an energy minimum along a certain reaction coordinate well if there is only one dominant ratelimiting step. More recently, and with improved understanding of the molecular mechanisms, also non-linear dependencies have been predicted and observed [8-10]. However, all of these models are only consistent with a unimodal rupture force distribution at a given force loading rate. Recently, we have reported on the mechanostability of a family of cohesin-domains \* gaub@lmu.de

from Acetovibrio cellulolyticus (A.c.), which are part of the cellulose degrading machinery of these gram-positive bacteria [11–13]. We found, that while certain cohesins exhibited bimodal rupture force distributions, unimodality was recovered upon binding of a ligand dockerin. No dependence of this behavior on the loading rate was observed. We proposed, that a multitude of different folded states undergo a transition to a single folded state upon dockerin binding. In this study, we improve our understanding of this protein fold by resolving its unfolding behaviour under a very wide range of loading rates. We achieved this by observing cohesin 3 (Coh3) in both measurement modes of applying mechanical stress during a constant retraction velocity and under a constant force loading rate. We reveal that Coh3 adopts different folded states with distinct mechanical stabilities within the experimental timescale, a behavior which has been predicted[14], but was not experimentally observed for a single protein domain before.

### RESULTS AND DISCUSSION

To gain a more in-depth understanding of the dynamic force response of A.c.'s CipA cohesin domains, we chose Coh3 for further study. It has previously been characterized as being neither particularly strong nor weak in terms of force resilience when compared to the other cohesin domains of the same scaffold[12], and it had not exhibited any bimodality in its rupture force distribution, with or without a dockerin bound[13]. We performed a combined constant velocity and constant loading rate experiment, probing Coh3 at four different loading rates and four different constant retraction velocities, all with the same cantilever. This enabled us to obtain a dynamic force spectrum of Coh3 unfolding events that covers over five orders of magnitude in loading rate. Figure 1 illustrates the experimental configuration and shows example traces of Coh3 unfolding events, which match the contour length increment expected for full unfolding of Coh3.

To our surprise, the dynamic unfolding behaviour deviates markedly from the previously recorded log-linear behaviour[15]. At loading rates below  $\sim 1000 \,\mathrm{pN \, s^{-1}}$ , a downwards kink of the most probable rupture force at a given loading rate is observed (Figure 2A). To exclude the possibility for this to be an artefact from combining data of constant speed and constant loading rate measurements, we assembled the dynamic force spectrum of the CttA-XDoc:CohE (R.f) receptor-ligand system, which is used here as a specific pulling han-



dle [5, 16]. Despite the loading rates not being constant in a constant velocity experiment due to linker molecules that act like entropic springs (such as the elastin-like polypeptide linkers employed here [17]), fits of the Bell-Evans model to the most probable rupture forces were in good agreement between the two measurement modes (see suppl. Figure S1). By inspecting the unfolding force histograms of Coh3 for each measurement variation (Fig. 2B), we can gain a better understanding of the observed transition; while distributions at both the upper and the lower end of the loading rate spectrum appear unimodal, intermediate loading rates give rise to spectra that are distorted from what would be expected in the Bell-Evans model. Especially at  $100 \,\mathrm{nm \, s^{-1}}$ , a bimodal rupture force distribution emerges, which is similar to the mutant Coh1(T107S) as investigated by Verdorfer et al[13]. There, it was suggested that either different conformations or different unfolding pathways give rise

to the bimodality of these distributions, and it was shown that the fold could be stabilized upon binding to the ligand dockerin. Given that unfolding rates for a single barrier can be expressed as:

$$k(F) = k_{off}^0 \exp(\frac{F\Delta x}{k_B T}),\tag{1}$$

we can exclude that the bimodality can arise from one state with multiple unfolding pathways: As the unfolding probability  $p_i$  along a pathway *i* per time interval  $\Delta t$  is given as  $p_i = k_i(F) * \Delta t$ , it will increase monotonously with force. Consequently, the cumulative probability of n pathways

$$\sum_{n} p_i = \sum_{n} k_i(F) * \Delta t \tag{2}$$

will still result in a monotonously increasing rupture probability, and therefore a unimodal rupture force distribution. The data presented here not only require the existence of at least two different conformations for Coh3, which are distinct in their ability to withstand a stretching force. They also suggest that these states interconvert within the experimental timescale.

A simple model in line with the experimental data consists of two states with distinct single barriers to the unfolded state. Switching rates between these two states can be modeled as thermal crossings over a separating energy barrier with either a reaction coordinate affected by, or orthogonal to the applied force. The system is assumed to behave Markovian. To illustrate this, we performed Monte Carlo simulations for the simplest case: Figure 3 shows simulated data for a two-state system with fixed, non-force dependent rates for switching between the two folded states. The barriers'  $\Delta x$  and  $k_{off}$  for each state are taken from fits to either the two highest or the two lowest loading rates of the experimental dynamic force spectrum as indicated by the dashed lines in Figure 2A.

As no Monte Carlo fitting to the experimental data was performed due to the computational cost, rates for switching between the two folded states were chosen by hand. Hence, Figure 3 is merely demonstrates the general applicability of this simple two-state model[18], given the experimental data. A rigorous theoretical analysis of different possible cases of multi-state systems is provided in recent work by Pierse and Dudko[14]. In conclusion, we were able to experimentally observe a protein fold which needs to be described by theory beyond the standardly employed Bell Evans model. Our results are consistent with a twostate model as described in [14]. By combining constant loading rate and constant speed



FIG. 2. (A) Dynamic force spectrum of Coh3 unfolding events, assembled from constant loading rate experiments performed at loading rates of  $10 \,\mathrm{pN} \,\mathrm{s}^{-1}$ ,  $40 \,\mathrm{pN} \,\mathrm{s}^{-1}$ ,  $160 \,\mathrm{pN} \,\mathrm{s}^{-1}$ , and  $640 \,\mathrm{pN} \,\mathrm{s}^{-1}$  (mint green) and from constant velocity experiments performed at  $100 \,\mathrm{nm} \,\mathrm{s}^{-1}$ ,  $500 \,\mathrm{nm} \,\mathrm{s}^{-1}$ ,  $2500 \,\mathrm{nm} \,\mathrm{s}^{-1}$ , and  $12500 \,\mathrm{nm} \,\mathrm{s}^{-1}$  (blue). To illustrate the deviation from a simple log-linear behavior, the Bell Evans model (Equation 3) was fitted to either the two lowest loading rate values or the two highest loading rate values and the resulting fits were then extrapolated and plotted as dashed lines. (B) Rupture force histograms of the Coh3 unfolding events presented in (A). While distributions at both low loading rates and high pulling speeds appear unimodal, intermediate measurement variations show a bimodal behaviour, especially pronounced at a velocity of  $100 \,\mathrm{nm} \,\mathrm{s}^{-1}$ . Distributions were fitted with the Bell Evans model (suppl. Equation 4), describing the unfolding probability at a given force. It is evident that the model is a poor fit to the experimental data , as the width of the distribution should not change for different loading rates.



FIG. 3. Simulated unfolding events at different pulling speeds and loading rates for a complex that switches between two different folded states, independently of the applied force. Parameters  $\Delta x$  and  $k_{off}^0$  for the two states were chosen from the dynamic loading rate fits (see Figure 2A):  $\Delta x_A = 0.25 \text{ nm}$ ,  $k_{offA}^0 = 2.84 \times 10^{-5} \text{ s}^{-1}$ ,  $\Delta x_B = 0.22 \text{ nm}$ ,  $k_{offB}^0 = 1.1 \times 10^{-2} \text{ s}^{-1}$ . The switching rates between the states are  $k_{AB} = 4.8 \text{ s}^{-1}$  and  $k_{BA} = 19.2 \text{ s}^{-1}$ , chosen by hand to qualitatively illustrate the consistency of the model with the experimental data.

experimental modes, we were able to capture the non-monotonic trend in force distributions consistent with a two-state model. Thereby we could cover the wide range of loading rates[19] that is required to properly capture the dynamic force response of this protein domain, which had remained opaque in previous studies [12, 13].

The authors would like to thank Constantin Schöler for helpful discussions. We acknowledge funding from the German Research Foundation within the framework of the SFB1032.

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### Surface Functionalization

Glass surfaces and silicon nitride cantilevers (BioLever mini BL-AC40TS-C2, Olympus, Tokio, Japan) were silanized with (3-aminopropyl)-dimethyl-ethoxysilane (APDMES, ABCR GmbH, Karlsruhe, Germany). Then a sulfosuccinimidyl 4-(N-maleimidomethyl)cyclohexane-1-carboxylate) cross-linker (sulfo-SMCC, Thermo Fisher Scientific Inc.) was conjugated to the silanes amine groups in 50 mM 4-(2-hydroxyethyl)-1- piperazineethanesulfonic acid (HEPES) at pH 7.5. Next, the SMCC's maleimide groups were coupled to 50 µM Cys-

 $6\mathrm{xELP}\mbox{-}\mathrm{LPETGG}$  in case of the surfaces and to  $50\,\mu\mathrm{M}$  GGG-6xELP-Cys[7] in case of the cantilevers. This reaction was carried out in 50 mM sodium phosphate, 50 mM NaCl, 10 mM EDTA, with a pH of 7.2.

### Single-Molecule Force Spectroscopy

Measurements were performed in Ca-TBS buffer using a custom built AFM instrument (sample driven vertically by PI-313 piezo actuator and laterally by a 100x100 nm stage (P-734.2CL), Physik Instrumente, Germany) in conjunction with a MFP-3D AFM controller (Asylum Research, Santa Barbara, USA) and a custom control software. Upon approaching the sample surface with the cantilever tip, the complex between CohE and CttA-XDoc was formed and the cantilever was retracted from the surface at constant velocities of  $100 \,\mathrm{nm} \,\mathrm{s}^{-1}$ ,  $500 \,\mathrm{nm} \,\mathrm{s}^{-1}$ ,  $2500 \,\mathrm{nm} \,\mathrm{s}^{-1}$ , and  $12\,500 \,\mathrm{nm} \,\mathrm{s}^{-1}$  or at constant loading rates of  $10 \,\mathrm{pN} \,\mathrm{s}^{-1}$ ,  $40 \,\mathrm{pN} \,\mathrm{s}^{-1}$ ,  $160 \,\mathrm{pN} \,\mathrm{s}^{-1}$ , and  $640 \,\mathrm{pN} \,\mathrm{s}^{-1}$ . After each force-extension curve was acquired, the sample was moved laterally by  $100 \,\mathrm{nm}$  in order to not probe the same molecule twice on the surface side. Single-molecule interaction traces were identified by filtering the datasets using contour length analysis, and identifying only those traces in which a Coh3 unfolding events was observed[1, 8]. Loading-rates prior to Coh3 domain unfolding or complex dissociation were extracted by applying a linear fit to the last 3 nm in case of a constant velocity experiment. For constant loading rate experiments, a linear fit to the force over time was applied.

Dynamic force spectra were fitted to the Bell Evans model according to Equation 1.

$$F_{most\_probable}(\dot{F}) = \frac{k_B T}{\Delta x} \log(\frac{\dot{F}\Delta x}{k_{off}^0 k_B T})$$
(1)

Rupture force distributions were fitted to Equation 2, which gives the rupture force probability for a constant loading rate  $\dot{F}$ 

$$p(F) = \frac{k_{off}^0}{\dot{F}} \exp(\frac{\Delta xF}{k_B T} - k_{off}^0 k_B T \frac{\exp(\frac{\Delta xF}{k_B T}) - 1}{\Delta x \dot{F}})$$
(2)

### Monte Carlo simulations

We follow the approach described in Refs. ([9, 10]) by integrating  $k(F) = k_{off}^0 \exp(\frac{F\Delta x}{k_B T})$ over a timestep  $\Delta t$ . Here,  $k_{off}^0$  denotes the zero force off-rate and  $\Delta x$  is the distance to the

free energy barrier. As we are simulating a two state system, the probability to be in state A at t = 0 is given by  $p_{A,initial} = \frac{1}{(k_{AB}^0/k_{BA}^0)+1}$ . We allowed for the timestep to be variable in order to keep the probability of overcoming the barrier within a certain timestep below a value of 0.0001. This enabled us to retain good precision without resorting to extremely small timesteps throughout a simulation run. For the constant speed mode, we chose to model the entropic stretching response of a 240 nm polymer linker according to Livadaru et al.[11]:  $F(x) = \frac{k_B T}{cb(1-\frac{T}{L_c})}$ . The loading rate prior to rupture was determined using the derivative  $\dot{F}(x,v) = \frac{k_B T v}{cbL(1-\frac{T}{L_c})^2}$ . For the constant loading rate mode, the current force at a certain time t was simply given by  $F(t) = \dot{F}t$ .

### Protein Sequences

### ${\rm MGGG-HIS-G-Coh3(A.c)-CttA-XDoc(R.f)}$

MGGGHHHHHHGTGFTVSVDSVNGNVGEQIVIPVSFANIPANGISTADMTITYDSSKL EYVSGVPGSIVTNPDVNFGINKETDGKLKVLFLDYTMSTGYISTSGVFTKVTFKVLS SGGSTVGITGATFGDKNLGNVSATINAGSINGGVVPNTVTSAVKTQYVEIESVDGFY FNTEDKFDTAQIKKAVLHTVYNEGYTGDDGVAVVLREYESEPVDITAELTFGDAT PANTYKAVENKFDYEIPVYYNNATLKDAEGNDATVTVYIGLKGDTDLNNIVDGRD ATATLTYYAATSTDGKDATTVALSPSTLVGGNPESVYDDFSAFLSDVKVDAGKELT RFAKKAERLIDGRDASSILTFYTKSSVDQYKDMAANEPNKLWDIVTGDAEEE\*

### CohE(R.f)-HIS-HRV3C-LPETGG

MALTDRGMTYDLDPKDGSSAATKPVLEVTKKVFDTAADAAGQTVTVEFKVSGAE GKYATTGYHIYWDERLEVVATKTGAYAKKGAALEDSSLAKAENNGNGVFVASGA DDDFGADGVMWTVELKVPADAKAGDVYPIDVAYQWDPSKGDLFTDNKDSAQGK LMQAYFFTQGIKSSSNPSTDEYLVKANATYADGYIAIKAGEPELKLPRSRHHHHHH GSLEVLFQGPLPETGG\*





## 3.6 Immobilization Strategies for SMFS

When studying how protein domains react under force, careful design of the experimental arrangement is key. This article utilizes a combination of two orthogonal enzymatic coupling strategies, to covalently immobilize a protein of interest to a surface at one terminus, and to ligate a high force pulling handle to the other terminus. This allows to express the protein of interest with only two short tags, obviating the need for the classically employed polyprotein strategies. The high force receptorligand system CttA-XDoc:CohE enables specific and reliable probing of the anchored molecules. This article successfully demonstrates the compatibility of this immobilization and probing scheme with cell-free translation and transcription in microwells, directly on a sample slide, greatly reducing the amount of time and work spent on sample preparation. Design and execution of the study was performed by myself, as well as the drafting of the manuscript.

# Post-Translational Sortase-Mediated Attachment of High-Strength Force Spectroscopy Handles

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force–extension curve was recorded, the sample was moved laterally by 100 nm to probe a different molecule. For data analysis, force–distance curves were transformed into contour length space using a freely rotating chain model with quantum mechanical corrections for peptide backbone stretching<sup>21</sup> and then sorted by contour length increments.<sup>22</sup> Loading rates prior to domain unfolding or complex dissociation were extracted by applying a linear fit to the last 3 nm before the respective event and then used in fitting the rupture-force histograms with the Bell–Evans model.<sup>23</sup>

#### RESULTS AND DISCUSSION

To test for successful surface functionalization, we incubated surfaces that had been prepared as described in the Materials and Methods section with fluorescently labeled cohesin. Figure 2A confirms that if sortase is ommitted, no dockerin functionalization is achieved, whereas if sortase was present to perform the ligation reaction, binding of CoE-CBM-ybBR-CoA647 is observed. Figure 2B demonstrates successful ligation of GGG-dockerin to ybbR-Titin-LPETGG and illustrates the superior performance of the evolved sortase mutant d95/ P94R/D160N/D165A/K190E/K196T<sup>18</sup> in comparison with wild-type sortase A.

Typical single-molecule force–distance unfolding patterns for the sortase-incubated spots are shown in Figure 3A. They exhibit the unbinding pattern of CohE–Doc dissociation as characterized in previous publications,<sup>10</sup> where dissociation can occur with or without unfolding of the dockerin subdomain called x-module. The resulting force–distance curves were transformed into contour length space and then sorted by comparing the observed unfolding increments (3B). Only curves exhibiting the 56 nm increment corresponding to a full unfolding of the CBM-domain were classified to be the result of probing a CohE–Doc complex. Furthermore, the curves were sorted to exhibit nm omer than one increment corresponding to the unfolding of sfGFP 79 nm and no more than seven increments corresponding to Titin-Ig unfolding 28 nM. These increments result from the added free contour length of the peptide chain upon unfolding the folded protein domains and match the previously reported values  ${}^{i_{A},i_{A},2}_{A}$  For these traces, unfolding forces of the domains of interest were histogrammed with a bin width of 20 pN (Figure 3C).

Despite its narrow tip apex, each cantilever is typically functionalized with multiple cohesin-anchors; hence, multiple receptor-ligand complexes can form if dockerin-decorated surface is densely populated. Therefore, we went for a rather sparse surface functionalization which can be tuned by the incubation times of Sfp and sortase-catalyzed ligation reactions and/or the substrate concentrations. Alternatively, cantilevers with blunter tips could be used when more interactions are desired. The achieved surface densities were in a suitable range for SMFS, sparse enough to avoid multiple interactions tu dense enough to acquire good statistics. Probing attempts (1.24%) resulted in single molecule unfolding traces satisfying the outlined criteria. In total, 142 Titin-Ig and 92 sGGP single molecule traces were obtained within 11 h of measurement with a single cantilever (spring constant: 0.093 N m<sup>-1</sup>). If sortase had been omitted, no traces showing unfolding of CBM and one of the POI were recorded. For probing of in vitroexpressed Titin-Ig, 0.33% of attempts were successful, yielding 72 Titin-Ig unfoldings in 9 h of measurement, which was also probed with a single cantilever (spring constant: 0.097 N m<sup>-1</sup>).

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Figure 3C shows force histograms for unfolding events of sfGFP, the last of seven Titin-Ig domain to unfold and the complex dissociation itself. This was carried out for bulk-expressed and purified sfGFP and Titin-Ig, as well as for Titin-Ig expressed in the cell-free system. Complex dissociation events cluster into two populations that are characteristic of Doc/Coh unbinding.<sup>10</sup> The most probable forces at which the POI unfold are  $124\begin{pmatrix} + 24\\ -16 \end{pmatrix}$  pN for sfGFP,  $257\begin{pmatrix} + 36\\ -34 \end{pmatrix}$  pN for

the first, and  $365\left(+\begin{array}{c}33\\-\begin{array}{c}22\end{array}\right)$  pN for the last Titin-Ig domain to

unfold  $(271\left(\frac{+42}{28}\right)$  and  $404\left(\frac{+45}{-30}\right)$  pN for Titin-Ig expressed in the cell-free system), the asymmetrical full widths at half maximum of the distributions are given in brackets. The most probable forces were determined by fitting each histogram of unfolding forces with the Bell–Evans model.<sup>23</sup>

The differences between the most probable unfolding forces observed for the POI expressed in the cell-free system and the bulk-expressed proteins are within tolerance of errors resulting from cantilever calibration.25

This method can be easily applied to any recombinantly expressed protein by adding the terminal peptide tags necessary for covalent surface attachment and post-translational sortasemediated ligation. Owing to the terminal location of these tags, only nondigested and fully expressed proteins are probed. This is especially advantageous for cell-free expression systems, where the small quantity of expressed protein often makes the usually necessary affinity purification cumbersome.

### CONCLUSIONS

We developed a method that enables acquisition of SMFS datasets of specifically probed and covalently immobilized single molecules. By post-translationally modifying the POI with the high-force interactions of the Coh/Doc receptorligand system via sortase ligation, we can probe even resilient proteins such as Titin-Ig domains with high specificity and throughput, improving on the nonspecific polyprotein method and eliminating the requirement of expressing the POI as large fusion constructs with handle domains. The modular system of post-translational attachment of the mechanostable pulling handle allowed us to probe different proteins with the same cantilever. We also applied this approach to proteins expressed in cell-free systems without further purification while still selecting for only fully expressed proteins owing to the specificity provided by the high-affinity pulling handle.

### ASSOCIATED CONTENT

### Supporting Information

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsomega.7b00478. Amino acid sequences (PDF)

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The authors declare no competing financial interest.

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### 3.7 Linkers for SFMS

While linkers are beneficial to SMFS experiments in that they provide spacing of the molecules of interest from the surface they are anchored to, the most commonly employed PEG linkers can introduce artifacts themselves. When probing the elastic response of a polypeptide chain to force, polymer elasticity models allow us to access the total contour length of the stretched chain. If the stretched system is a mixed one, containing both polypeptide and PEG linkers, the differing elastic responses of each component, as well as the deviation of PEG from a purely entropic behaviour at higher forces, analysis of SMFS data becomes intricate. Here, we develop a genetically encoded ELP linker for force spectroscopy experiments. Not only does this ensure monodispersity in length, it is also ideal for force spectroscopy experiments on protein systems, as the ELP linkers are polypeptides themselves. The linkers allow for easy, covalent surface modification via a terminal Cysteine, and provide a Sortase recognition sequence at the other terminus, which is used to ligate a protein of interest under mild reaction conditions. We show that this linker system can be used to acquire high quality data, with a stretching response that more closely follows established polymer elasticity models than a mixed system. I contributed to this work by providing the pentamutant Sortase A variant used to covalently couple the ELP linkers to the molecules of interest. I also contributed to discussions about data analysis and polymer elasticity.

## Elastin-Like Polypeptide Linkers for Single Molecule Force Spectroscopy

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separated are pulled apart using the same receptor-modified cantilever. In such a configuration, Coh:Doc is used as a binding handle to successfully and continuously unfold target proteins for over 24 h of measurement time without significant loss of binding activity. Data sets of typically several tens of thousands of force-extension curves can easily be obtained using type III Coh:Doc, dramatically outperforming other mechanostable interactions (*e.g.*, biotin-avidin). The ability to measure with a single cantilever over several

The ability to measure with a single cantilever over several days allows interrogation of different types or variants of proteins immobilized on different positions of the same substrate (*i.e.*, protein microarrays) and to achieve statistical significance over the course of a single experiment. This leads to large data sets and requires the use of sophisticated algorithms to identify and extract specific single-molecule interactions among a large number of traces with poor signal, such as empty traces, multiple interactions in parallel, or nonspecific interactions. Independent of the size of the data sets though, elasticity models whether applied as part of elaborate algorithms or fitted manually to single curves have in the past been required to account for the different elastic contributions stemming from heterogeneous stretching behavior of mixed poly(ethylene glycol) (PEG)–protein polymer backbone.

Conformational Changes of PEG Linker Molecules Obscure Molecular Characteristics of Interest. When performing SMFS in an elevated force regime using PEG as linker molecules, additional challenges arise. A conformational transition of PEG occurs in a force range of up to *ca.* 300 pN, resulting in an approximately linear force–extension regime...<sup>10–18</sup> In aqueous solutions, PEG exhibits a trans–trans– gauche conformation. With rising force on the polymer, the occupancy of conformations is shifted to all-trans, effectively increasing the net polymer contour length. Analysis methods such as fitting standard elasticity models to the data or detecting contour length increments within said force range are therefore compromised and would, for a quantitative description, require improved heterogeneous elasticity models. PEG is a highly flexible polymer with a low persistence

PEG is a highly flexible polymer with a low persistence length, while peptide bonds have restricted degrees of freedom. These restrictions alter the stretching behavior and give rise to marked differences in comparison to PEG. Furthermore, the ratio of PEG linker length to unfolded protein backbone length is not constant over the course of an unfolding trace, which means fitting parameters must be optimized for different sections of the curve as more domains unfold. This issue becomes particularly significant and noticeable when probing protein unfolding and receptor—ligand unbinding in a high force regime and is also problematic when unfolding occurs across a broad range of forces. Benefits of ELP Linkers in SMFS. In this study we

Benefits of ELP Linkers in SMF5. In this study we investigate the feasibility of biological peptide polymers to circumvent this problem. We selected well-characterized elastin-like polypeptides (ELPs) as a suitable candidate for this purpose. The progression of cloning techniques of repetitive genes in recent years has set the stage for precisely defined protein polymers and opened up the ability to design, produce, and purify protein spacers of well-defined contour length and chemical composition for single-molecule experiments.<sup>19–22</sup> ELPs exhibit similar elasticity behavior as unfolded protein backbone and are completely monodisperse, a key advantage compared to synthetic polymers such as PEG. Monodisperse ELP linkers fused directly to a protein of interest allow for complete control of the lengths of a nanomechanical system from the surface up to the force transducer, which is not rue for the chemically synthesized PEG polymers with nonnegligible polydispersity. Since ELPs are expressed recombinantly in *Escherichia coli* (*E. coli*), their production is easily scaled up, resulting in lower costs compared to commercially available heterobifunctional PEGs. Furthermore, ELPs can be produced with N-/C-terminal protein ligation tags, which can be used for specific and bio-orthogonal surface chemistry in SMFS sample preparation. ELPs are synthetic biopolymers derived from tropoelastin

ELPs are synthetic biopolymers derived from tropoelastin domains. They are composed of a repetitive amino acid heptamer "Val-Pro-Gly-Xaa-Gly",<sup>33</sup> where Xaa is a guest residue that can be any amino acid apart from proline. The guest residue influences the hydrophobicity of the protein and impacts the lower critical solution temperature, the point at which the ELP undergoes a soluble-to-insoluble phase transition. At this environment-dependent cloud point, ELPs change their conformation and precipitate, resulting in clouding of the solution.

ELPs are intrinsically disordered proteins that do not fold into well-defined secondary and tertiary structures, but rather remain unfolded and flexible, a property that is ideally suited to their application as spacer/linker molecules for SMFS.<sup>24</sup> We hypothesized that ELPs would therefore be a suitable choice to achieve both surface passivation and site-specific immobilization in single-molecule nanomechanical experiments. The bulky yet flexible features of ELPs inhibit nonspecific protein binding to the surface, while enabling ligation of other proteins due to the high degree of accessibility of N- or C-terminally fused peptide tags. Post-translational protein ligation methods have made it possible to move from organic chemical conjugation methods toward enzyme-mediated covalent immobilization, for example utilizing sortase A or Sfp.<sup>14,25</sup> Both enzymes catalyze sequence- and site-specific reactions yielding uniform protein orientation at the surface.

ELPs have previously been the subject of atomic force microscopy (AFM) studies. For example, AFM was used to support theoretical predictions about the behavior of ELPs above and below their cloud point, as well as to study ELP elasticity.<sup>26–28</sup> This study was carried out entirely below the cloud point, so that intermolecular interactions between ELPs were negligible. In contrast to prior studies, we employ ELPs as spacer molecules with other protein domains attached. Our results show that ELPs provide several benefits over PEG linkers in SMFS attributable primarily to the features of having uniform elastic properties and monodisperse linkers.

This study offers an attractive substitute for established PEG systems using all-protein ELP linkers. The immobilization strategy provides precise control over the elastic properties of multicomponent protein mechanical systems linked between a glass surface and a force transducer. Our approach transfers advances in smart polymer research to SMFS experiments and describes the improvements achieved through this alternative surface anchoring strategy.

#### RESULTS AND DISCUSSION

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SMFS with Receptor-Ligand Polyproteins Employing Site-Specific Immobilization. Typically PEG linkers with an N-hydroxysuccinimide (NHS) group are linked to an aminosilanized surface. The other end of the PEG contains a reactive group for protein immobilization, which in most cases is a thiolreactive maleimide group. Figure 1A illustrates a Coh:Doc-

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orientation of the proteins at the surface. Such uniformly immobilized proteins lead to a well-defined propagation of the applied force through the molecular complex under investigation and to well-defined distributions of the unfolding/ rupture events in the force-extension curves. The use of Nand C-terminal tags for surface chemistry also ensured that only full-length (*i.e.*, fully translated) ELPs were measured in the experiment. AFM experiments performed with ELPs as linkers showed a

AFM experiments performed with ELPs as linkers showed a higher percentage of clearly identifiable single-molecule unfolding traces. We attribute this to the bulky character of the ELPs. They provide a less dense surface immobilization of the biomolecules of interest when compared to PEG-based immobilization. This behavior is advantageous since high surface density frequently causes multiple interactions between surface- and cantilever-bound molecules in SMFS experiments (Supplemental Figure S1). Multiple interactions are generated when more than one receptor–ligand interaction is formed in parallel. The complicated unfolding and unbinding traces that result from multiple bonds pulled in parallel are hardly interpretable and therefore discarded from the analysis (Supplemental Figure S2). Efficient passivation of glass surfaces against nonspecific adhesion of proteins requires a dense PEG surface layer, to prevent proteins from nonspecifically sticking to the glass surface. Approaches such as titrating functional (*i.e.*, maleimide end-groups) with nonfunctional (*i.e.*, CH<sub>3</sub> endgroups) PEG or changing the concentration of binding agents or proteins of interest can improve the process. In our experience, however, surface immobilization with ELP instead of PEG linkers leads to better passivation of the surface and a higher percentage of single-molecule traces without the need for any titration of functional and nonfunctional linkers.

Comparison of Dispersity between PEG and ELP Linkers. All unfolding traces were presorted by an automated analysis routine, selecting for single interactions that display two consecutive CBM unfolding events. Following the automated sorting, deletion of obviously erroneous curves (typically 10%) caused by, for example, baseline drift was performed manually.<sup>7,29</sup> PEG unfolding traces showed wildly varying initial extensions prior to the first CBM unfolding event. This is likely caused by the non-negligible polydispersity of PEG, as we did not observe multiple discrete populations with ELP experiments. The intrinsic monodispersity of ELP molecules is a clear advantage. Since they are produced recombinantly in *E. coli* with functional tags *in vivo*, only fulllength protein sequences have the necessary terminal peptide tags that allow for surface immobilization. Additionally, ELPs were purified with inverse transition cycling (ITC), a method developed for ELP purification based on their reversible precipitation behavior. Possibly shorter ELPs are removed during the process, since their cloud point is higher than for ELP<sub>120 nm</sub>. Although the polydispersity of chemically synthesized PEGs (mass distribution ~10–20 kDa) is sufficiently low for many applications, it leads to a noticeable impact in SMFS.

The influence of PEG polydispersity on the SMFS data is illustrated in Figure 3A, which shows SMFS traces recorded with both PEG and ELP linkers and also shows example traces of the shortest and largest extensions found in a typical type III Coh:Doc data set. Figure 3B shows a histogram of extension values at which the first CBM unfolding event occurred. For ELPs, the distribution shows one peak centered at an extension value that is expected based on the known ELP linker length. In the case of the PEG experiment, however, three distinct



Figure 3. Comparison of dispersity of PEG and ELP linkers. (A) Typical force-extension traces for PEG (purple) and ELPs (blue). In the PEG linker experiment, the unfolding events occur over a wider range of absolute extension values, whereas unfolding events with ELP linkers occur over a narrow range. (B) Histograms showing the distribution of extension values corresponding to the first CBM unfolding event in each curve (PEG: N = 219; ELP: N = 521). Due to the polydispersity of the PEG linkers, three discrete populations with different extensions are clearly visible, while for ELPs only one population is observed.

populations are observed. This can be understood by considering that at the level of single molecules a polydisperse distribution results in discrete peaks representing the corresponding lengths of the discrete polymeric linkers on the cantilever tip. We interpret the distributions as being caused by three different PEG molecules with different lengths attached to the tip. Although the discrete distributions could conceivably be caused by different positions of the molecule attachment points to the AFM cantilever tip, this effect should be the same for ELPs. Moreover, varying linker lengths also reflect in varying steepness of the force-extension trace peaks, which would not occur simply because of attachment geometry (Figure 3A, PEG traces). We exclusively observed monomodal distributions for ELPs; therefore an anchor position effect seems not to play a major role. This polydispersity is clearly disadvantageous, since multiple linker lengths render data analysis more difficult. Curves cannot simply be overlaid in force-distance space due to varying loading rates. Furthermore, for constant-speed SMFS experiments, loading rate populations in dynamic force spectra will be broadened due to the probabilistic nature of the thermally driven rupture events. We note that the PEG-modified surfaces are softer than ELP-

We note that the PEG-modified surfaces are softer than ELPmodified surfaces during indentation of the tip into the polymer brush, as determined by the curvature at the beginning of each trace. The firmer ELP-modified surfaces require a lower indentation force to reach a linear force-distance regime after the initial soft indentation. For calibrating the inverse optical lever sensitivity, this is advantageous since high indentation forces can damage the molecules attached to the tip through adsorption and denaturation processes.<sup>31</sup>

Uniform ELP Stretching Behavior Minimizes Artifacts. We hypothesized that by replacing synthetic PEG linkers with biological ELP linkers, and thereby having a single type of polymer backbone throughout the mechanical system, better defined elasticity properties for the recording of force curves would be achievable. The persistence lengths of ELP peptide backbones should be comparable to those of unfolded protein

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isomerization that also extends the contour length.27,34 However, the low number of prolines in the overall sequence (every fifth amino acid) in the ELP motif renders this effect much smaller compared to the conformational change of PEG and will be camouflaged by signal noise in typical experiments with proteins.

Figure 4C shows the transformation into contour length space using the qmWLC model. A kernel density estimate (KDE) was used (Gaussian kernel, bandwidth of  $2.5~\mathrm{nm}$ ) to generate smooth functions describing the contour length increments observed between unfolding or rupture events, which in this case included 2× CBM unfolding and Coh:Doc dissociation. In the case of PEG linkers, the KDE–contour length distribution shows several peaks. This is because of the failure of the qmWLC model to accurately describe the force response of the polymer. Determining the contour length increments between the peaks of the KDE proves problematic even for this relatively simple exemplary case of two large even for this relatively simple exemplary case of two large fingerprint unfolding events and a receptor ligand dissociation. Smaller unfolding steps or even folding intermediates, which appear as substeps, would be even harder to pinpoint with the PEG system. In the case of ELP-immobilized proteins, only three distinct peaks appear, with much more clearly identifiable contour length increments between the peaks.

#### CONCLUSION

PEG linkers have successfully been employed in numerous studies to anchor biomolecules of interest to surfaces for SMFS. In the low-force regime (below 100 pN) the extended WLC model describes their elastic properties with sufficient accuracy for the majority of applications. For elevated forces, however, the conformational transitions in the PEG backbone would necessitate further development of elasticity models for a convincing description.<sup>16</sup> Moreover, the inherent polydispersity of PEGs, together with their complex elasticity, complicates data analysis and reduces the amount of information that can be deduced from SMFS.

The ELP-based linkers, however, have proven in our studies to be significantly improved linker molecules for surface immobilization and passivation purposes in single-molecule force experiments. ELPs are monodisperse, are highly flexible, and readily allow for direct, site-specific tethering. We showed that these features lead to more accurate measurements of contour length increments in receptor-ligand polyprotein force spectroscopy experiments. A well-established elasticity model suffices for the data analysis.

Even at low forces, the PEG subunits already start to change their conformational state occupancy. At 50 pN, the probability for their elongated state is already above 10%.<sup>16</sup> Therefore, the findings we present here are also relevant for investigations at lower forces or in systems that should be analyzed over a large range of forces. PEG linkers may still deliver satisfying results, as long as data in similar force ranges can be compared. In some cases, elasticity parameters such as the Kuhn length or persistence length can heuristically compensate for effects not explicitly described by the model. As soon as different force explicitly described by the model. As soon as different force ranges of multiple domains need to be compared, though, the varying proportions of elongated (all-trans) versus non-elongated (trans-trans-gauche) PEG subunits cannot simply be accounted for by the elasticity parameter, and therefore measured contour length increments get distorted. Different biochemical approaches like those described here are thus accounted to gain and the transfer to the second to the second necessary to gain meaningful insights. These scenarios include,

for example, shielded unfolding events or small substeps, where

The ELPs investigated here represent only one formulation of the vast variety of smart polymer linkers that could be utilized in SMFS experiments. Further studies are required to evaluate other nonstructured, non-proline-containing protein linkers to determine their suitability for SMFS studies, since the amino acid side chain composition may affect the persistence length  $^{36,37}$  or give rise to nonentropic behavior. Biotechnological characteristics, *i.e.*, recombinant production yields and ease of purification, are as important as the biophysical requirements, which renders the easily produced ÉLPs particularly attractive. Other smart polymers should be similarly accessible to perform as suitable alternatives. The reported approach can be applied to enhance SMFS studies with purified proteins on functionalized surfaces as shown here or alternatively to modify cantilevers for chemical recognition imaging and force spectroscopy on artificial membranes or cell surfaces. It can easily be adopted by standard molecular biology equipped laboratories to streamline the procedure and improve data quality for resolving smaller unfolding features with high accuracy. Studies on smart polymers as tethers for SMFS experiments might also help to develop environmentally responsive surfaces, which bear potential for exciting applications in the nanobiosciences.

the force cannot drop sufficiently in between stretching events.

### MATERIALS AND METHODS

All reagents were at least of analytical purity grade and were purchased from Sigma-Aldrich (St. Louis, MO, USA) or Carl Roth GmbH (Karlsruhe, Germany). All buffers were filtered through a 0.2  $\mu$ m poly(ether sulfone) membrane filter (Nalgene, Rochester, NY, USA)

poly(ether suitone) memorane inter (Naigene, Kochester, NY, USA) prior to use. The PH of all buffers was adjusted at room temperature. A 300 amino acid long ELP was the basis for the AFM linker constructs used in this study, and the underlying cloning and protein purification procedure of the ELP is described in detail elsewhere.<sup>19</sup> The ELP sequence was [(VPGVG)<sub>5</sub>-(VPGAG)<sub>2</sub>- (VPGGG)<sub>3</sub>]<sub>6</sub> and is sence was  $[(VPGVG)_{5}-(VPGAG)_{2}-(VPGGG)_{3}]_{6}$  and is The ELP sequence was  $[(VPGVG)_{5^{-}}(VPGAG)_{2^{-}}(VPGGG)_{5^{-}}(VPGAG)_{2^{-}})$ referred to as  $ELP_{120 \text{ nm}}$ . Standard molecular biology laboratories capable of

producing recombinant proteins are equally capable of expressing ELPs, since both rely on the same principles, reagents, and instrumentation. With our plasmids provided at Addgene, cloning can even be avoided and production of ELP linkers for protein immobilization can be

production of ELP linkers for protein immobilization can be performed right away. **Cloning**. A detailed description of the cloning procedure of the constructs can be found in the Supporting Information (Figures SS-S11). ELP sequences used in this study, along with 40 nm length variants and binding handles, are deposited at Addgene and available upon request (Addgene accession numbers: 90472: Cys-ELP,<sub>120 nm</sub>-Cys, 91572: GGG-ELP,<sub>120 nm</sub>-Cys, 91572: GGG-ELP,<sub>120 nm</sub>-Cys, 91572: GGG-ELP,<sub>120 nm</sub>-Cys, 91598: GGG-HIS-CBM-Xmod-DocIII). **Transformation of Cells**. A 2  $\mu$ L amount of Gibson assembly or ligation reaction transformed DH5 $\alpha$  cells (Life Technologies GmbH, Frankfurt, Germany; 30 min on ice, 1 min at 42 °C, 1 h at 37 °C in SOC medium) was used. The cells were plated on 50  $\mu g/mL$  kanamycin-containing LB agar and incubated overnight at 37 °C. Clones were analyzed with Colony PCR, and clones with amplicons of appropriate lengths were sent to sequencing.

Clones were analyzed with Colony PCR, and clones with amplicons of appropriate lengths were sent to sequencing. **Protein Expression**. Chemically competent *E. coli NiCo21(DE3)* (New England Biolabs, Ipswich, MA, USA) were transformed with 50 ng of plasmid DNA for the expression of all constructs used in this study. Transformed cells were incubated in autoinduction ZYM-5052 media (for ELP containing constructs supplemented with 5 mg/mL proline, valine, and 10 mg/mL glycine; 100 µg/mL kanamycin) for 24 h (6 h at 37 °C, 18 h at 25 °C).<sup>56</sup> Expression cultures were harvested via centrifugation (6500g, 15 min, 4 °C), the supernatant was discarded, and the pellets were stored at -80 °C until further lysis.

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Throughout the whole purification process, for ELPs containing a cysteine, 1 mM tris(2-carboxyethyl)phosphine (TCEP, Thermo Fisher Scientific Inc., Waltham, MA, USA) or 1 mM of dithiothreitol (DTT) was added to the respective buffers. Cell pellets with proteins containing no HIS-tag were solubilized in 50 mM Tris-HCl pH 7.5 (supplemented with cOmplete, EDTA-free protease inhibitor cocktail, Sigma-Aldrich, St. Louis, MO, USA), and all other pellets in lysis buffer (50 mM Tris, pH 8.0, 50 mM NaCl, 10% (w/v) glycerol, 0.1% (v/v) Triton X-100, 5 mM MgCl<sub>2</sub>, DNase I 10 µg/mL, lysozyme 100

(v/v) Triton X-100, 5 mM MgCl<sub>22</sub> DNa8e 1 10 µg/mL, 19505/mE 200 µg/mL). Cys-ELP<sub>120 nm</sub>-LPETGG and GGG-ELP<sub>120 nm</sub>-Cys were purified with the ITC method.<sup>39</sup> After resolubilization, the cells were lysed by sonication (Bandein Sonoplus GM 70, tip: Bandein Sonoplus MS 73, Berlin, Germany; 40% power, 30% cycle, 2 × 10 min). The cells were kept on ice during the sonication procedure. The soluble fraction was separated from the insoluble cell debris by centrifugation (15000g, 4 °C, 1 h). In a first heating step (60 °C, 30 min) of the supernatant, most of the *E. coli* host proteins precipitated. The fraction of the collapsed ELPs was resolubilized by cooling the suspension for 2 h to 4 nost of the D contrast process processing the supersisting of D to A of C on a reaction tube roller. The insoluble host proteins were pelleted by centrifugation (15000g, A °C, 30 min). Further purification steps were necessary to increase the purity of the ELP solution. This was done by repeated thermoprecipitation of the ELP followed by redissolution. The ELP solution was clouded by adding 1 M acetate buffer (final

concentration 50 mM, pH 2.5) and 2 M NaCl. A heating step (60  $^{\circ}$ C, 30 min) ensured all ELPs were collapsed. A hot centrifugation (3220g, 40  $^{\circ}$ C, 75 min) was necessary to separate the high-salt, low-pH 40 C, 75 min ) was necessary to separate the high-sait, low-pri-solution from the ELP pellet, which was resolubilized in 50 mM Tris-HCl (pH 7.0) after discarding the supernatant. The solution was incubated for 2 h at 4 °C to resolubilize all ELPs completely. A cold centrifugation step (3220g, 4 °C, 60 min) isolated the remaining insoluble fraction of the suspension. After decanting the supernatant, the salt concentration was increased and pH lowered, to precipitate the ELPs again. This guide user remarched these times or attended if the ELPs again. This cycle was repeated three times or extended if the purity of the solution was not high enough. The constructs CohIII-CBM-HIS-LPETGG and GGG-HIS-CBM-

Xmod-DocIII were expressed and lysed as described above. After the first centrifugation, the supernatant was, however, filtered (0.45  $\mu m)$  and applied to a HisTrap FF (GE Healthcare Europe GmbH, Freiburg, and applied to a Hisi rap FP (GE Healthcare Europe GmbH, Freiburg, Germany). Unspecifically bound proteins on the column were removed by washing five column volumes (25 mM Tris-HCl pH 7.8, 500 mM NaCl, 20 mM imidazole, Tween 20 0.25% (v(v)), 10% (v(v) glycerol). Finally, the desired HIS-tag containing protein was eluted (25 mM Tris-HCl pH 7.8, 500 mM NaCl, 300 mM imidazole, Tween 20 0.25% (v(v)), 10% (v(v) glycerol). For long-term storage the protein solutions of the different constructs were concentrated (Amicon Ultra-15 centrifugal filter write 10W WWCO March VCoA. Demoty Comparison and reduced

constructs were concentrated (Amicon Ultra-15 centrilugal filter units 10K MWCO, MercK KGaA, Darmstadt, Germany) and reduced with 5 mM TCEP overnight (at 4 °C) for constructs that contained a cysteine. The buffer of the reduced ELP solution was exchanged (Zeba spin desaling columns 7K, Thermo Fisher Scientific Inc.) to 50 mM sodium phosphate, 50 mM NaCI, 10 mM EDTA, with a pH of 7.2, and 100 ( $\mu$  ( $\mu$ ) downed and  $\mu$  day for any in the initial bitments in small dimension

solium phosphate, 50 mM NaCl, 10 mM EDTA, with a pH of 7.2, and 10% (v/v) glycerol and flash frozen in liquid nitrogen in small aliquots to be stored at  $-80^{\circ}$ C. All other proteins were exchanged with 25 mM Tris-HCl, 75 mM NaCl, and 5 mM CaCl<sub>2</sub> with a pH of 7.2 and supplemented with a final glycerol concentration of 20% (v/v). No loss of functionality of the ELPs (cross-linking and passivation capability) could be detected, when stored buffered or lyophilized in small aliquots at  $-80^{\circ}$ C, over the duration of more than one year. SDS-PACE (Any kD Mini-PROTEAN stain-free gels, Bio-Rad Laboratories GmbH, Hercules, CA, USA) was employed to detect any impurities. Since ELPs could not be stained with the stain-free technology, an Alexa Fluor 647-C<sub>2</sub>-maleimide dye (Thermo Fisher Scientific Inc.) was incubated for 1 h at room temperature with the ELP solution. An appropriately diluted protein solution was mixed with 5x loading buffer (250 mM Tris-HCl, pH 8.0, 7.5% (w/v) SDS, 25% (v/v) glycerol, 0.25 mg/mL bromophenol blue, 12.5% (v/v) 2-mercaptoethanol) and heated for 5 min at 95 °C.

ELP concentration was photometrically determined at 205 nm (Ultrospec 3100 Pro, Amersham Biosciences, Amersham, England, and TrayCell, Hellma GmbH & Co. KG, Müllheim, Germany). For all other constructs an absorption measurement at 280 nm led to the concentration (NanoDrop UV-vis spectrophotometer, Thermo Fisher Scientific Inc.). The extinction coefficient was determined theoretically for ELPs at 205 nm<sup>40</sup> and 280 nm<sup>41</sup> for all other fusion

AFM Sample Preparation. Force spectroscopy samples, measure-ments, and data analysis were prepared and performed according to previously published protocols.<sup>10,35</sup> Silicon nitride cantilevers (Biolever min BL ACOTS C2 Object Constructions of the construction of th mini, BL-AC40TS-C2, Olympus Corporation, Tokyo, Japan; nominal spring constant: 100 pN/nm; 25 kHz resonance frequency in water) spring constant: 100 pN/nm; 25 kHz resonance frequency in water) were used as force probes. Surface chemistry for cantilevers was similar to that for coverslips (Menzel Gläser, Braunschweig, Germany; diameter 24 mm). Surfaces were amino-silanized with 3-(aminopropyl)dimethylethoxysilane (APDMES, ABCR GmbH, Karls-ruhe, Germany). a: Maleinimidohexanoic-av-NHS PEG (NHS-PEG-Mal, Rapp Polymere, Tübingen, Germany; PEG-MW: 15 kDa) was used as a linker for the sortase peptides (GGGGG-C and C-LPETGG, Centic Biotec, Heidelberg, Germany) in PEG-linked experiments. The cysteine-containing ELPs were linked to the surface with a sulfosuccinnindyl 4.(-Maleinimidohevil)cyclohexane-1-carboxylate) cross-linker (10 mM) was dissolved in 50 mM 4-(2-hydroxyethyl)-1-piperazineethanesulfonic acid (HEPES) pH 7.5.

piperazineethanesulfonic acid (HEPES) pH 7.5. Sortase-catalyzed coupling of the fingerprint molecules (GGG-CBM-Xmod-DocIII and CohIII-CBM-LPETGG) was done in 25 mM

CBM-Amod-DocIII and CohIII-CBM-LPETGG) was done in 25 mM Tris-HCl, pH 7.2, 5 mM CaCl<sub>2</sub> and 75 mM NaCl at 22 °C for 2 h. Typically, 50  $\mu$ M ELP or sortase peptide was coupled with 25  $\mu$ M fingerprint molecule and 2  $\mu$ M sortase enzyme. In between both of the cross-linking steps (PEG, SMCC, or ELP, peptide reaction) surfaces were rinsed with water and dried with nitrogen. After immobilization of the fingerprint molecules, surfaces were rinsed in measurement buffer (25 mM Tris-HCl, pH 7.2, 5 mM CaCl. 75 mM NaCl). CaCl<sub>2</sub>, 75 mM NaCl). The reaction of the different surface chemistry was done spatially separated by using silicone masks (CultureWell reusable gaskets, Grace Bio-Laboratories, Bend, OR, USA). The mask was applied after silanization and removed under buffer after the last

Immobilization step. AFM-SMFS Measurements. Data were taken on custom-built instruments (MFP-3D AFM controller, Oxford Instruments Asylum Research, Inc., Santa Barbara, CA, USA; piezo nanopositioners: Physik Instrumente GmbH & Co. KG, Karlsruhe, Germany, or Attocube Systems AG, Munich, Germany). Instrument control coherage area custom written in Iror. Pro 6.3.

Instrument control software was custom written in Igor Pro 6.3 (Wavemetrics Inc., Portland, OR, USA). Piezo position was controlled with a closed-loop feedback system running internally on the AFM controller field-programmable gate array. A typical AFM measurement took about 12 h and was done fully automated and at room temperature Retraction velocity for constant-speed force spectroscopy measurements was 0.8  $\mu$ m/s. Cantilever spring constants were calibrated after completing all measurements on different spots on the surface using the same cantilever. This was done by utilizing the thermal method applying the equipartition theorem to the one dimensionally oscillating lever.<sup>31,42</sup>

Force–Extension Data Analysis. Obtained data were analyzed with custom-written software in Python (Python Software Foundation, Python Language Reference, version 2.7, available at http://www. python.org), utilizing the libraries NumPy, SciPy, and Matplotlib.

Raw voltage data traces were transformed into force distance traces with their respective calibration values after determining the zero force value with the baseline position. A correction of the force-dependent cantilever tip z-position was carried out. Force distance traces were filtered for traces showing two CBM unfoldings and a subsequent type III cohesin–dockerin dissociation, without preceding Xmodule unfolding.<sup>7</sup> This screening was carried out by detecting maximum-to-maximum distances of kernel density estimate (Gaussian kernel, bandwidth 1 nm) peaks in contour length space in each single trace, after applying thresholds for force, distance, and number of peaks. For

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sorting data sets, transformation of force distance data into contour length space was done with a manually fixed persistence length of 0.4 nm, to measure distances of energy barrier positions.<sup>29,43</sup> Sorting was nm, to measure distances of energy barrier positions.<sup>2705</sup> Sorting was done allowing generous errors to the expected increments to account for the conformational stretching of the spacer molecules. Fits to the force–extension data with the WLC model had the following parameters additionally to the values mentioned in the figure captions, if not stated otherwise: initial guess for persistence length: 0.4 nm; fit precision:  $1 \times 10^{-7}$ . For assessment of transformation quality, the inverse worm-like-chain model was applied for transformation of force distance traces into the contour langth space in a force window of 10. distance traces into the contour length space in a force window of 10 to 125 pN and with a persistence length previously fitted to each peak separately: The global mean value of each data set for each peak was used. Final alignments of the whole data sets were assembled by crosscorrelation

correlation. Master Curve Assembly. The master curves were assembled by cross-correlation of each force-distance trace of a presorted data set with all previous curves in contour length space, starting with a random curve. Each curve was shifted on its x axis to fit the maximum correlation value and added to the set assembly in contour length space. Subsequently, a second run was performed, cross-correlating each curve with the previously assembled set, to facilitate an equal correlation template for every curve, independent of its occurrence. Finally, the most probable shift was calculated by a KDE and subtracted from each curve to get representative absolute distances with respect to the origin. Distance and correlation value thresholds were applied to filter out less probable PEG populations and otherwise badly fitting data. In a final step, all overlaid raw data points in force-distance space were binned on the x axis into nanometer-sized slices, distance space were binned on the x axis into nanometer-sized slices, and their densities on the y axis were estimated by a KDE for each slice. Near the rupture events, where the kernel density estimates cannot unambiguously identify maxima of the data slices, the value was set to zero. Therefore, after each rupture, a small "gap" is visible, which was not included in data points used for fitting. Their most probable value and the corresponding full width at half-maxima then assembled the master curve. Although by this procedure representative absolute rupture forces for the domains are not necessarily reproduced to the highest accuracy, the most probable and most representative pathway of the elastic behavior in between peaks is resembled well. **qmWLC model.** For WLC fits and transformations into contour

length space, a recently improved approximation, solved for the extension, was used,<sup>32</sup> adding correction terms for quantum mechanical backbone stretching.<sup>33</sup> With the abbreviations

$$f = FL_p/kT$$

$$b = \exp\left(\sqrt[4]{\frac{900}{f}}\right) \tag{2}$$

WLC fits were done with the model formula

$$x = L_{\rm corr} \left( \frac{4}{3} - \frac{4}{3\sqrt{f+1}} - \frac{10b}{\sqrt{f(b-1)^2}} + \frac{f^{1.62}}{3.55 + 3.8f^{2.2}} \right)$$
(3)

With the quantum mechanical correction,

$$L_{\rm corr} = \frac{L_{c,0}}{2y_2} (\sqrt{4Fy_2 + y_1^2 - y_1 + 2y_2})$$

(4) where  $y_1$  and  $y_2$  are the *ab initio* parameters from the original publication. Transformations were performed with the model contour length:

$$L_{\rm c} = \frac{x}{\frac{4}{3} - \frac{4}{3\sqrt{f+1}} - \frac{10b}{\sqrt{f(b-1)^2}} + \frac{f^{1.62}}{3.55 + 3.8f^{2.2}}}$$

With the reverse quantum mechanical correction for zero force contour length,

$$=\frac{L_c}{\frac{1}{2y_2}(\sqrt{y_1^2+4y_2F}+2y_2-y_1)}$$

with x being the extension,  $L_c$  the model contour length, F the force,  $L_p$  the persistence length, k Boltzmann's constant, T the temperature,  $L_p$  the persistence length, k boundarins constant, r the emperature,  $y_1$  and  $y_2$  the quantum mechanical correction parameters,  $L_{our}$  the qm-corrected contour length, and  $L_{c,0}$  the reverse qm-corrected contour length at zero force. As a nonlinear fitting algorithm, a Levenberg–Marquardt least-squares minimization method was applied.  $y_1$  and 1

(6)

### ASSOCIATED CONTENT

L<sub>c,0</sub>

**Supporting Information** The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsnano.7b02694. Further details on experimental methods, supplementary results, and sequence information (PDF)

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### Author Contributions

W.O.: experiment design, sample preparation, measurements, data analysis, writing of manuscript; M.A.J.: experiment design, data analysis, writing of manuscript; M.S.B.: data analysis; E.D.: sample preparation; L.F.M.: data analysis; M.A.N.: experiment design, writing of manuscript; H.E.G.: experiment design, writing of manuscript.

### Notes

(1)

(4)

(5)

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The authors declare no competing financial interest.

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For this analysis, only forces in a range from 10 to 125 pN were taken into account, to minimize the effects of conformational stretching. The elastic properties of the first stretching event of a data trace are dominated by the linker molecules. As more protein domains unfold, the peptide backbone of the unfolded domains contributes increasingly to the overall elastic response. Contour length transformations of force distance data were performed with the mean fitted persistence lengths of each peak, as shown in **Supplemental Figure S3**, **Panel A** (0.35, 0.44, and 0.49 nm for ELP data peaks; 0.20, 0.25, and 0.27 nm for PEG data peaks), to account for

















Name	Sequence (5'-3')
Construction of Cys-ELP <sub>120 nm</sub> -LPETGG	· · ·
FW N-Cys Bsal	GACTCTCTGGAATTCATCGCTTCTAAACTGGC TGGTCTCCTGCGTGCCGGGAGAAGGAG
REV Bsal ybbR	CCCGGCACAGCCAGTTTAGAAGCGATGAATTC CAGAGAGTCGGTCTCACATATGTATATC
Construction of TEV-GGG-ELP60 nm-LPETG	G
QuikChange Primer ybbR to TEV-GGG	GACACCAGGGACTCCTTCTCCCGGCACACCG CCCCCTCCCTGGAAGTACAGGTTTTCCATATG TATATCTCCTTC
Construction of TEV-GGG-ELP <sub>60 nm</sub> -ybbR	
QuikChange Primer LPETGG to ybbR	GACACCAGGGACTCCTTCTCCCGGCACACCG CCCCCTCCCTGGAAGTACAGGTTTTCCATATG TATATCTCCTTC
Construction of TEV-GGG-ELP <sub>120 nm</sub> -ybbR	
FW backbone Bsal	GAAAACCTGTACTTCCAGGGAGGGGGGGTCTC GGGGTGTGCCGGGAGAAGGAG
REV backbone Bsal	ATATATGGTCTCGACCGCCCCCTCCCTGGAAG TACAGGTTTTC
FW insert TEV-GGG Bsal	CCAGGGAGGGGGGGTCTCGCGGTGTGCCGGG AGAAGGAG
REV insert Bsal	TCGAGTTAAGCCAGTTTAGAAGCGATGAATTC CAGAGAGTCGGTCTCCACCCTCACCCGG
Construction of GGG-ELP <sub>120 nm</sub> -ybbR	·
FW ELP GGG	GGGGGCGGTGTGCCGGGAG
REV Bsal TEV	GGCACACCGCCCCCCCCCGGAAGTACAGGT TTTCGGTCTCACATATGTATATCTCCTTC

Г

Construction of GGG-ELP <sub>120 nm</sub> -Cys	
FW backbone Cys	GCCAGTTTAGAAGCGATGAATTCCAGAGAGTC GGTCTCCACCTTCACCC
REV ybbR Bsal	TGCTAACTCGAGTAAGATCCGGCTGCTAACAA AGCCC
Construction of GT-CohllI-CBM-HIS-	LPETGG
FW backbone	TAACTCGAGTAAGATCCGGCTGC
REV CBM LPETGG	GCCGCCGGTTTCCGGCAGCGGACCCTGGAAC AGAAC
Construction of CohIII-CBM-HIS-LPE	TGG
FW Cohlli	GCGCTCACAGACAGAGGAATG
REV backbone without GT	CATATGTATATCTCCTTCTTAAAGTTAA
Construction of TEV-GGG-HIS-CBM-	XDocIII
FW backbone	CTCGAGTAAGATCCGGCTGC
REV backbone	ACCGGGTTCTTTACCCC
FW insert	GTATGGGGTAAAGAACCCGGTGGCAGTGTAG TACCATC
REV insert	CGGATCTTACTCGAGTTATTCTTCTTCAGCATC GCCTG
Construction of GGG-HIS-CBM-XDo	
FW CBM	ATGGCCAATACACCGGTATCA
REV backbone	TCCGTGGTGGTGGTGGTGGTGACCGCCCCC

repeats (5) <sub>x</sub>	ε <sub>205</sub> [1/M cm] <sup>7</sup>	Molecular weight [Da] <sup>8</sup>	Isoelectric point	Amino acids in ELP repeats (total) <sup>8</sup>	Total Length [nm] <sup>9</sup> (.365 nm per aa)
Cys-ELP <sub>120 nm</sub> - LPETGG	851370	24763.08	3.20	300 (307)	112.06
GGG-ELP <sub>120 nm</sub> - Cys	843030	24379.63	3.23	300 (304)	110.96
	GVGVPGV( GVGVPG\	GVPGVGVP(	GVGVPGA	GVPGAGVPGGG	VPGGGVPGEGVF
VPGVGVPGVGV GVPGAGVPGAG AGVPGGGVPGG GGGVPGEGC <u>Cys-ELP<sub>120 nm</sub>-LP</u> Cysteine ELP Sortase C-Tag	PGAGVPG SVPGGGVF GVPGEG' <u>ETGG</u>	GAGVPGGG GAGVPGGGVPGE VPGEGVPG	PGAGVPG VPGGGVF GVPGEG\ VGVPGVC	GGVPGGGVPGE GEGVPGEGVPG /PGVGVPGVGVPG GVPGVGVPGVGVP	GVPGEGVPGVGV VGVPGVGVPGVG GVGVPGVGVPGA PGAGVPGAGV

MGGG-HIS-CBM-Xmod-Dockerin III Sortase N-Tag His<sub>6</sub>-Tag CBM **Xmod** Dockerin III MGGGHHHHHHGMANTPVSGNLKVEFYNSNPSDTTNSINPQFKVTNTGSSAIDLSKLTLRYYYT VDGQKDQTFWSDHAAIIGSNGSYNGITSNVKGTFVKMSSSTNNADTYLEISFTGGTLE PGAHVQIQGRFAKNDWSNYTQSNDYSFKSASQFVEWDQVTAYLNGVLVWGKEPG<mark>G</mark>S PATTKPPATTIPPSDDPNAVVPNTVTSAVKTQYVEIESVDGFYFNTEDKFDTA QIKKAVLHTVYNEGYTGDDGVAVVLREYESEPVDITAELTFGDATPANTYKAVENKFDYE IPVYYNNATLKDAEGNDATVTVYIGLKGDTDLNNIVDGRDATATLTYYAATSTDGKDATT VALSPSTLVGGNPESVYDDFSAFLSDVKVDAGKELTRFAKKAERLIDGRDASSILTFYTK SSVDQYKDMAANEPNKLWDIVTGDAEEE Cohesin III-CBM-HIS-LPETGG Cohesin III СВМ His<sub>6</sub>-Tag Sortase C-Tag MALTDRGMTYDLDPKDGSSAATKPVLEVTKKVFDTAADAAGQTVTVEFKVSGAEGKYATT GYHIYWDERLEVVATKTGAYAKKGAALEDSSLAKAENNGNGVFVASGADDDFGADGVMWTV ELKVPADAKAGDVYPIDVAYQWDPSKGDLFTDNKDSAQGKLMQAYFFTQGIKSSSNPSTDEYL VKANATYADGYIAIKAGEP<mark>GSVVPSTQPVTTPPATTKPPATTIPPSDDPNA</mark>MANTPVSGNLKVE FYNSNPSDTTNSINPQFKVTNTGSSAIDLSKLTLRYYYTVDGQKDQTFWSDHAAIIGSNGSYNGI TSNVKGTFVKMSSSTNNADTYLEISFTGGTLEPGAHVQIQGRFAKND WSNYTQSNDYSFKSASQFVEWDQVTAYLNGVLVWGKEPGELKLPRSR<mark>HHHHHH</mark>GSLEVLFQ GPLPETGG



the calculated polymer sizes, in black the data the calculation is based on. Number of subunits always round to the next integer.					
Molecular Weight [Da]	Number of Subunits	Length [nm]			
513.3	4	2.5			
601.6	6	3.2			
689.71	8	3.9			
865.92	12	5.3			
1394.55	24	9.5			
1000	15	6.4			
5000	106	38.3			
10000	220	78.1			
15000	333	118.0			



## 3.8 Force Activation of Enzymes

This publication investigates the mechano-activation of focal adhesion kinase. In a combined in-vitro and in-silico approach, we show that mechanoactivation can occur at forces below the unfolding forces of FAK. Our findings are consistend with, and suggest that force-induced conformational changes of FAK may induce focal adhesion signaling. My contribution to this work was the construction of a high-speed AFM for SMFS experiments. This enabled experiments at retraction velocities of 12 800 nm s<sup>-1</sup>, which resulted in loading rates beneficial in elevating unfolding events above the noise level.

# Structural and mechanistic insights into mechanoactivation of focal adhesion kinase

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on an enzyme was previously described for the mammalian titin kinase, the related twitchin kinase in nematode (8, 9), and the smooth muscle myosin light-chain kinase (10), which are located





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integrin signaling layer in FAs in close proximity to the plasma membrane (12). Accordingly, the FERM domain contains a stretch of basic residues forming a basic patch that interacts with phosphatidylinositol 4,5-bisphosphate (PIP2) (13, 14), which is clustered in the cell membrane at FA sites (15, 16). Activation of FAK can be initiated by an orchestrated process starting with PIP2 binding to the basic patch, resulting in exposure of a linker region containing the autophosphorylation site (14). After autophosphorylation, this site becomes a docking site for the Src kinase, which in turn, phosphorylates tyrosine residues in the activation loop of the FAK kinase. This results in full activation of FAK. The last step of Src phosphorylation is strongly promoted by prior separation of FERM and kinase domains (11). Here, we test the hypothesis that force is a key stimulus driving domain separation for FAK activation. With the N-terminal fAT domain tethered via paxillin and structural FA components to the actin cytoskeleton, FAK is likely exposed to stretching forces in FAs. Force-induced domain opening and activation were indeed already proposed and supported by molecular dynamics (MD) simulations, showing that forces applied to FAK result in rupture of autoinhibitory interactions and exposure of Src phosphorylation sites in the FAK activation loop (17). Furthermore and consistent with this hypothesis, FAK is known to undergo domain opening when localized to FAs in cells (13). Additionally, mutational disruption of autoinhibitory interactions and resulting FAK opening has experimentally been shown to strongly promote phosphorylation by Src (11).

Initiational display of a sexperimentally been shown to strongly promote phosphorylation by Src (11). In this study, we aim to obtain insight into the force response of FAK on a single-molecule level by means of atomic force microscopy (AFM)-based force spectroscopy in conjunction with force probe molecular dynamics (FPMD) simulations. We are able to measure highly reproducible force profiles recorded during FAK stretching and map-detected force events to structural features in FAK. By adapting experimental conditions, we are able to identify interface rupture of FERM and kinase domains as a discrete force peak. This is verified by control force profiles of FAK mutants lacking autoinhibitory FERM-kinase interactions and comparison with MD simulations. We conclude that the inhibitory interface ruptures at forces of around 25 pN for pulling speeds of 12,800 nm/s (or at smaller forces for slower pulling)—significantly before any domain unfolding—and that functionally important regions in FAK start to unfold at an extension of about 50 nm. Hence, we demonstrate that force activation of FAK can occur at physiological forces in FAs and that FAK maintains an active state when exposed to extensions expected to occur in FAs.

#### Results

Relative Mechanical Stabilities of FERM and Kinase Are ATP Dependent. In this study, we record force-distance traces for FAK containing FERM and kinase domain (residues 1–686; referred to hereafter as FK-FAK) (Fig. IC) attached via short PEG linkers (425.39 Da) and pulling speeds of 800 nm/s. The curves show force peaks between 20 and 50 pN, with reliably recurring unfolding features and a final Strep-Tag II (SII) rupture at around 70 pN (Fig. 2). Force application guides the protein through a sequence of conformational states obeying a strict hierarchy as suggested by the heatmap, and the most probable unfolding curve in black (Fig. 2). All detected peaks above the FWHM of the noise level are highlighted with a colored triangle (the process is described in *Data Analysis*) (Fig. 2). These unfolding peaks indicate transitions between these conformational states and mark the unfolding of stable structural modules within the protein structure leading to a release of previously hidden contour lengths. To determine contour length increments of the polypeptide chain unfolded for each peak, we fitted the data before each rupture event with the worm-like chain model (18, 19).

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Fig. 2. Assignment of force peaks to structural modules of FERM and kinase domain (residues 1–686; FK-FAK) using AFM unfolding patterns. (A) The heatmap obtained from an overlay of 224 curves shows the typical unfoldheatmap obtained from an overlay of 224 curves shows the typical unfold-ing pattern of KF-RK in the absence of ATP, revealing low-force unfolding below 50 pN and final SII rupture at around 70 pN. In *Upper*, the increments between identified peaks are depicted, allowing assignment of the rupture events to structural modules in the crystal structure. (8) Depiction of the KK-RK unfolding pattern by an overlay of 115 curves in the presence of 3 FK+AK unfolding pattern by an overlay of 115 curves in the presence of 3 mM ATB. Both plots are created from one dataset recorded with the same cantilever, and therefore, they are directly comparable in absolute force. Although the unfolding pattern in B looks different from the one in A, the increments stay conserved. Comparison of the two conditions shows that the increment of f1 is swapping its position with k1 and k2 on ATP binding. This can be attributed to stabilization of the kinase domain on addition of ATP, The balance of the balance of the theory, changing the force hierarchy between FERM and kinase. (A and B) The black lines show the most probable unfolding patterns with all detected peaks above the FWHM of the noise level (described in Data Analysis) highlighted with colored triangles according to their assigned domain.

We determine contour length increments of 20 nm (k1), 68 nm (k2), and 48 nm (f1) for FK-FAK (Fig. 24 and Table 1). Additionally, FK-FAK was stretched in presence of 3 mM ATP to probe the effect of ATP binding to the kinase domain. Although the presence of ATP substantially changes the unfold-ing pattern of FK-FAK (as observed in Fig. 2), the identified

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increments remain remarkably conserved [48 nm (f1<sup>ATP</sup>), 11 nm (k1<sup>ATP</sup>), and 68 nm (k2<sup>ATP</sup>)] (Fig. 2*B* and Table 1). This suggests that the unfolding hierarchy changes on ATP addition as increment f1 moves from the back to the beginning of the curve (Fig. 2). The most probable rupture force of peak k1 (Fig. 2) significantly increases on ATP binding, corroborating its associ-ation with the kinase domain (20, 21). This is probably due to a stabilization of the part of the kinase domain involved in ATP binding, thereby yielding higher forces for peak k1 and conse-quently, changing the force hierarchy between FERM and kinase domain. The kinase domain maintains the internal sequence of unfolding (k1 then k2a) but unfolds after the FERM domain in the presence of ATP. Their sums of all contour lengths (88 and 79 nm, respectively) agree with expected values for the FAK kinase domain [91 nm for residues 422–686 (19, 22), subtracting 5 nm to account for the initial distance between residues 422 and 686 in the folded domains according to the crystal structure]; the cover a signed 1 and k2 to the kinase domain. The increment f1 is assigned to FERM unfolding, and a contour length of 48 nm reproduces an expected length of 50 nm for folded FERM from residues 216–362 plus an additional 12 residues for the inserted ybbR-tag, subtracting an initial folded distance of 6 nm. We, therefore, conclude that stabilizing the kinase domain by ATP binding results in a reversed unfolding sequence, with the kinase binding results in a reversed unfolding sequence, with the kinase unfolding first in absence of ATP but as the last event in the pres-ence of ATP. All determined contour length increments are in good agreement with defined structural features in the crystal structure of FK-FAK in Fig. 1B (11).

To further validate our assignment of unfolding increments, we used single-domain constructs with only the isolated FERM or kinase domains. The curves were analyzed as previously described (Figs. 3 and 4). The data confirm our assignment of the structural modules, with determined contour length incre-ments matching well with the ones assigned in FK-FAK plots (Table 1). In agreement with our assignment in FK-FAK, the (Table 1). In agreement with our assignment in FK-FAK, the isolated kinase likewise unfolds in two main modules (k1 and k2a). Furthermore, we identify additional unfolding intermedi-ates that were previously hidden in the noise. For the FERM domain, we now resolve three peaks with contour lengths 7 nm (f1a), 25 nm (f1b), and 14 nm (f1c). For the kinase domain, k2a is followed by another peak k2b, but it is too short to determine a proper contour length increment. Close inspec-tion reveals that these intermediate states (k2b, f1b, f1c) are in fact also present in FK-FAK (Fig. 2). They can be detected in the beginning of the unfolding curve where the total free length is still short enough leading to hisher loading rates (as length is still short enough, leading to higher loading rates (as property of the worm-like chain model) and consequently, to higher forces, thereby allowing us to resolve more subtle peaks. These subtle peaks are, however, too small (and their extension is too short) to determine their contour length accurately. determined contour length increments are summarized All in Table 1.

# Table 1. All measured contour length increments of the various FAK constructs used in this study

Protein segment	f1	k1	k2
FK-FAK —ATP	48	20	68
FK-FAK +ATP	48	11	68
FERM - ATP	46		
f1a/f1b/f1c	7/25/14		
FERM +ATP	46		
f1a/f1b/f1c	7/21/18		
Kinase – ATP		15	66
Kinase +ATP		14	66

determined as most probable values from a KDE.



hidden within prominent peaks. The interface opening between kinase and FERM likely represents such a small peak that is expected to release a contour length of only 20 nm. Assuming physiological activation under stress in vivo, the opening has to occur at relatively small forces and before any other unfolding, since it is shielding the remaining structure. However, since in this region of the force profiles, we do not detect an additional distinct force peak in FK-FAK profiles with a pulling speed of 800 nm/s, we conclude that domain separation may occur anywhere before the first unfolding event (i.e., below 40 pN in this loading rate regime) but likely is hidden in the measurement noise.

Aiming to resolve domain separation, we used higher pulling speeds of 12,800 nm/s to increase force responses and used long PEG (5,000 Da) for FAK attachment to facilitate analysis early in the force profile. This indeed enabled us to resolve an additional force peak at the beginning of the unfolding pattern (yellow triangles in Fig. 5.4), where we expect domain separation to occur. To verify that this peak originates from the domain opening, we probed a mutant of FK-FAK (FK-FAKmut) with point mutations at the domain interface (Y180 A, M183 A). These mutations cause FERM and kinase domains to be permanently dissociated (14). As shown in Fig. 5.8, this mutant lacks the low-force rupture event, which we identified as the interface opening in the FK-FAK wild-type profile



Fig. 4. Verification of the assignments of structural modules by probing single-kinase (residues 411–686) domains. The heatmaps show single-kinase unfolding events without (*Upper*) and in the presence of 3 mM ATP (*Lower*). The first peak k1 increases in force by addition of ATP suggesting a binding event of ATP to the kinase domain. The same peak detected in detail here (*k*2b) can also be found in the kinase domain in Fig. 2A. Due to longer linkers at the end of the unfolding curve and therefore, lower loading rates leading to lower forces, they cannot be detected in Fig. 2B.

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Fig. 5. Resolving the FERM-kinase domain interface rupture by comparing FK-FAK with a permanently open FK-FAKmut—recorded with long PEG (5,000-Da) linkers at a pulling speed of 12,800 nm/s. (A) The heatmap shows the same construct as measured in Fig. 2 but with improved force resolution in the investigated region. This way, an additional subtle peak in the beginning of the curve can be observed highlighted by yellow triangles. This peak is not detected in the permanently open FK-FAKmut construct measured in *B*. We conclude that this rupture event corresponds to the interface opening between FERM and kinase domain. With measurements in the presence of 3 mM ATP (*Bottom*), the same behavior is observed, hence showing no association of the rupture event with a protein domain and validating its assignment to interface opening. *SI Appendix*, Fig. S2 provides additional contour length

(yellow triangles and Int label in Fig. 5.4), but still retains all other unfolding features. As expected, domain dissociation in the FK-FAK wild type precedes any unfolding events. In conclusion, these data indicate that forced domain dissociation and thereby, conformational activation happens before any other subdomain unfolding at a force around 25 pN at pulling speeds of 12,800 nm/s.

MD Simulations Confirm Conformational Activation of FAK. Fig. 6 summarizes our unfolding simulations (detailed unfolding data are given in *SI Appendix*, Figs. S3–S11). Consistent with previously simulated results (17), the first event is invariably dissociation of the FERM-kinase interface. After this, there is considerable heterogeneity in our data, with most of the simulations featuring both domains being in the process of unfolding. However, it should be considered that pulling speeds in simulations are significantly faster than in AFM experiments, and Fig. 6B shows a clear trend of slower simulations featuring the experimentally observed hierarchy: the FERM domain unfolds before the kinase. We, therefore, restrict the following analysis to these cases (five at 0.1 m/s and two at 0.33 m/s).

The first force-induced conformational change is the loss of the FERM-kinase interface, giving 10 nm of extension (Fig. 6 and *SI Appendix*, Table S1). This is followed by another 12-nm extension due to the linker losing contact with the F1 lobe in the FERM domain (Fig. 6B). In agreement with simulations, domain separation of FAK was also experimentally observed as the first event (Fig. 5), and FERM-linker separation, which simulations show to require less force (*SI Appendix*, Fig. S11) and to occur shortly after domain separation, is not detected in experiments.

The FERM domain unfolding was observed experimentally to occur in three steps: one at 7–8 nm, a second one around 21–25 nm, and a third one around 14–19 nm. In our MD simulations, the FERM domain unfolding happens in two stages: first, lobe F3 unfolding is observed, amounting to 30 nm, and

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second, lobe F2 stretching is observed, corresponding to an increase of 9 nm. Since the F2 stretching never happens before F3 unfolding in our trajectories, we assign the first two events observed experimentally to F3 unfolding. The F2 stretching contributes to an increase in end-to-end distance of the protein by about 9 nm, which in experiments, amounts to around 13 nm considering the additional loop contained in the handle of the experimental FK-FAK construct. Kinase domain unfolding leads to FAK deactivation independent of the detailed sequence of events. We, therefore, discuss kinase unfolding events of the kinase subdomains observed in MD pulling simulations vs. AFM in *SI Appendix*. Importantly, the same unfolding sequences of in simulations performed on the subdomains only, further validating the experiments on individual domains. *SI Appendix* has details.

#### Discussion

Here, we report a detailed mechanical characterization of FAK by using an AFM setup to apply stretching forces on single FAK molecules and record force-extension profiles with high sensitivity. Combining our measurements with structural information of the FERM-kinase region of FAK (11) and FPMD simulations has allowed us to assign measured force peaks to unfolding of defined structural factures in FAK. Importantly, increasing the force loading rate by applying high pulling speeds has enabled the detection of a low-force event corresponding to rupture of the autoinhibitory FERM-kinase interaction. We show that domain separation occurs at a low-force regime around 25 pN for pulling speeds of 12,800 nm/s. The fact that forces required for domain unfolding supports the hypothesis that tensile forces in FAs applied to N- and C-terminal regions in FAK can trigger activation via domain separation. Mutational dissociation of FERM and kinase domains has previously been shown to activate FAK (11).



Fig. 6. Domain dissociation and linker detachment precede domain unfolding. (A) Rupture events across the 30 simulations are depicted. The green pentagons show the moment of activation (i.e., FERM-kinase dissociation; measured by a sudden increase in V215-K630 distance), and the three different types of bars show when unfolding of these elements begins and ends (monitored by end-to-end distance changes); (B) Preferred unfolding mechanism as a function of pulling velocity: FERM unfolds first before the kinase (blue), kinase unfolds first before the kinase (red), or a "mixed" unfolding pathway (white). The bars show proportions of the 10 simulations performed at that pulling velocity. Snapshots 1–4 show the process of conformational activation and linker release observed in MD simulations as illustrated by a sample trajectory at 0.1 m/s. In the native state (snapshot 1), the FERM domain is in contact with the kinase. After an initial kinase C-terminal helix unwinding (SI Appendix, Fig. 56), these two domains lose contact (snapshot 2). The interdomain linker gradually detaches and elongates, leading to a release of Tyr397 (snapshot 3). After an additional extension of 12 nm, the linker is fully stretched (snapshot 4). The approximate positions in time of these snapshots are shown in A.

We propose that tension forces in FAK are built up between the membrane-bound N-terminal FERM domain and the Cterminal FAT domain engaged via paxillin to vinculin and the actin cytoskeleton (Fig. 7). This scenario was also suggested by our previous MD simulations of the first conformational activation step (17), where FERM and kinase domains detached from one another via these attachment sites. These previous simulations explicitly included a PIP2-containing bilayer and indicated that, for membranes enriched in PIP2, such as is the case at FA sites (15, 16), the membrane–FERM interaction is mechanically significantly more stable than the autoinhibitory FERM-kinase interaction. Force applied at the FAK C terminus in the opposite direction of the membrane resulted in FERM-kinase separation independent of the pulling angle or loading rate (17). Consistent with AFM measurements and the MD simulations also suggest that domain separation occurs before domain unfolding. In the previous simulations as well as in the full unfolding simlations shown here, the force required for separating FERM

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and kinase domains is 150 pN; however, the much higher pulling speed in simulations (6 mm/s or higher) is known to cause overestimation of rupture forces. Our combined AFM experiments and MD simulations suggest that the FAK domain organization protects against unfolding of functional domains of FAK: that is, the kinase domain and the PIP2 binding site in the F2 lobe of FERM. First, FERM-kinase domain dissociation yielding around 10-nm extension is followed by around 12-nm-long extension (*SI Appendix*, Table S1, Linker-F1) due to the stretching of the interdomain linker including the Tyr397 phosphorylation site. This would indicate that the length of the linker further protects the kinase domain from forced unfolding after the dissociation of the FERM and kinase domains. Second, F3 unfolding involves an increase in extension of about 30 nm (*SI Appendix*, Table S1, F3 unfold) and happens when the F2 lobe is still capable of binding PIP2. All in all, this would mean that FAK has a "safety margin<sup>2</sup> up to a total length of 50 nm (10-nm domain dissociation, 12 nm of linker stretching plus at least 28 nm for F3 unfolding) where it is still catalytically active and also, capable of binding PIP2. Consistent with our model where force to the C termi-In 2: Consistent with our model where force to the C termines of FAK is applied via paxillin and vinculin (Fig. 7), vinculin is found to transition from a signaling layer close to the membrane in FAs (which also contains FAK) to a force transduction layer closer to actin (23). Both layers have an approximate thickness of 30 nm measured vertical to the membrane; therefore, the average movement of vinculin approximates 30 nm toward the force transduction layer. This suggests that the 50-nm safety margin appears sufficient to protect the average engaged FAK molecules from force-induced deactivation. The fraction of FAK molecules norm force-induced deactivation. The fraction of PAK molecules exceeding this margin would expect to unfold their F2 lobe, hence losing contact with the membrane but retain-ing an active kinase. At an average of 30-nm extension, our data suggest that FAK molecules, after they are extended, no longer experience significant stretching forces other than what is required to keep FAK in an extended conformation. Likely, forces generated in FAs are mainly carried by structural components, such as talin and vinculin. In contrast for FAK, force seems to act as an activation catalyst by operating a digital distance switch, which is "on" when forces stretch FAK into an open conformation

In contrast to force activation of FA signaling, force-induced changes on structural FA components, such as talin or vinculin, have been characterized in detail. Mean forces experienced by talin and vinculin in FAs in cells have been determined to be in the range of 7-10 pN and around 2.5 pN, respectively (24, 25). However, force estimates from bulk measurements can be deceiving, since the load could be carried by only a fraction of molecules, whereas many others might not be engaged. Indeed, for integrins, average forces have initially been estimated at 1–2 pN (10); however, DNA-based tensions sensors that are irreversibly ruptured above a threshold force indicate peak forces of up to 40 pN for single integrin molecules (26). In our experiments, FERM-kinase separation occurs at around 25 pN at a pulling speed of 12,800 nm/s. Considering that average cellular force application in FAs is likely slower or even constant over certain time periods, it is highly likely that engaged FA molecules build up sufficient force (at a maximum of 40 pN for integrins) to separate FERM and kinase domains in FAK. Intriguingly, we find that unfolding of the FAK kinase domain in the more physiological ATP-loaded state occurs at around 50 pN; therefore, cellular forces in FAs seem well suited to allow FERM-kinase separation but not kinase unfolding, which are both prerequisites for FAK activation.

Force-induced separation of FERM and kinase domains will expose both the autophosphorylation site in the linker and the Src phosphorylation sites in the kinase activation loop. Mechanical extension of the Tyr397-containing linker might enhance Tyr397 exposure and autophosphorylation. This would likely

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Fig. 7. Model of force-induced FAK activation. FAK is recruited into FAs via C-terminal FAT interactions with paxillin and talin. The N-terminal FERM domain docks via PIP2 to the lipid membrane to promote a primed FAK state where Tyr397 in the linker between FERM and kinase is autophosphorylated. Forces generated via the actin cytoskeleton pull FAK's C terminus away from the membrane, resulting in kinase release from the FERM domain and membrane. Src is recruited to autophosphorylated FAK and phosphorylates the exposed FAK activation loop to trigger full FAK activity.

only be the case in transphosphorylation mode, as force would act against folding back of Tyr397 into the active site of the same FAK molecule. However, we showed previously that membrane binding and resulting FAK oligomerization are sufficient to promote highly efficient FAK autophosphorylation also in the absence of force (14). This study also indicated significant membrane-induced conformational changes that expose the autophosphorylation site but apparently not the kinase active site, since contrary to domain separation, membrane binding did not catalytically activate FAK (14). Together, the two studies, therefore, support a model where initial membrane binding promotes a primed state of FAK by exposing the linker for efficient autophosphorylation, but subsequent buildup of tensile forces in FAK exposes the active site for efficient phosphorylation of the activation loop by Src (Fig. 7). It is the latter event that promotes full catalytic activity of FAK. It was recently shown that the kinase domain of FAK also contributes to binding to PIP2 membranes (27); therefore, force might be responsible for removing the kinase from both the FERM domain and the membrane. MD simulations indeed support such a scenario and found that the pulling angle can dictate which occurs first (17). In conclusion, our mechanical analysis of FAK supports a

In conclusion, our mechanical analysis of FAK supports a model where physiological stretching forces in FAs can cause conformational changes in FAK, promoting its catalytic activation and thereby, triggering of FA signals. Multiple cellular studies have previously shown that FAK is activated in response to various mechanical stimuli (4–6), and our analysis on single FAK molecules demonstrates the feasibility of direct force activation of FAK. Force-induced activation of FA signals is highly relevant in disease. In tumors, stiffening of the stroma that allows increased force generation triggers strong adhesion signals that promote tumor invasion (28). Understanding the direct relation between tumor stiffness, force-induced adhesion signaling, and tumor invasion can, therefore, provide the basis for the development of specific agents targeting this mechanism.

#### **Materials and Methods**

AFM Setup for Characterizing FAK. To mimic physiological FAK stretching as occurring in FAs and to identify force-induced structural changes in FAK under stress, we developed an AFM-based single-molecule force spectroscopy assay. This allows for the detection of subtle force-induced events

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for FAK with high sensitivity during its guided stretching. We engineered FAK proteins to harbor affinity pulling handles for attachment to sample surface and cantilever. We introduced a ybbR-tag (29) for covalent link-age to the glass surface and an SII (30) for reversible tethering to an AFM cantilever tip functionalized with a monovalent Strep-Tactin (monoST) (26). Both attachments are formed via heterobifunctional PEG linkers (Fig. 1C has a schematic of FAK attachment to AFM). Before the experiment, FAK proteins are covalently immobilized to the glass surface. Typically, several thousand single-molecule AFM manusurements are then performed, and force extension profiles are recorded by repeatedly approaching and retracting the functionalized and overlaid to generate heatmaps highlighting recurring features in the plots. These recurring unfolding events were identified by creating a most probable unfolding curve as described in *Data Analysis*.

In a previous study, the autoregulatory region of FAK was defined as FERM interacting with the kinase domain (14). Initial AFM experiments were conducted with full-length FAK (residues 1–1052 in Fig. 18, Bottom) and a construct containing only the FERM and kinase domains (residues 1–686, FK-FAK in Fig. 18, Bottom), both equipped with affinity tags at their N and C termini. These experiments indicated that the FAT domain does not contribute to the force profile of the autoregulatory region (*SI Appendix*, Fig. S1); hence, subsequent experiments were only performed with FK-FAK. To mimic the physiological force path through the molecule, we introduced the N-terminal tag close to the lipid binding site, which in vivo attaches to the cell membrane. To prevent perturbation in protein folding, we inserted the 11-residue ybbR-tag into an unstructured loop immediately before the K216AKTLRK PIP2 binding site in the FERM domain. We confirm that these insertion mutants retain basal activity of vild-type FAK and that the FERM domain still maintains the ability to autoinhibit the catalytic activity of FAK (SI Appendix, Fig. S13). Previous MD simulations confirmed the PIP2-FERM linkage to be significantly more mechanically robust than the FERM-kinase interaction at relevant PIP2 concentrations (17), rationalizing the choice of a covalent handle to mimic of the FERM-membrane interaction.

FAK Expression. Chicken FAK constructs were engineered to contain the 11-aa ybbR-tag after V215, just before the K216-AKTLRK basic patch sequence (29), and the 8-aa SII tag (30) at the C terminus. FAK constructs containing full-length, FERM and kinase, or kinase-only regions were expressed by transient transfection of HEK293GnT1 cells using polyethyleneimine as a transfection agent (31). FERM-only constructs were expressed in Escherichia coli BL21 (DE3) as in ref. 32. All proteins were expressed with an N-terminal GxHis tag. Initial purification was performed by Ni-chelate affinity purification (GE Healthcare) followed by protease cleavage to remove the 6xHis tag. Proteins containing an SII tag were further purified by

Strep-Tactin (GE Healthcare) affinity and size exclusion (Superdex 200; GE Healthcare) chromatography. Proteins without SII tag were further purified by anion exchange (Source 15Q; GE Healthcare) and size exclusion chromatography.

Sample Preparation for Surface and Cantilevers. The preparation of the experiment includes specifically immobilizing (29, 33) the FAK construct on the glass surface and functionalizing the cantilever with an monoST. This ensures a well-defined pulling geometry for minimizing multiple interactions. All FAK constructs harbored a ybRot-tag for covalent immobilization on a glass surface and an SII for binding to the monoST-functionalized cantilever (Fig. 1C). All measured constructs were derived from chicken FAK and expressed in HEK cells (compare with FAK Expression).

Both cartilevers and glass surfaces were passivated by short 425.39-Da [SM(PEG)2; PEGylated SMCC cross-linker; Thermo Scientific Pierce] or long 5,000-Da (molecular mass 5,000; Rapp Polymere) heterobifunctional PEG spacers to avoid unspecific interactions between the cartilever and the glass surface. The PEG spacers offer an N-hydroxy succinimide group on one side for attachment to the amino silanized surface of the cantilever. The other end provides a Maleimide (Mal) group for attachment of the thiol group found in the Cysteine of the monoST.

found in the Cysteine of the monosi. For silanization, the cantilevers were first oxidized in a UV ozone cleaner (UVOH 150 LAB; FHR Anlagenbau GmbH) and subsequently silanized for 2 min in (3-aminopropyldimethylethoxysilane (ABCR; 50% (col/vol) in MilliQ and afterward, dried at 80 °C for 30 min. After that, the cantilevers were incubated in a solution of 25 mM heterobifunctional PEG spacer and 50 mM Hepes for 30 min (for short PEG first solved in half DMSO and then filled to 50 mM Hepes). Finally, the monosT was bound to the cantilevers for 1 h at room temperature followed by a washing step in 1 × PBS. The functionalized cantilevers were stored in measurement buffer (40 mM Hepes, pH 7.4, 10 mM MgCl2, 200 mM NaCl, 1 mM DTT) until use.

The preparation of the glass surfaces is in a lot of steps similar to the functionalization of the cantilevers as seen in Fig. 1C. The glass surfaces are amino silanized followed by a passivation with PEG linkers. The Mal of PEG offers a binding site for the thiol group of CoA. Via an Sfp-catalyzed reaction, the CoA can bind the ybbR-tag harbored by the FAK protein construct. This way, the protein gets attached and tethered in an uncompromisingly specific way. Before silanization, the glass surfaces have to be cleaned by sonification and the protein gets attached and tethered in the surface surfaces have to be cleaned by sonification and the surface surfaces have to be cleaned by sonif

Before silanization, the glass surfaces have to be cleaned by sonification is 50% (vol/vol) IPA in MilliQ for 15 min. For oxidation, the glass surfaces are soaked for 30 min in a solution of 50% (vol/vol) hydrogen peroxide (30%) and sulfuric acid. Afterward, they have to be thoroughly washed in MilliQ and then blown dry in an N2 stream. Then, the glass surfaces get silanized by incubating them in ABCR [1.8% (vol/vol) in ethanol]. Thereafter, they were washed again in IPA and MilliQ and then dried at 80 °C for 40 min. Then, the PEG is applied as described for the cantilevers. Subsequent to rinsing, the surfaces were incubated in 20 mM CoA (Calbiochem) dissolved in coupling buffer (sodium phosphate, pH 7.2) to react with Mal. After washing the glass surfaces, 8  $\mu$ L of the FAK construct (20  $\mu$ M) was mixed with 1  $\mu$ L 5fp-synthase (132  $\mu$ M) and 1  $\mu$ L of 10× reaction buffer (100 mM Tris, pH 7.5, 100 mM MgCl<sub>2</sub>); then, it was pipetted on the surfaces and incubated for 2 h at room temperature. Finally, the surfaces were rinsed thoroughly in measurement buffer (40 mM Hepes, pH 7.4, 10 mM MgCl<sub>2</sub>, 200 mM NaCl, 1 mM DTL).

Force Spectroscopy Experiments. The AFM measurements were conducted on an Asylum research controller (Asylum Research) providing analog-todigital converter and digital-to-analog converter channels as well as a digital signal processor board for setting up feedback loops. The controller operated either a custom-built AFM head (34) or an xyz-movable piezo-driven sample stage. Data were recorded automatically by cycling through the following steps: (*i*) approach of the functionalized AFM tip to the surface to allow coupling to the SII of surface-immobilized FAK; (*ii*) retraction of the AFM cantilever with nanometer precision at a constant speed and simultaneous recording of the mechanical force response with pN precision; and (*iii*) after monoST:SII spearation, the piezo stage of the AFM is moved to probe a new spot on the sample surface in the next cycle. This process was operated by using an IgorPro6 (Wavemetrics) program controlling the z piezo in the AFM head (or sample stage) and the *xy* piezos. The surface is sampled in steps of 100-nm distance in a snail trace to avoid probing a spot multiple times. The Biolever Mini (BLAC40TS) contilevers (Olympus; 10-nm nominal tip radius, sharpened probe) were indented with 180 pN, applying no additional dwell time. Cantilevers were chemically modified (compare with Sample Preparation) for Surface and Cantilevers) and were

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calibrated after the measurement using the equipartition theorem method (28). The datasets were recorded in the course of a few hours and contained around 50,000-90,000 curves saved in hdf5 files for additional data analysis.

Previous studies on smooth muscle myosin light-chain kinase (10) were conducted with 5,000-Da PEG linkers at a pulling speed of 800 nm/s. For this study, these parameters did not provide sufficient force resolution for clearly identifying contour length increments of the subtle peaks measured. By using shorter PEG linkers (35) (in this case, 425 Da) and thereby, increasing the loading rate as seen by the molecule, rupture forces of the detected events could be increased (*SI Appendix*, Fig. S13). These higher forces enable reliable worm-like chain fits for accurate analysis of the contour length increments.

Data Analysis. To show the characteristic unfolding patterns of the probed FAK construct, heatmaps were assembled with all curves that showed the correct total contour length (indicating correct site-specific attachment) as well as the presence of the characteristic unfolding peaks. Denoised (based on Savitzky–Golay, length 35, two polynomial for 800 nm/s) and length 21, two polynomial for 12,800 nm/s) force spectroscopy data were aligned manually in force-distance space (only by translating along the distance anis to account for length differences in PEG); they were binned from -15 to 150 nm in distance and from -15 to 150 pN in force for measurements with short PEG linkers (425.39 Da) and binned from -15 to 250 pN in force for Measurements with long PEG linkers (5,000 Da) to create a heatmap. The number of bins (equal for both distance and force axis) is dependent on the curves contained in the heatmap (Fig. 2, 150 bins); a, 3, 250 bins; Fig. 4, 250 bins; and Fig. 510 bins).

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The denoised data points (savitzky–Golay) in force–distance space were binned on the distance axis into 2.5-nm- (for Fig. 2), 3-nm (for Figs. 3) and 4), and 3.5-nm-sized slices (for Fig. 5) (moving the slice window by 0.2 nm each step), and their densities on the force axis (*y* axis) were estimated by a kernel density estimate (KDE) with a bandwidth of 0.2 pN (compare with *SI Appendix*, Fig. S14). The resulting most probable values are then assembled to form the most probable unfolding curve (shown as a black line in *SI Appendix*, Fig. S14). The resulting most probable values are then assembled to form the most probable unfolding curve (shown as a black line in 5*I Appendix*, Fig. S14) and analyzed to find the most probable unfolding peaks. The FWHMs of the distance slices were then taken as selection criteria for the unfolding peaks. It's lices contain a rupture event, the drop in force results in broad distributions, thereby clearly deviating from the noise level. The peaks were first detected by a simple peak detection of the distance slices. To be accepted as a peak, the FWHM of the distance slices has to be above the FWHM of the KDE of the accumulated PWHMs of the distance slices. To be accepted as a good representation of the noise level of the curve (compare with *SI Appendix*, Fig. S14). The procedure of assembling the most probable curve does not necessarily reproduce absolute rupture forces but yields a good result for the most probable and most representative pathway (36).

For additional analysis of the contour length increments, each stretch preceding an unfolding event is fitted with the worm-like chain model. This is done for every single curve contained in the heatmap. The most probable contour length for each peak is determined using a KDE. The increments between these most probable contour lengths were used to compare them with structural elements of the crystal structure.

with structural elements of the crystal structure. Data analysis was completely carried out in Python 2.7 and is available online together with all used datasets (https://gitlab.physik.uni-muenchen. de/Magrus.Bauer/fak.analysis).

**MD Simulations.** We use the FK-FAK construct developed previously (17) and solvate the protein in a 150 × 10 × 10-nm box. The total system contains or 1.5 million atoms, including 908 Na+ and 903 Cl-ions, corresponding to a ionic strength of 0.1 M. We use GROMACS (15), version 2016 for all of our simulations. As force field, we use Amber995B-ILDN\* force field (16) with Joung ions (27) and a transferable intermolecular potential with 3 points (TIP3P) water model (37). We use a time step of 2 fs and freeze all bonds in our simulations through a linear constraint solver (UINCS) procedure (8P Hess) of fourth order. Two Nosé-Hoover thermostats, one for protein and one for nonprotein atoms, were used with a time constant of 0.6 ps to keep pressure of 1 nn were used not expressibility of 4.5  $\times$  10<sup>-5</sup> with a reference pressure of 1.0 nm were used with an initial frequency of 0.03 ps. These parameters were automatically updated during the simulations by GROMACS for particle hest of such order.

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We perform a total of 30 pulling simulations, each with a spring constant of 830 pN/nm: 10 simulations each at the velocities of 1, 1/3, and 1/10 nm/ns. The simulations were performed in the presence of an ATP Into minus, the simulations were performed in the presence of an AIP molecule and Mg<sup>2+</sup> ion. To obtain single-domain pulling simulations, we start from the coordinates of the full FK-FAK construct and keep only the residues in the relevant domains. We relax the structures in 100-ns equi-librium simulations and solvate the FERM domain and the kinase domain in 67 × 9 × 9 and 100 × 8.5 × 8.5 nm, respectively. These correspond to 534,000 (FERM) and 706,000 (kinase) atoms. In both cases, we remove the ATP molecule and the  $Mg^{2+}$  ion from the simulation. We pull only using the fastest pulling velocity (1.0 m/s) and otherwise, keep all parameters unchanged.

We quantify domain unfolding by measuring distances between residues as follows: for FERM, we used residues 216-362; for the linker, we used residues 362-418, and for the kinase, we used residues 218-362; for the linker, we used the beginning and end of unfolding events as the times that the distances reach 10 and 45 nm for FERM, 7 and 15 nm for the linker, and 20 and 75 nm for the kinase. For the initial conformational activation, we also use a simple distance criterion: namely, whenever the distance 216-640 exceeds 10 nm. Since force profiles obtained in MD simulations include several intermediate

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ruptures, we identify peaks through a two-step procedure: (i) a Gaussian smoothing of the force profiles with an SD consistent with an extension of 0.1 nm and (ii) finding local maxima of the smoothed force profile in a window consistent with an extension of  $\pm 10$  nm.

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Event	MD (nm)	AFM (nm)	Assignment in AFM
			plots
Domain separation	10		first extension (prio
			to first unfolding)
Linker-F1 separation	12		first extension
F3 unfolding	30	29-32	f1a+f1b
F2 unfolding	13*	14-19	f1c
CK partial	10	13	k1
Kinase rest	70	68	k2
9 nm increase in the su included in the experime upplementary Table S1: Su istance changes) and AFM	imulations and an estin ental construct ummary of the length chai	ated 4 nm from a loop i ages observed in MD simul ath increments). Due to the	egion of the ybbR-ta
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Part III Appendix

## A List of Publications

- [1] C. Schoeler\*, K. H. Malinowska\*, R. C. Bernardi, L. F. Milles, M. A. Jobst, E. Durner, W. Ott, D. B. Fried, E. A. Bayer, K. Schulten, H. E. Gaub, and M. A. Nash, "Ultrastable cellulosome-adhesion complex tightens under load", Nature Communications, 2014, DOI: 10.1038/ncomms6635
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- [4] E. Durner, W. Ott, M. A. Nash, and H. E. Gaub, "Post-translational sortasemediated attachment of high-strength force spectroscopy handles", ACS Omega, 2017, DOI: 10.1021/acsomega.7b00478
- [5] W. Ott\*, E. Durner\*, and H. E. Gaub, "Site-Specific Protein Coupling Strategies for Surface-Based Binding Assays", Angewandte Chemie, 2018, DOI: 10.1002/ange.201805034
- [6] M. S. Bauer, F. Baumann, C. Daday, P. Redondo, E. Durner, M. A. Jobst, L. F. Milles, D. Mercadante, D. A. Pippig, H. E. Gaub, F. Gräter, and Daniel Lietha, "Structural and mechanistic insights into mechanoactivation of focal adhesion kinase", Proceedings of the National Academy of Sciences, 2019, DOI: 10.1073/pnas.1820567116
- [7] R. C. Bernardi\*, E. Durner\*, C. Schoeler, K. H. Malinowska, B. G. Carvalho, E. A. Bayer, Z. Luthey-Schulten, H. E. Gaub, and M. A. Nash, "Mechanisms of Nanonewton Mechanostability in a Protein Complex Revealed by Molecular Dynamics Simulations and Single-Molecule Force Spectroscopy", Journal of the American Chemical Society, 2019, DOI: 10.1021/jacs.9b06776
- [8] E. Durner, M. S. Bauer, M. A. Jobst, W. Ott, L. F. Milles, and H. E. Gaub, "Conformational Transitions of a Protein Fold Observed by Dynamic Single-Molecule Force Spectroscopy", under review at Physical Review Letters

\* These authors contributed equally to this work

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